

Challenges from experiment



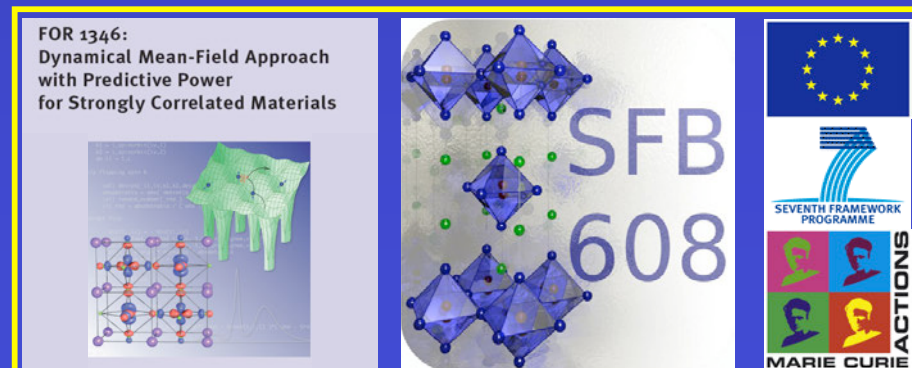
Hao Tjeng

Physics of Correlated Matter

*Max-Planck-Institute
Chemical Physics of Solids
Dresden*



- Zhiwei Hu, Jonas Weinen, Yi-Ying Chin, Nils Hollmann, Stefano Agrestini – *Dresden*
- Maurits Haverkort – *MPI FKF, Stuttgart*
- Thomas Koethe, Hua Wu, Roger Chang, Tobias Burnus, Christian Zobel, Marco Reuther, Helena Hartmann, Daniel Khomskii, Markus Braden, Thomas Lorenz, – *Cologne*
- Arata Tanaka – *Hiroshima, Japan*
- J. Cezar, N. Brookes – *ESRF, France*
- H.-J. Lin, K.-D. Tsuei, H.-H. Hsieh, C.T. Chen – *NSRRC, Taiwan*



Correlated electron systems:

- how does the spectral weight distribution change as a function of U / W ?

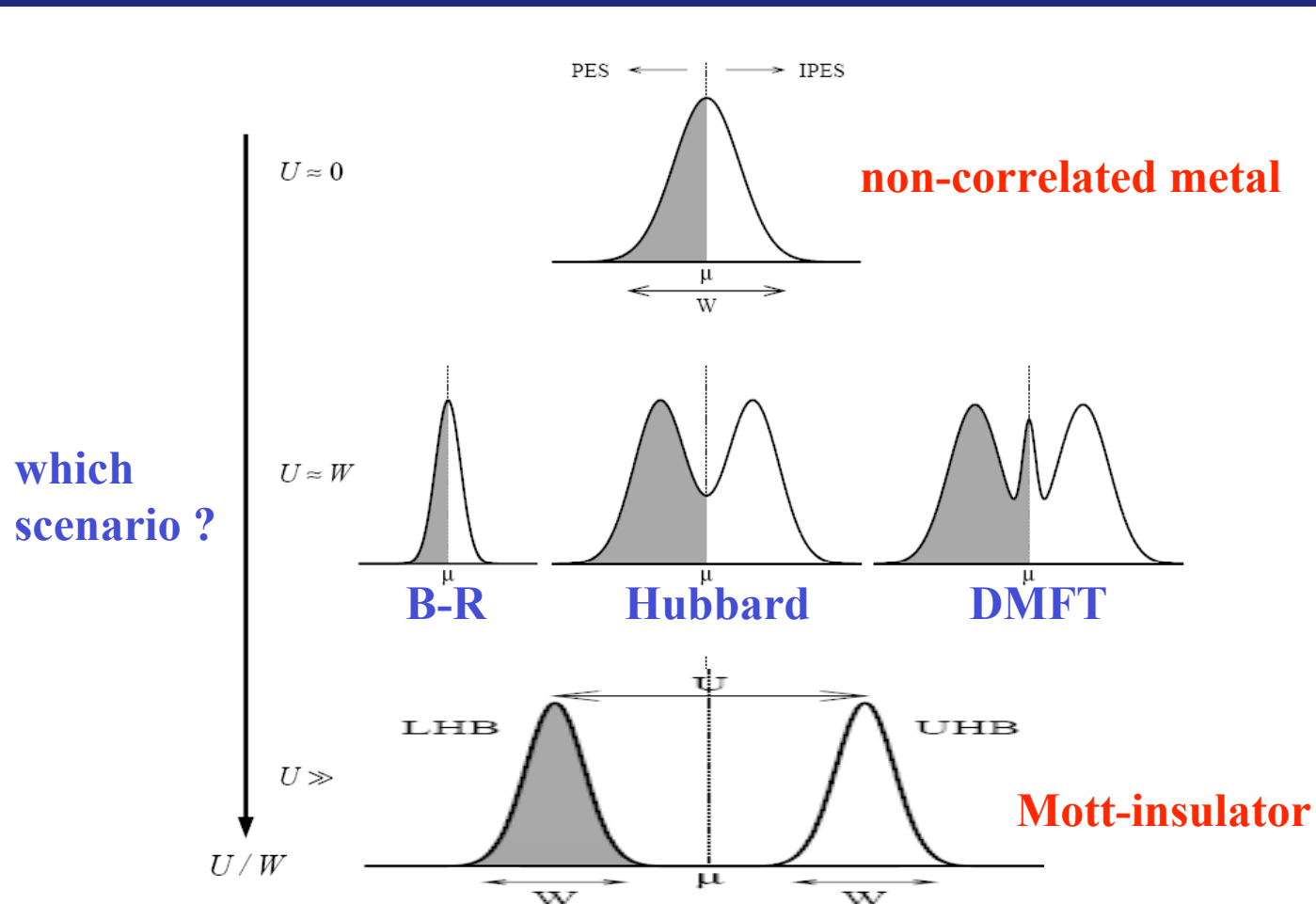


Figure 1.5: Photoemission and inverse photoemission spectra as expected in a Hubbard model for different values of U/W . The grey shading indicates occupied states. Adapted from Morikawa *et al.* [61]. *Phys. Rev. B* 54, 8446 (1996).

Transition metal oxides: wide range of properties

• d ¹ :		Ti ₂ O ₃	VO ₂	LaTiO ₃	
• d ² :	TiO	V ₂ O ₃	CrO ₂	YVO ₃	
• d ³ :	VO	Cr ₂ O ₃		LaCrO ₃	
• d ⁴ :				LaMnO ₃	Ca ₂ RuO ₄
• d ⁵ :	MnO	Fe ₂ O ₃		BaCoO ₃	
• d ⁶ :	FeO			LaCoO ₃	NaCoO ₂
• d ⁷ :	CoO			LaNiO ₃	
• d ⁸ :	NiO				La ₂ NiO ₄
• d ⁹ :	CuO				La ₂ CuO ₄

+ doping

- atomic multiplet interactions
- nearest neighbor interactions

→ determine effective U, W
and more !!!

(difficult to make a general scheme – each compound form ist own universe)

I. Importance of atomic multiplet interactions

Example: cobaltates

(a material class with very rich physics)

- Metal-Insulator-Transitions, Colossal-Magneto-Resistance, Thermoelectric power in:
 - layered: $\text{RBaCo}_2\text{O}_{5+x}$ (R=Eu,Gd)
 - “113”: $\text{La}_{1-x}\text{A}_x\text{CoO}_3$ (A=Ca,Sr,Ba)
- Superconductivity: $\text{Na}_x\text{CoO}_2 \cdot y\text{H}_2\text{O}$
- Unusual magnetic properties: $\text{Ca}_3\text{Co}_2\text{O}_6$ (*Ising/steps*)
- Multiferroicity: $\text{Ca}_3\text{MnCoO}_6$ (*Ising*)

What is special about Co ions ?

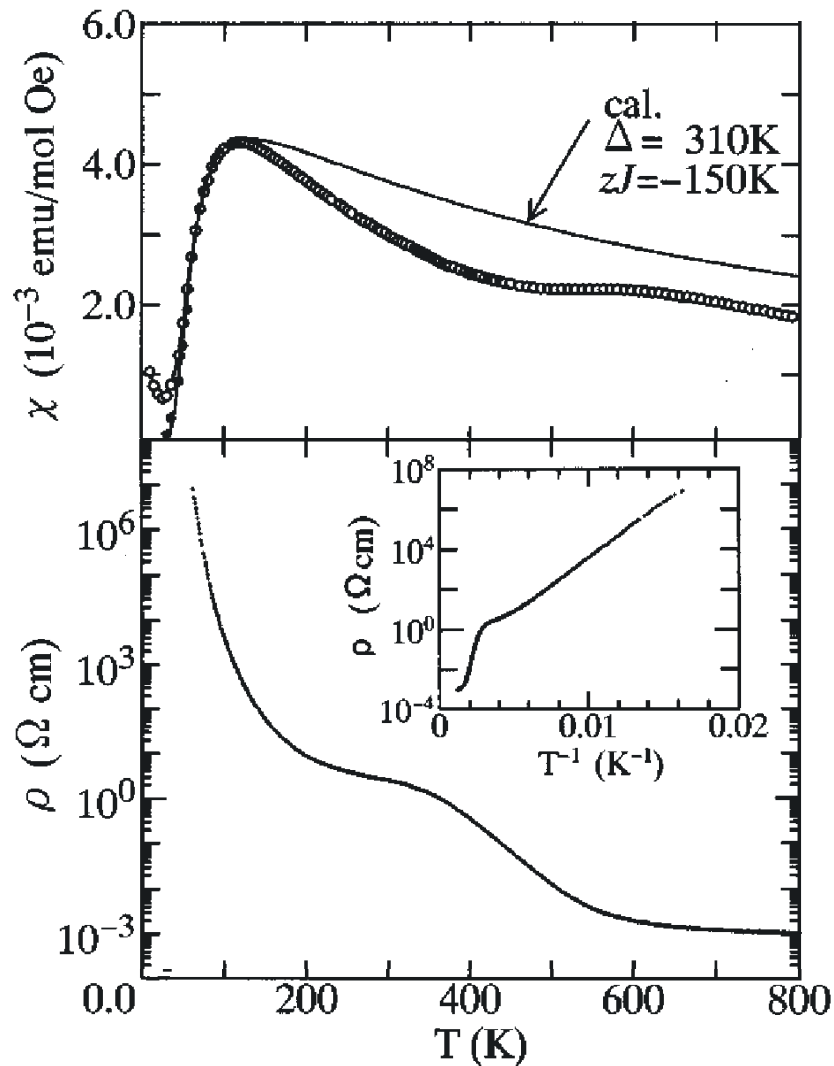
- valence state : 2+ 3+ 4+
- spin state: HS HS/IS/LS HS/IS/LS
- spin state transitions and *spin-blockade* mechanism
- local coordination: oct/pyramidal/trig/prismatic-trig
- *spin-orbit interaction*

Confusion in the literature about:

- valence, orbital, and spin state
- all kinds of possible scenarios are proposed from neutron and magnetic susceptibility experiments
- band theories give conflicting results

LaCoO₃ : a benchmark system

Co³⁺: 3d⁶

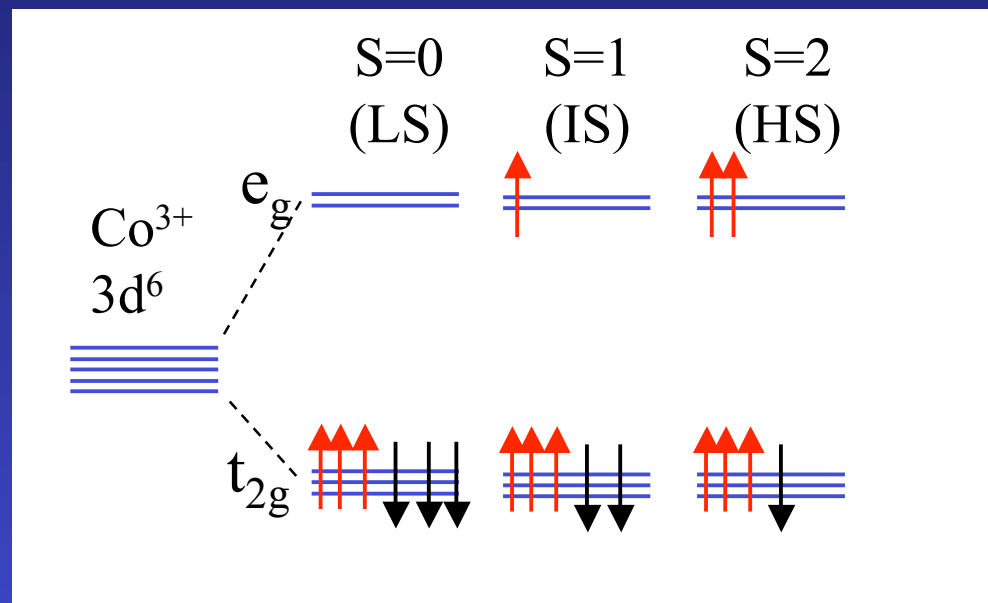


S. Yamaguchi *et al.*, PRB **53**, 2926 (1996)

- non-magnetic insulator at low T
- non-magnetic to paramagnetic transition for $T > 25$ K, with max. in magn. susceptibility at 100 K
- resistivity drop $T = 350$ K - 550 K, “metal-insulator transition”

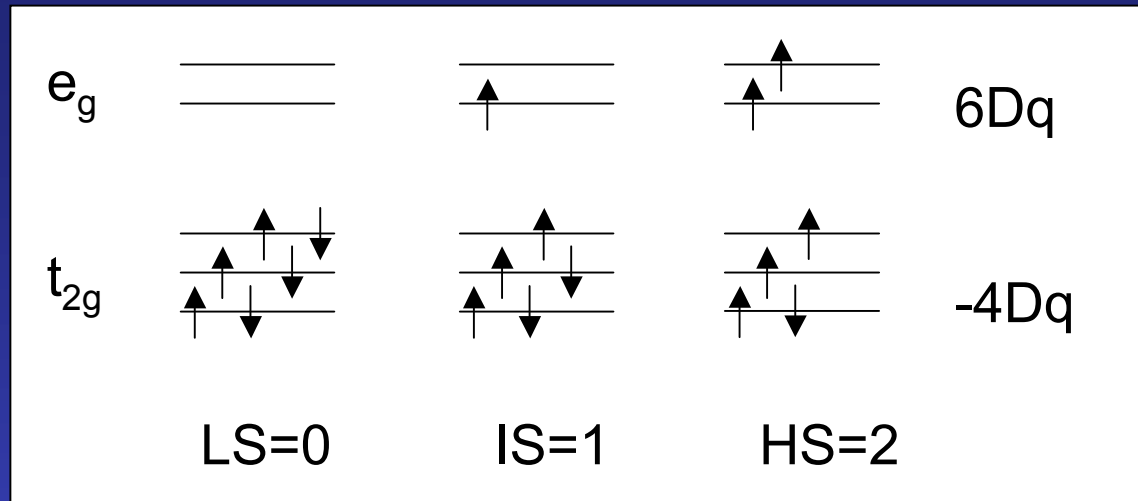
Spin-state transitions ?
Low – Intermediate – High spin ?

Puzzle: what is the spin state of Co^{3+} ??



competition:
crystal field - band formation - Hund's exchange

Energetics of 3d⁶ spin states: ionic picture



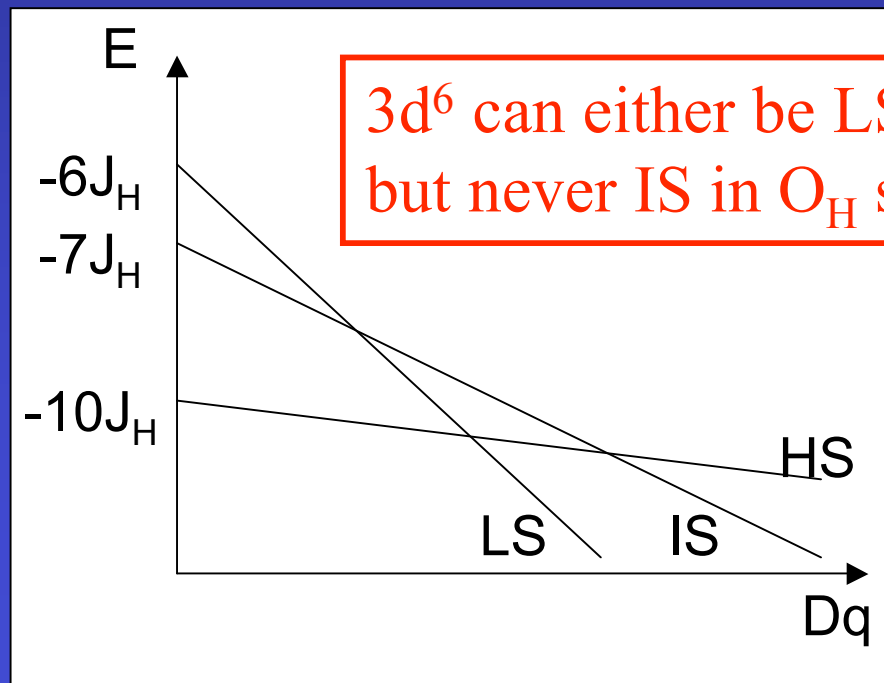
Simplest scheme:

$$E_{LS} = -24Dq - 6J_H$$

$$E_{IS} = -14Dq - 7J_H$$

$$E_{HS} = -4Dq - 10J_H$$

LS-HS crossing occurs
at $10Dq = 2 J_H \approx 1.6 \text{ eV}$



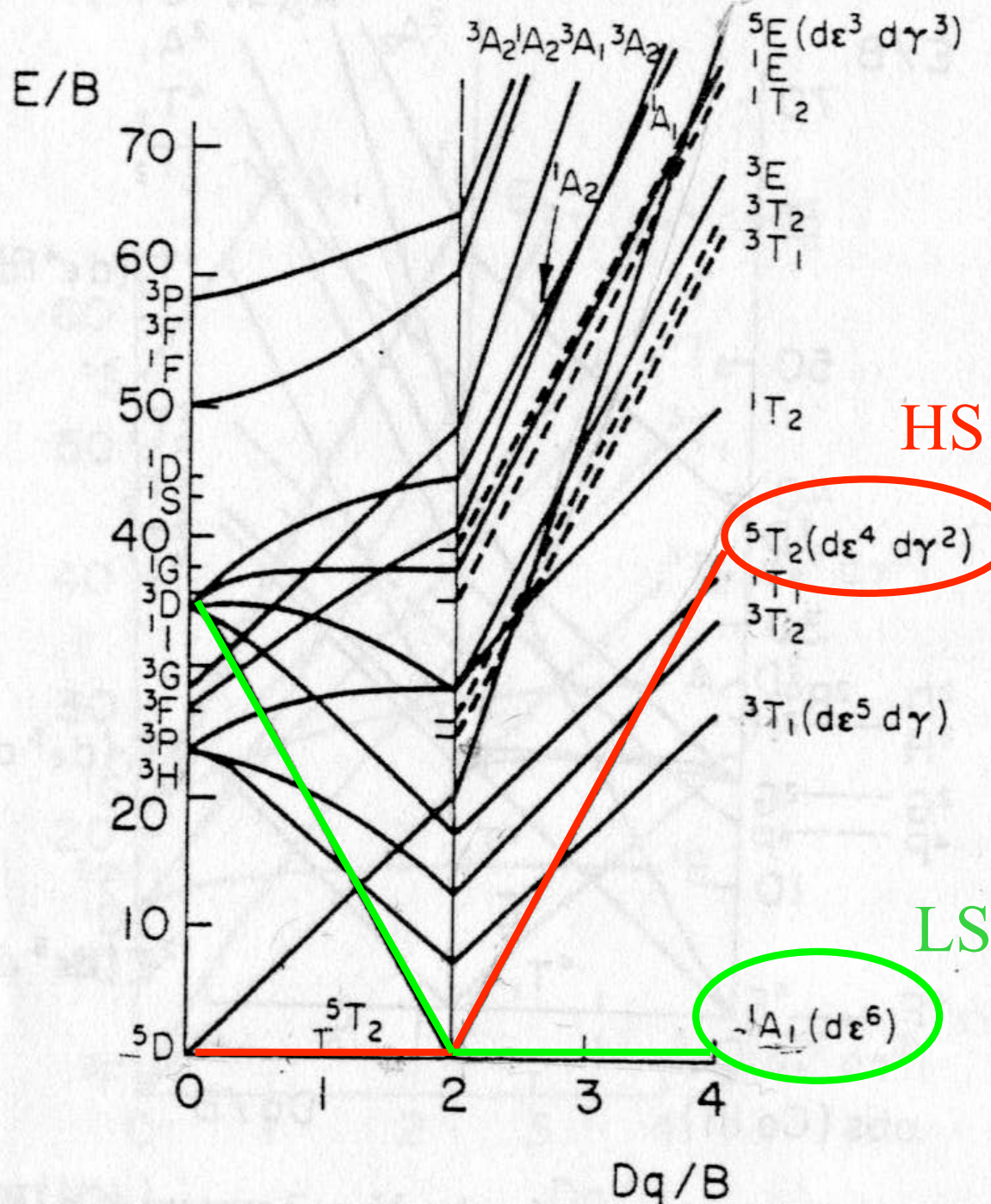
Energetics of $3d^6$ spin states: an ionic picture

full atomic
multiplet theory

Sugano, Tanabe,
Kamimura - 1970

$3d^6$ can either be
LS or HS
but never IS in O_H

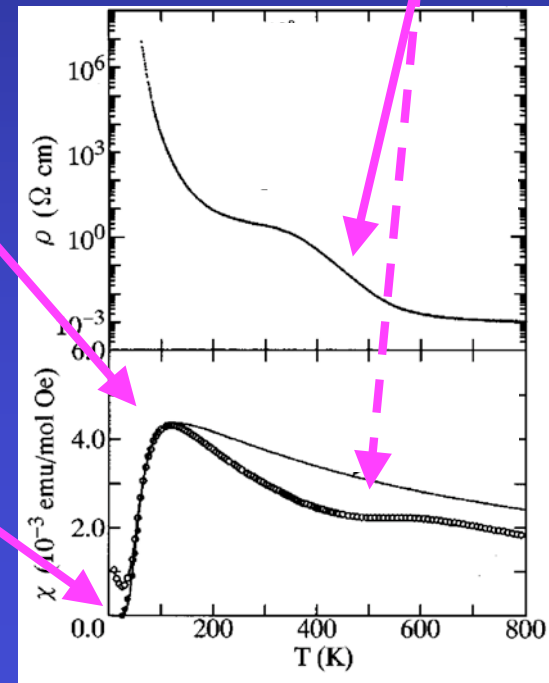
LS-HS crossing at
 $10Dq=20B \approx 2.2\text{eV}$



Classic model: LS-HS scenario

J.B. Goodenough *et al.* (1960s) proposed:

- LS ground state non-magnetic at $T = 0\text{K}$
- increasing paramagnetism through thermal population of HS, $E_{\text{HS}} - E_{\text{LS}} \leq 80\text{meV}$
- reduced resistivity through ordering of LS-HS Co-sites at higher temperatures



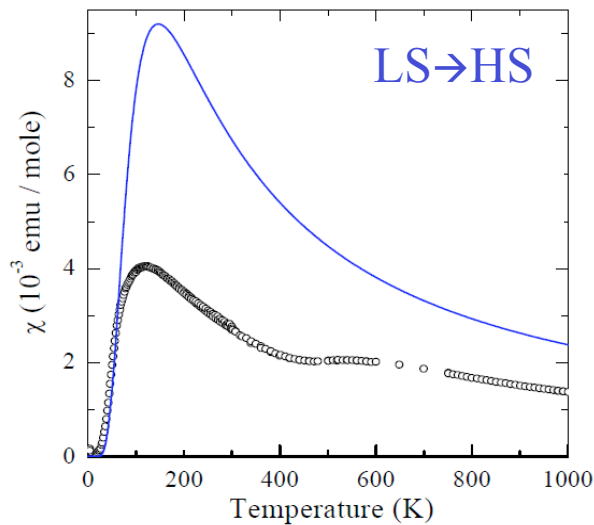
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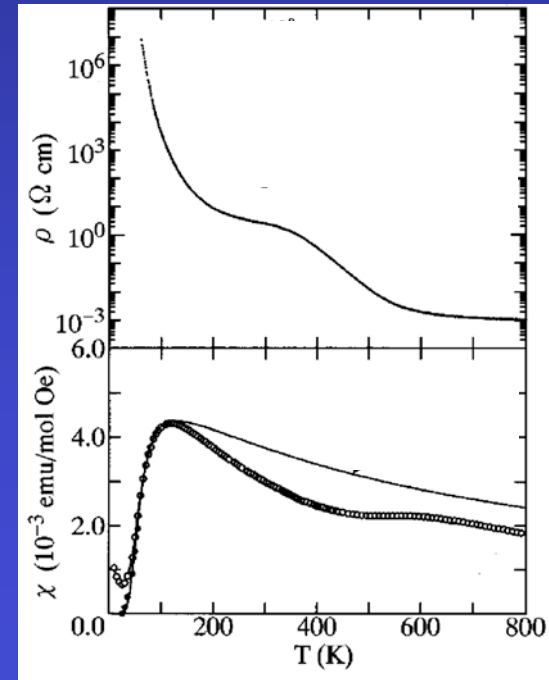
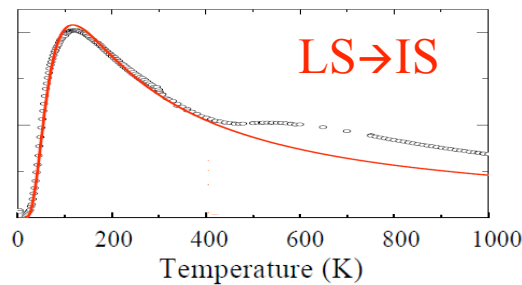
- LS ground state non-magnetic at $T = 0\text{K}$
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- reduced resistivity through ordering of LS-HS Co-sites at higher temperatures

T. Lorenz

But LS→HS does not fit quantitatively !!!



LS→IS better ???



Modern model: LS-IS scenario

Remark:

Kor

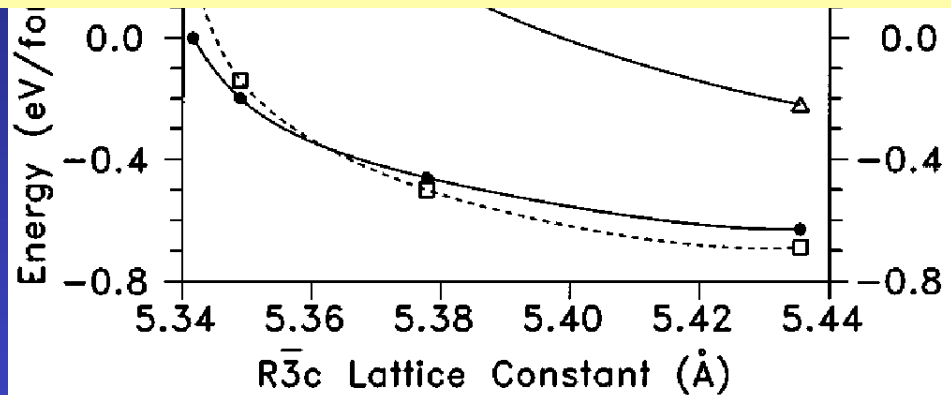
Saw

• ba

• spin state transition from LS to *orbitally ordered* IS at 150K

• metallic *disordered* IS at higher temperatures

LDA+U cannot implement full-multiplet effects $U_{ijkl} \hat{d}_i^\dagger \hat{d}_j \hat{d}_k^\dagger \hat{d}_l$
 LDA+U is a single Slater determinant approach: $U_{ij} \hat{d}_i^\dagger \hat{d}_i \hat{d}_j^\dagger \hat{d}_j$
 → incorrect energies for especially IS !!



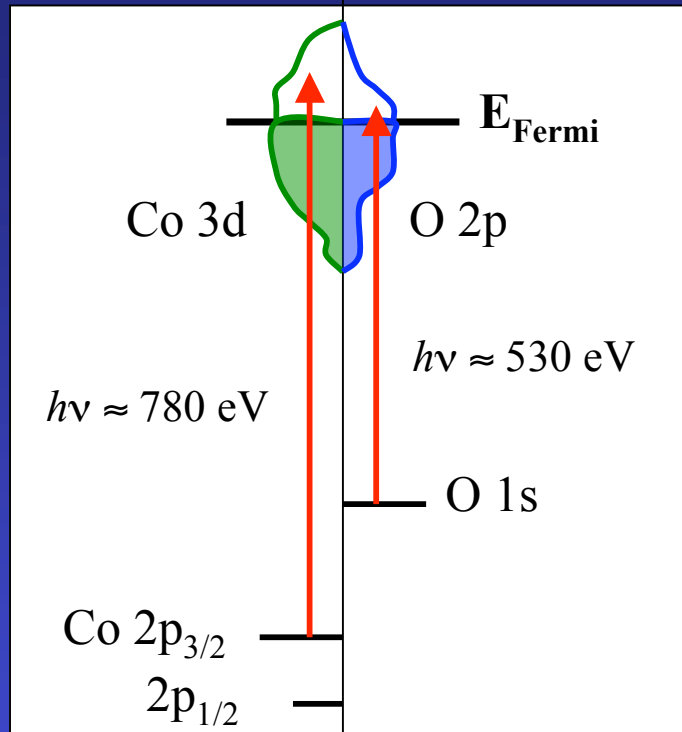
LDA + U: PRB 54, 5309 (1996)

Remark: results may strongly depend on details of calculations, e.g.

- assumed magnetic ordering
- value of Hubbard U parameter
- coexistence of spin states not considered
- **incomplete implementation of multiplet effects**

How to determine spin state of Co ions:

Soft-X-Ray Absorption and Magnetic Circular Dichroism (spectroscopically-resolved magnetic susceptibility measurement)



- use of core levels → local transitions → element and site specific
- involves most relevant orbitals:
 - $2p-3d$ (TM), $3d-4f$ (RE), $1s-2p$ (O,N,C)
- dipole allowed → very strong intensities
- dipole selection rules + multiplet structure give extreme sensitivity to symmetry of initial state: charge, spin and orbital

theory:

TM $2p-3d$: Cluster calculations with full atomic multiplet theory
 O $1s-2p$: LDA+U calculations

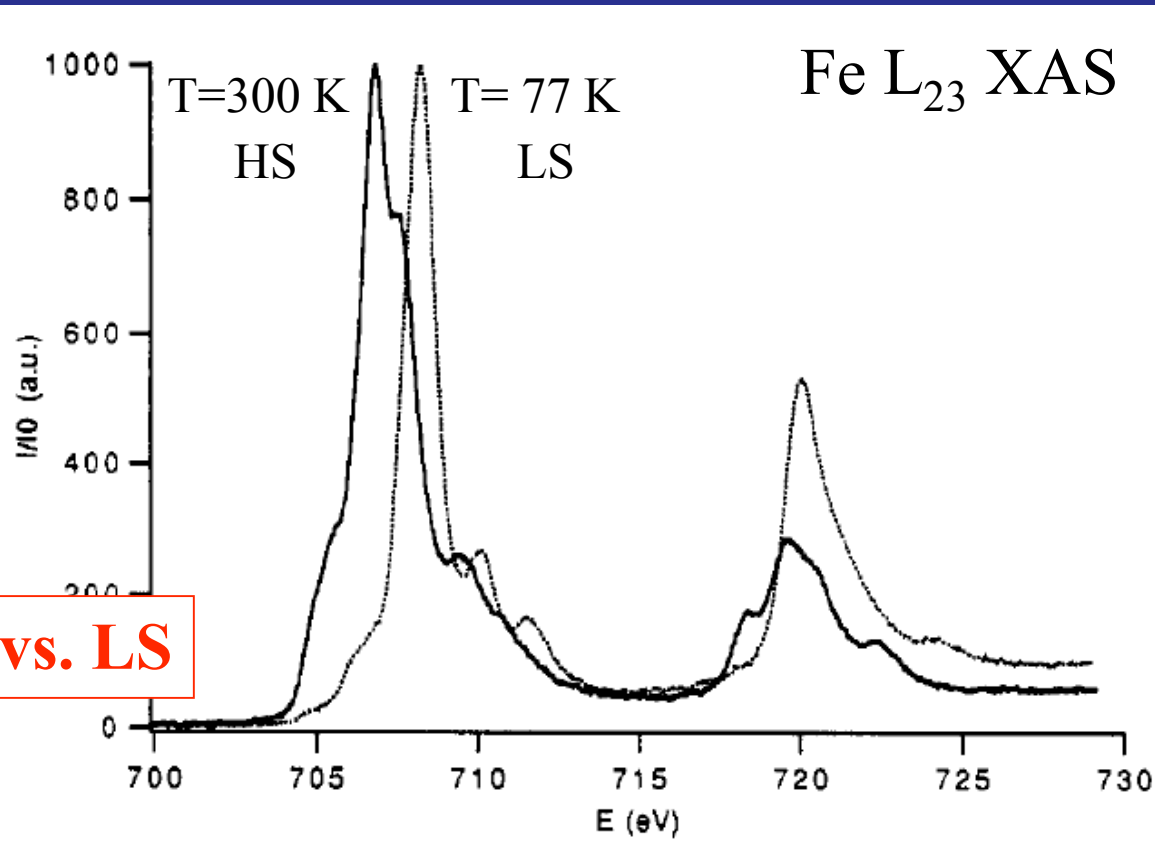
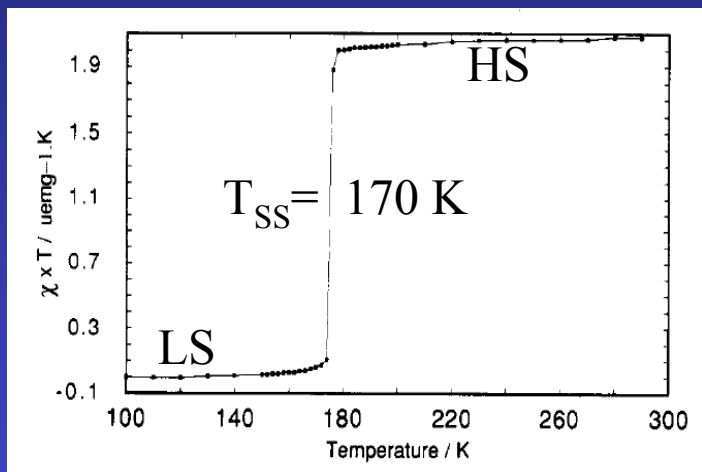
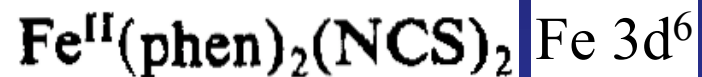
Technique developed in late 1980's:

- Fink, Sawatzky, Fuggle
- Thole, van der Laan
- Chen, Sette

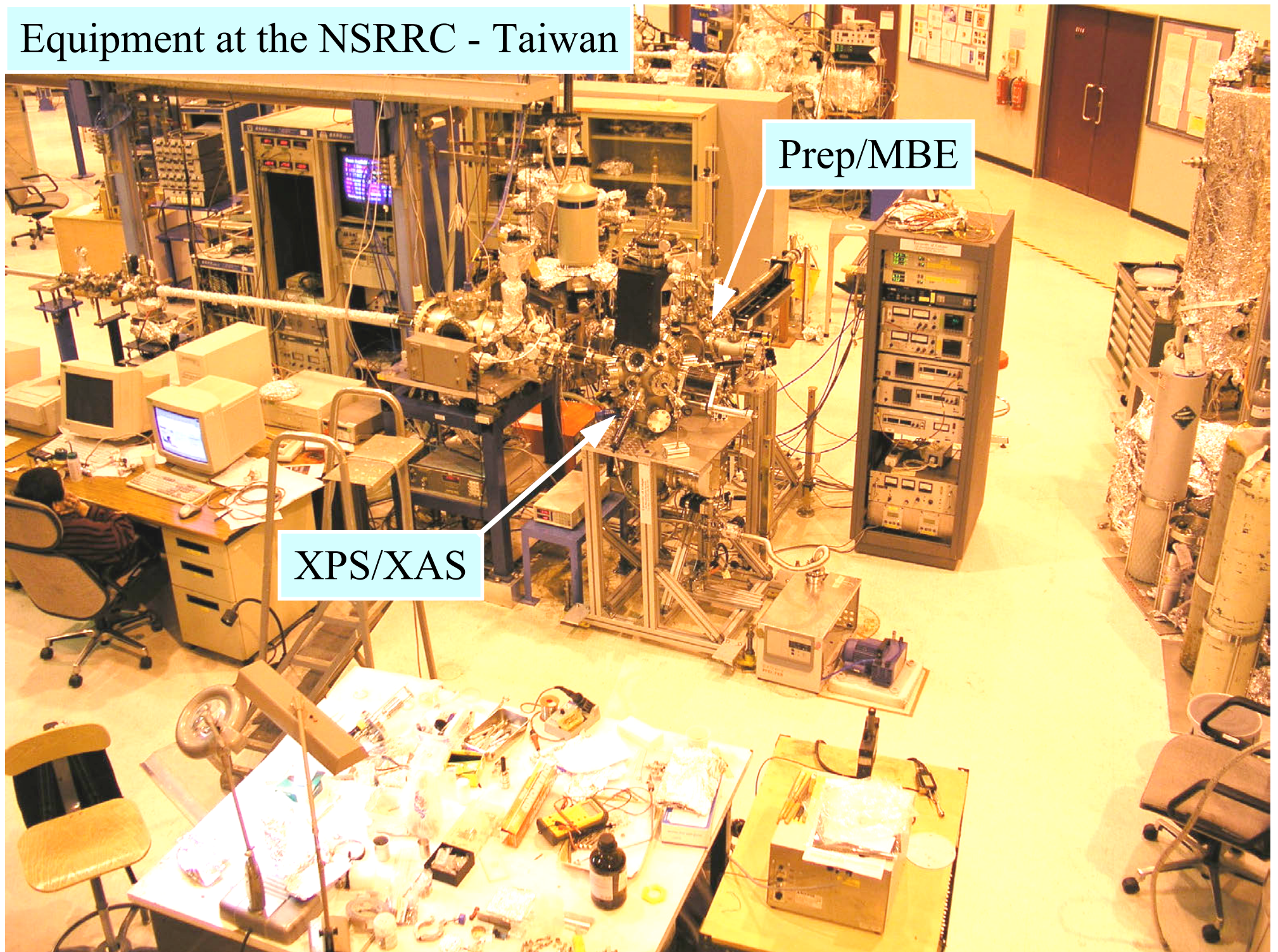
$$\text{Spectrum } (h\nu) = \sum_f |\langle i | e.r | f \rangle|^2 \delta(h\nu - E_f + E_i)$$

$|i\rangle$ = initial state, $|f\rangle$ = final state
e.r = dipole transition.

If experimental spectrum is simulated correctly,
then also initial state and final states are known.

Spin Transition Evidenced by Soft X-ray Absorption SpectroscopyChristophe Cartier dit Moulin,^{*,†} Petra Rudolf,^{‡§} Anne-Marie Flank,[†] and Chien-Te Chen[‡]*LURE, Bâtiment 209d, F-91405 Orsay, France, and AT&T Bell Laboratories, Murray Hill, New Jersey 079074 (Received: January 14, 1992; In Final Form: March 31, 1992)***XAS: sensitive to HS vs. LS**

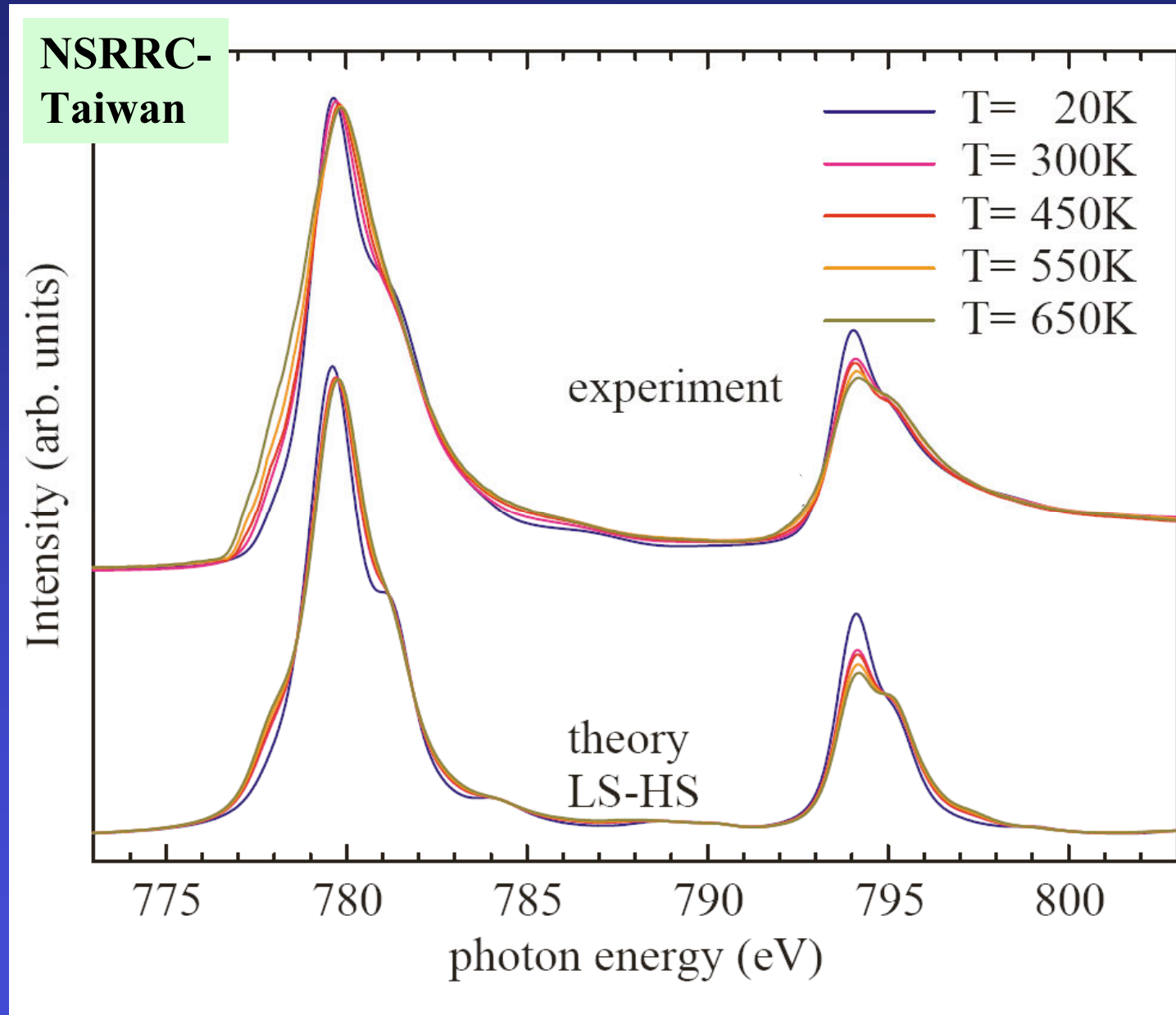
Equipment at the NSRRC - Taiwan



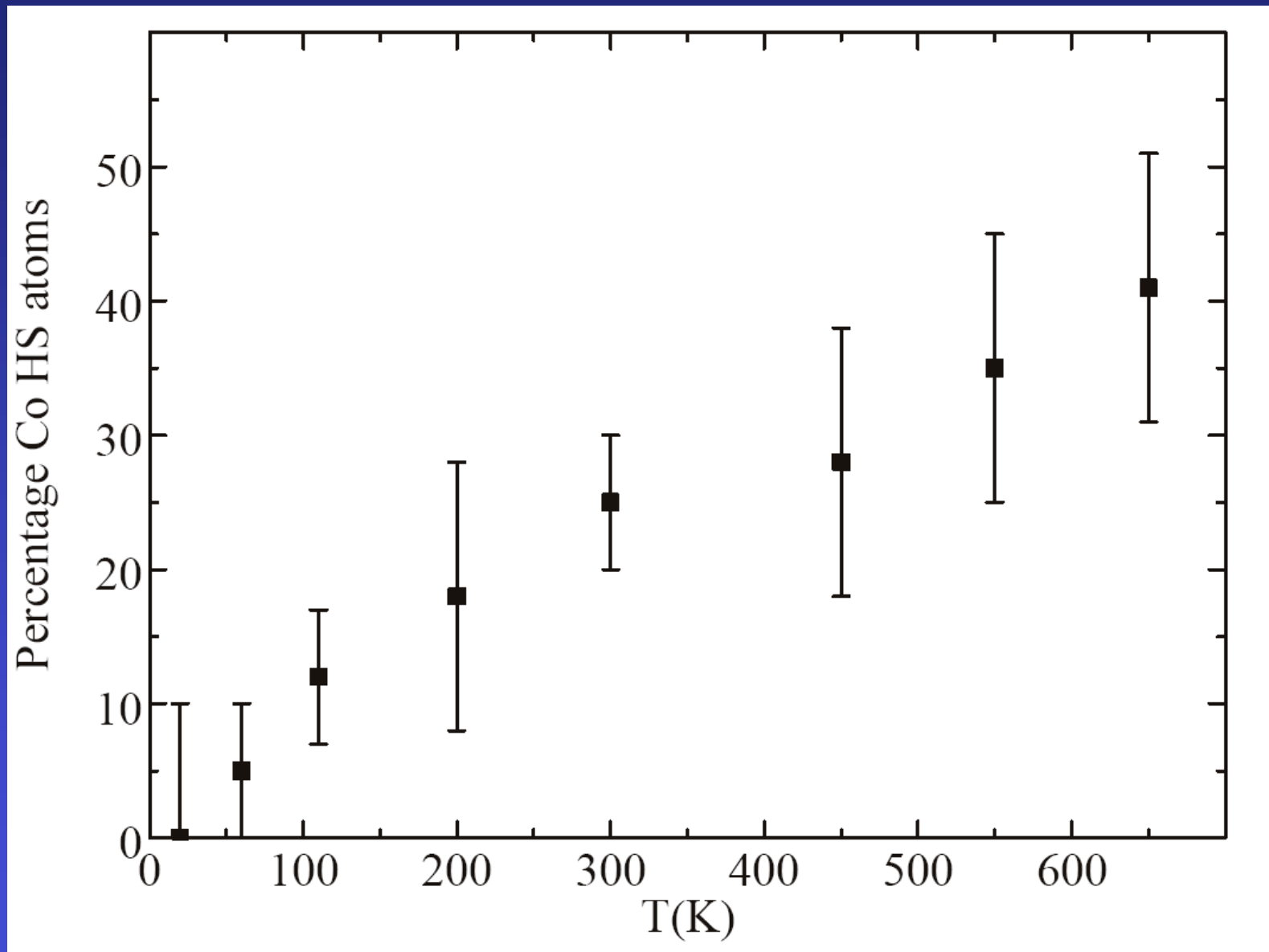
Prep/MBE

XPS/XAS

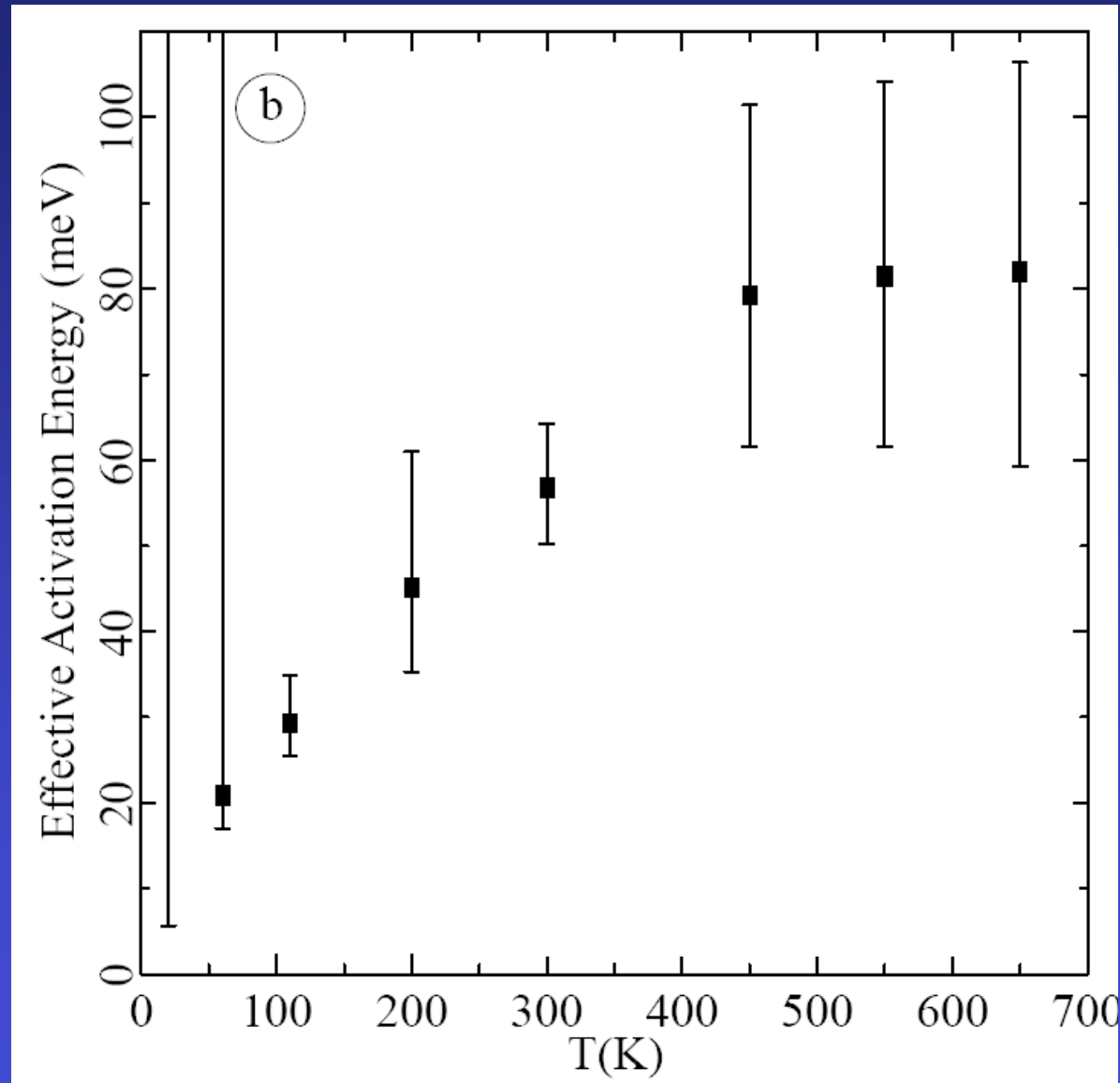
XAS study on the spin state of Co^{3+} ion in LaCoO_3



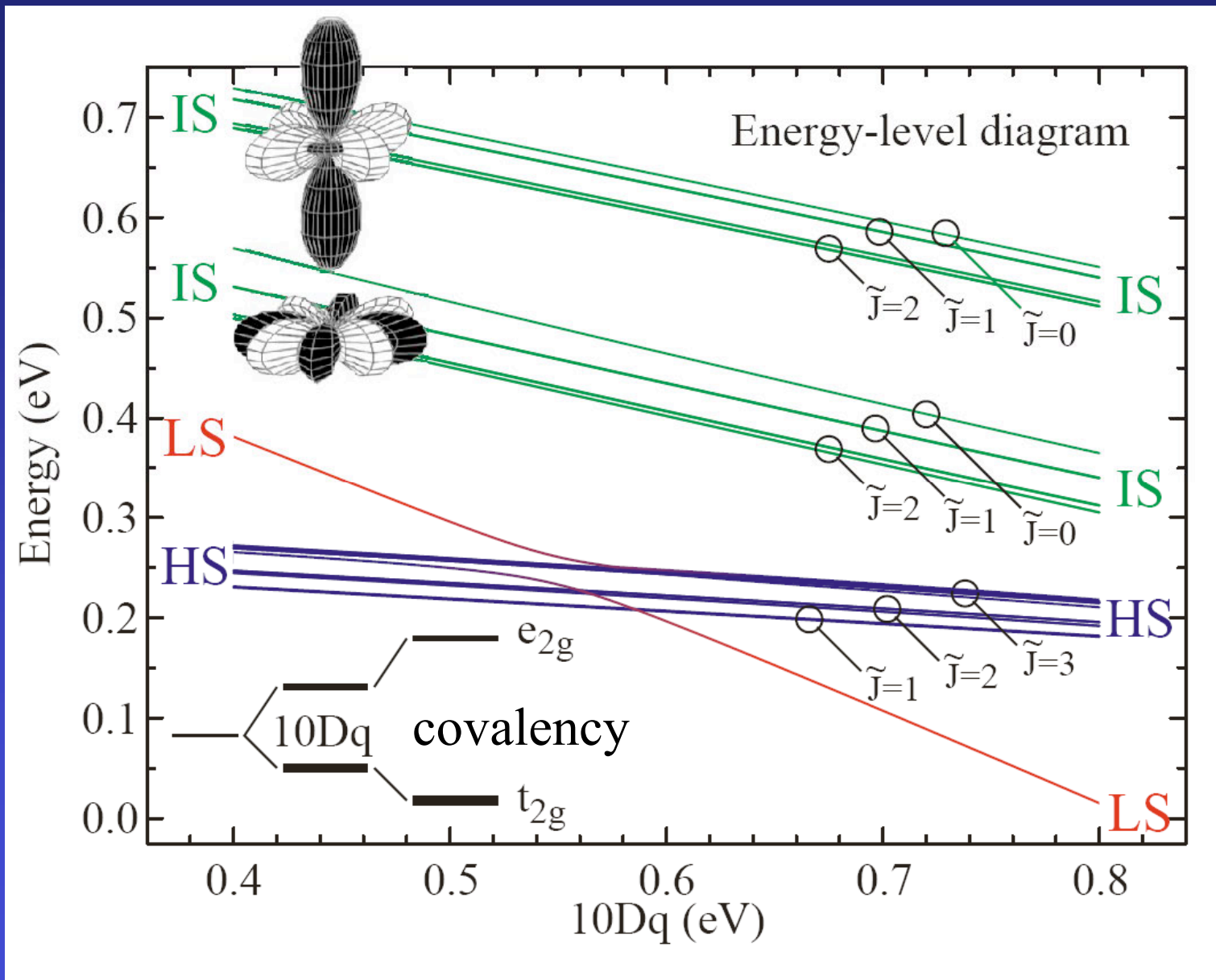
LaCoO₃: occupation of HS in the LS-HS scenario



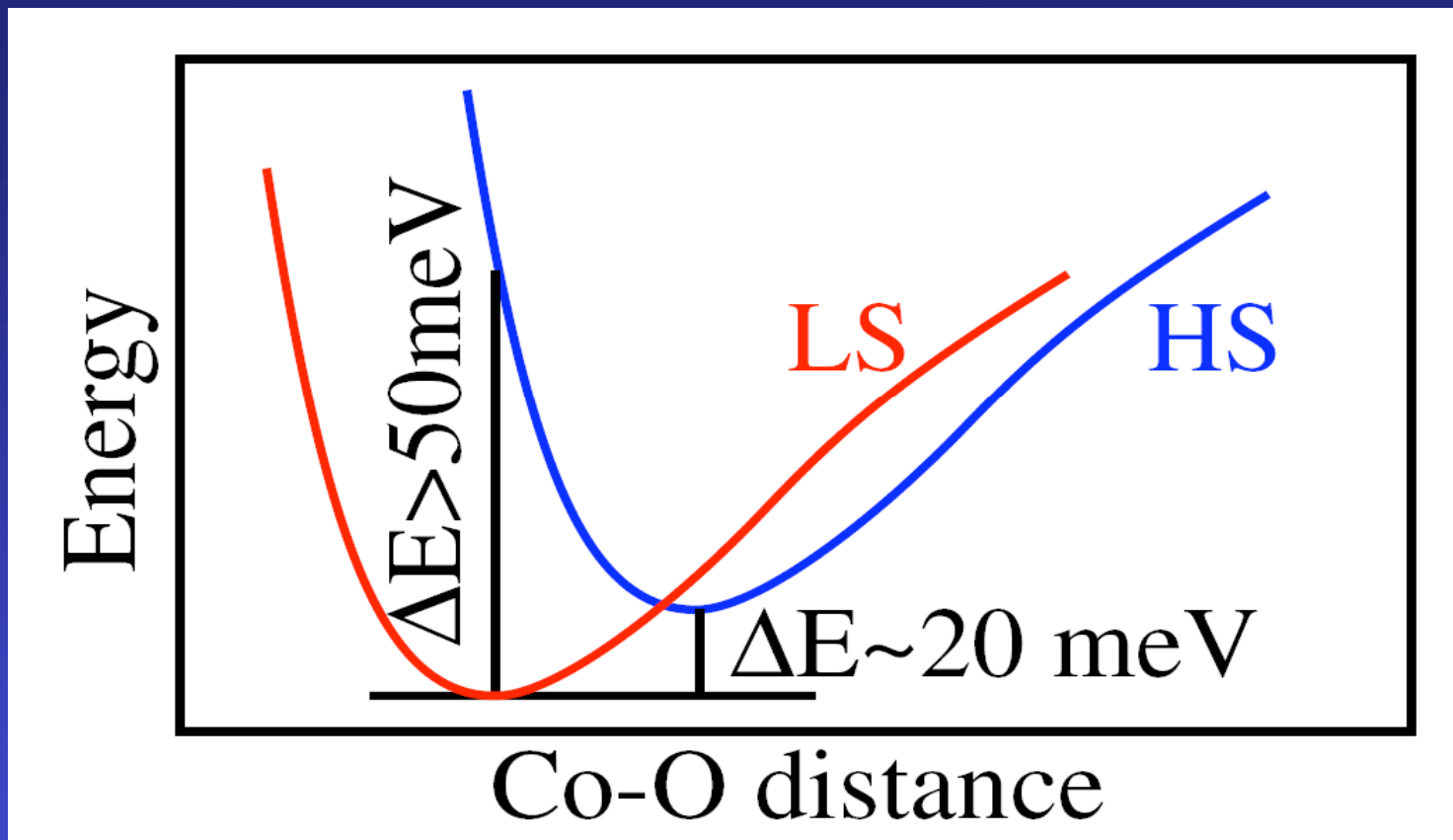
Effective activation energy: temperature dependent !



Energy level diagram: CoO_6 cluster incl. covalency

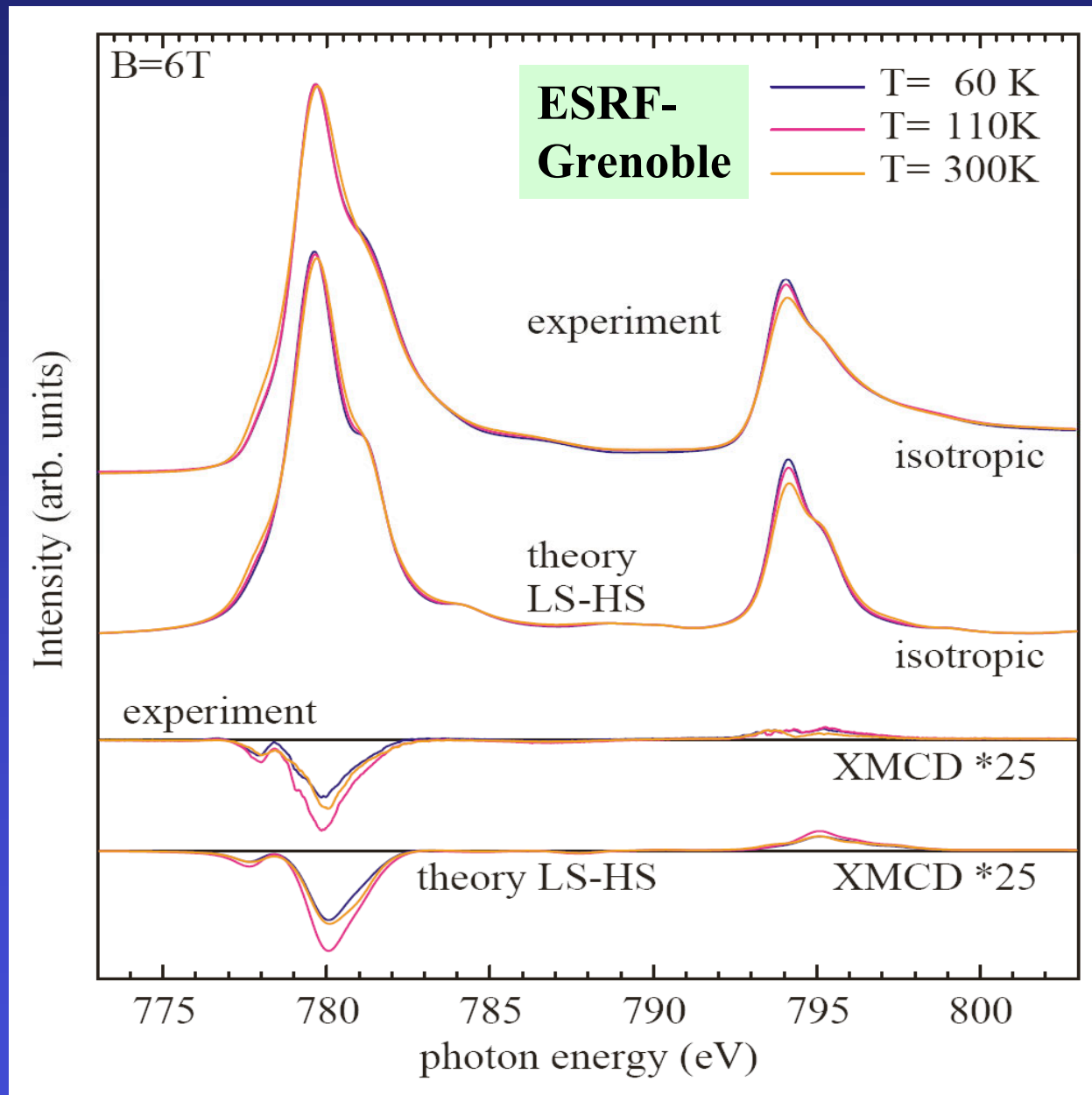


Spin state transition: local lattice relaxation

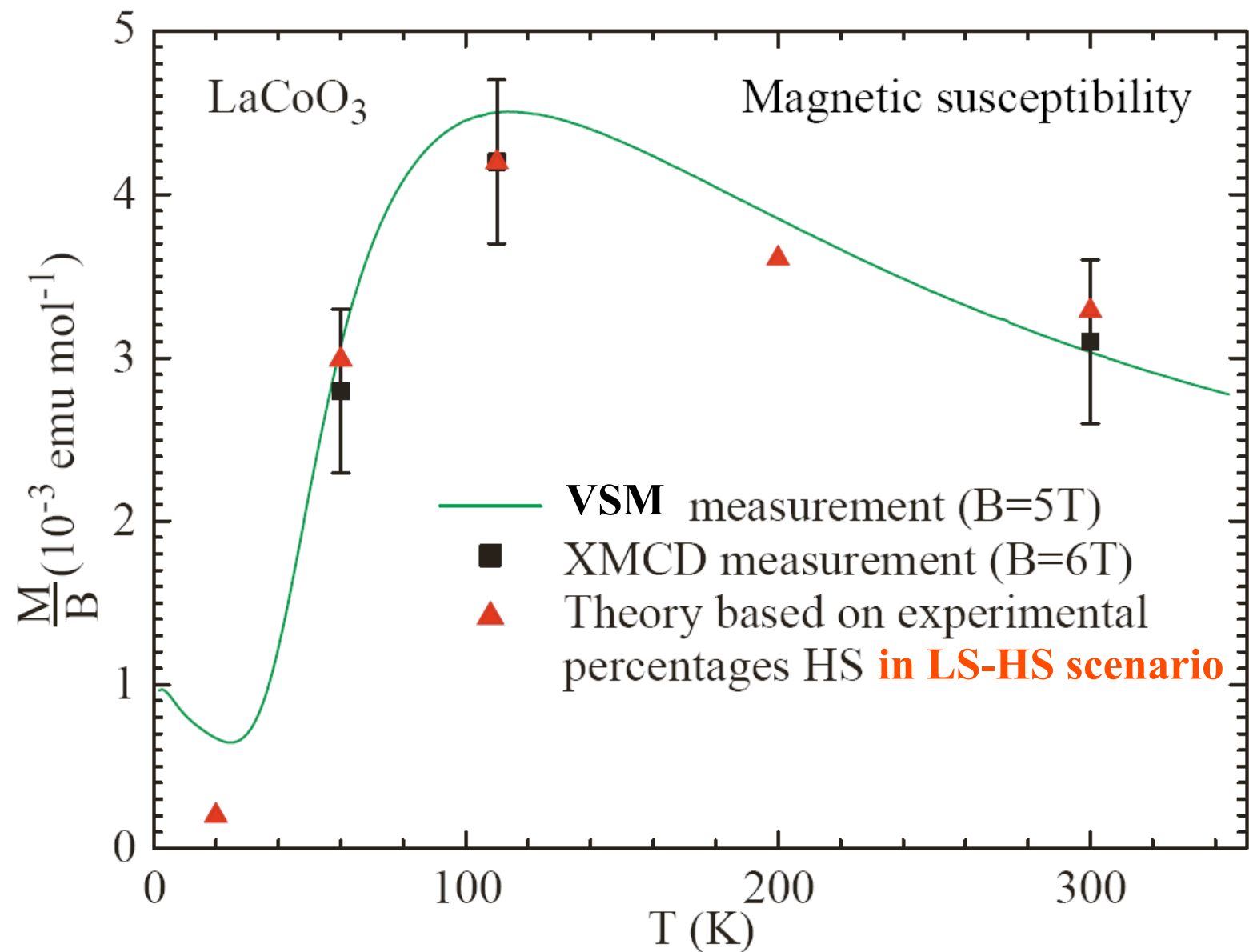


- frozen lattice: $\Delta E \gg k_B T$
otherwise too much Van Vleck and incorrect XAS spectra
- inhomogeneous mixed spin-state system

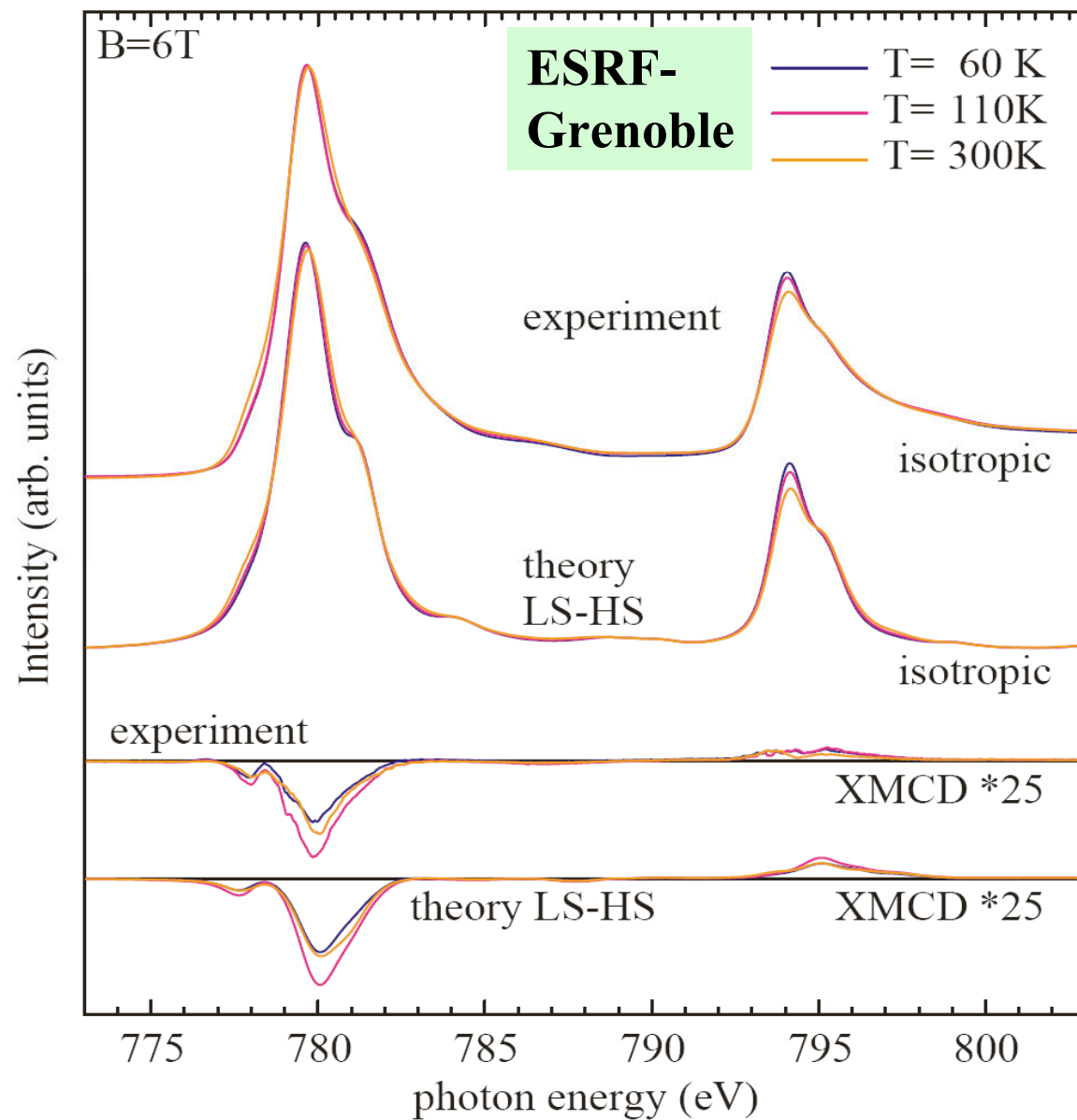
Magnetic Circular Dichroism: *paramagnetic* LaCoO₃



Magnetic susceptibility – MCD - XAS

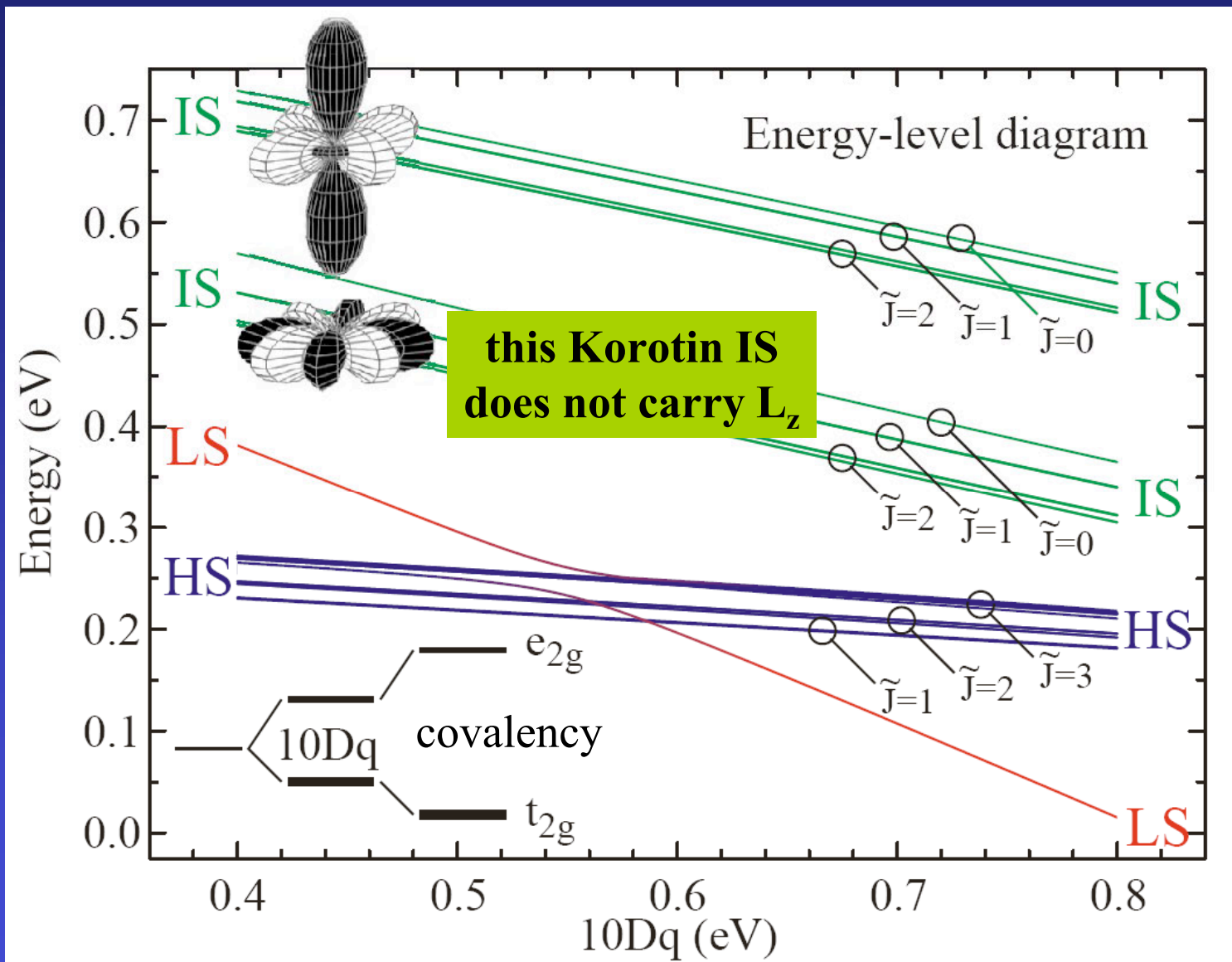


Can we rule out the LS-IS scenario ?!



crucial finding:
orbital/spin-moment
sum-rules:
 $L_z/S_z \approx 1/2$

Energy level diagram: CoO_6 cluster incl. covalency



Very difficult to find an IS state with L_z which also fits the XAS/MCD spectra

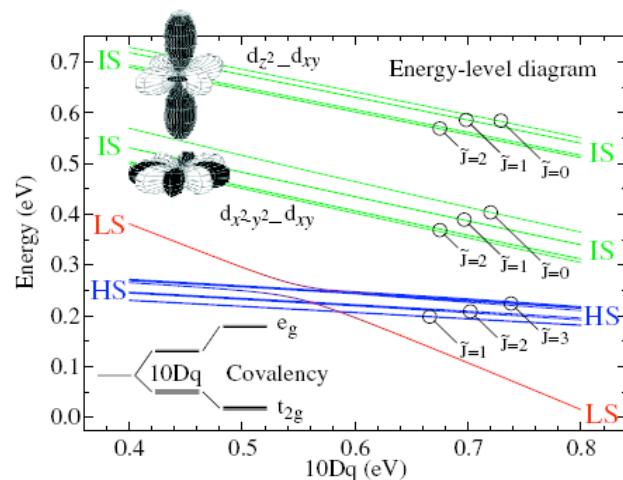
Summary: magnetic properties of LaCoO_3

- LaCoO_3 : non-magnetic at low T $[\text{Co}^{3+} = 3d^6]$
paramagnetic $T > 25 \text{ K}$
- 1960's : Low-Spin ($S=0$) \rightarrow High-Spin ($S=2$)
- 1996 : Low-Spin ($S=0$) \rightarrow Intermediate-Spin ($S=1$)

Korotin, Ezhov, Solovyeu, Anisimov, Khomskii, Sawatzky, PRB 54, 5309 (1996): LS-IS

claimed experimental proof: thermodynamic properties
best explained by *triply* degenerate excited state

Spectroscopy XAS/MCD

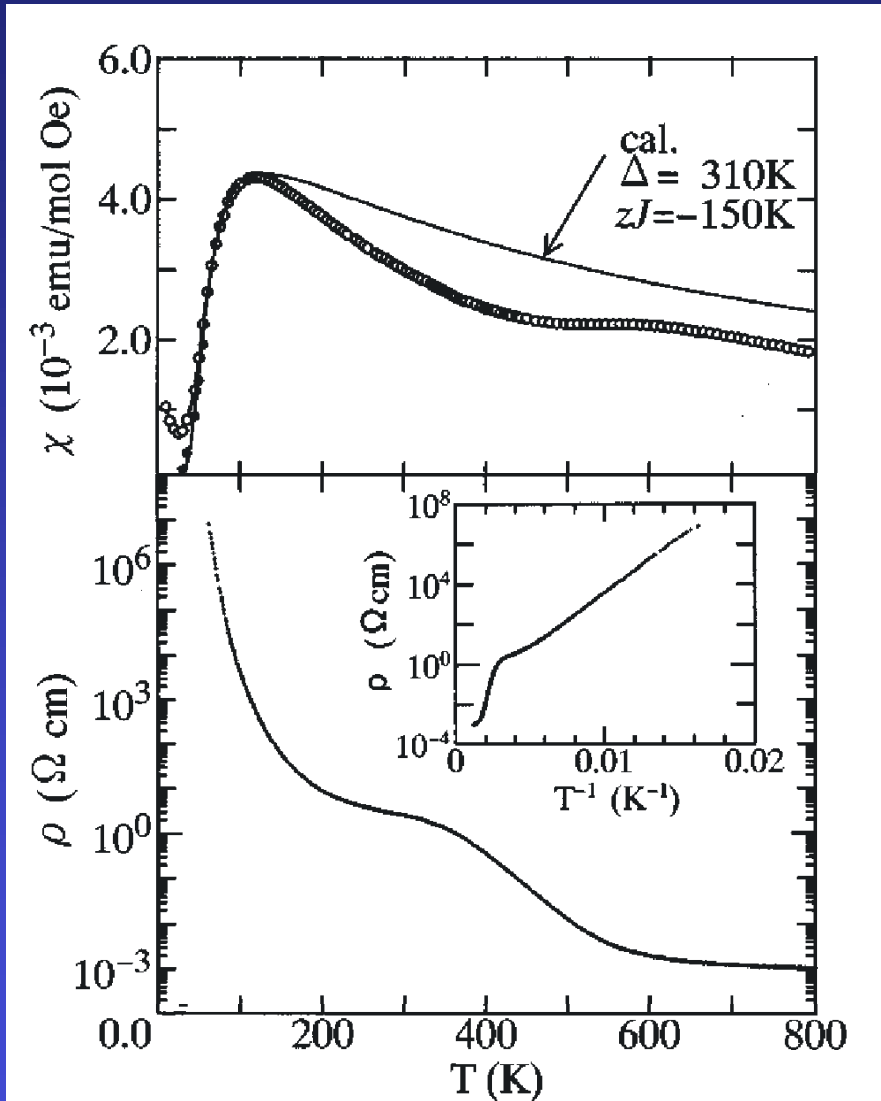


- Low-Spin ($S=0$) to High-Spin ($S=2$)
- temperature dependent activation energy

Full multiplet theory:

- High-Spin ($S=2$) is 3-fold degenerate !!
- Interm.-Spin ($S=1$) is 5-fold degenerate !!
- due to spin-orbit interaction

Temperature dependence of the resistivity of LaCoO_3



S. Yamaguchi *et al.*, PRB **53**, 2926 (1996)

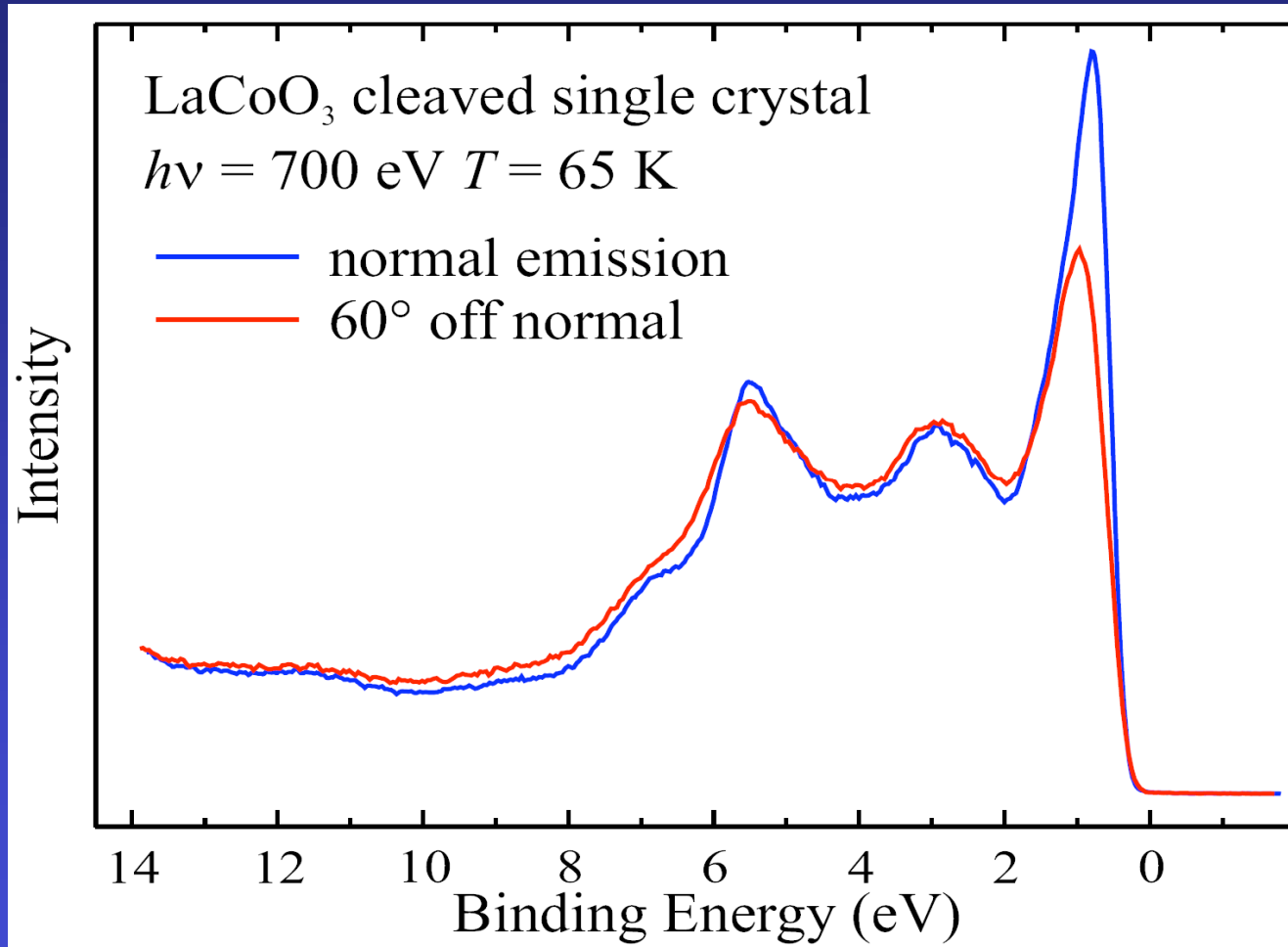
- non-magnetic insulator at low T
- non-magnetic to paramagnetic transition for $T > 25\text{K}$, with max. in magn. susceptibility at 100K

- resistivity drop $T = 350\text{K} - 550\text{K}$, “metal-insulator transition”

What happens with the band gap of LaCoO_3 ?

Photoelectron spectroscopy on LaCoO_3

T. Koethe -- PhD thesis



spectrum of **surface** different from spectrum of **bulk** !!

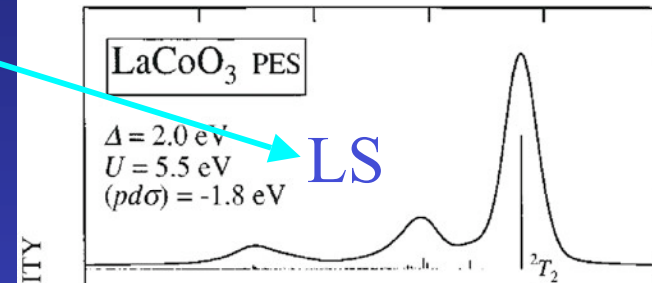
Photoelectron spectroscopy on correlated oxide systems

T. Koethe -- PhD thesis

Spin state transition in LaCoO_3

„bulk“

Cluster model calculations:
Saitoh *et al.*, PRB **55**, 4257 (1997)

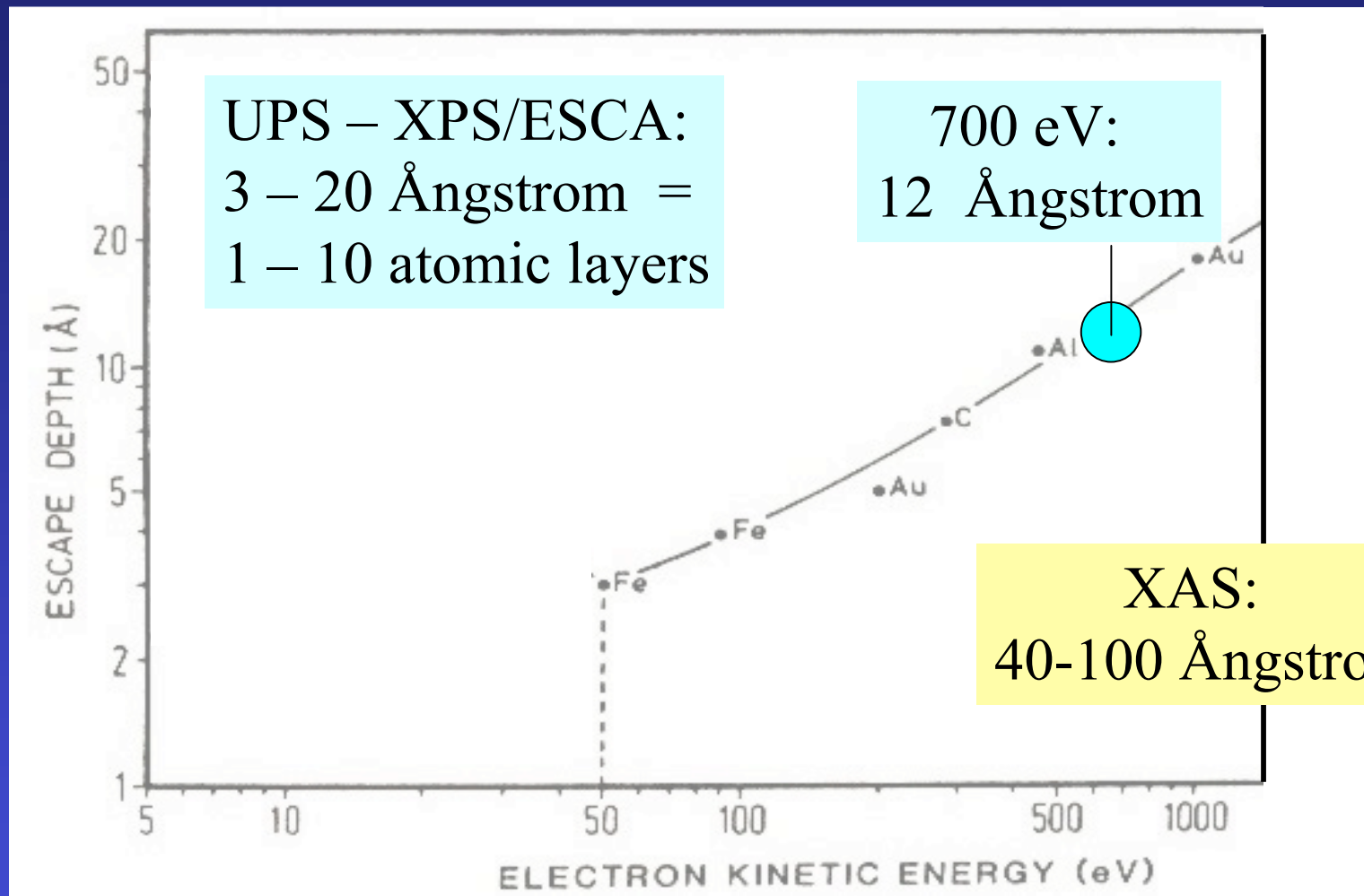


correlated electron systems:
surface tends to have a different electronic structure than **bulk**

this is especially acute for systems close to an instability,
e.g. metal-insulator and spin-state transitions

spin state at the surface
is different from bulk !!

Probing depth of photoelectron spectroscopy



Problematic for obtaining bulk PES electronic structure of LaCoO_3

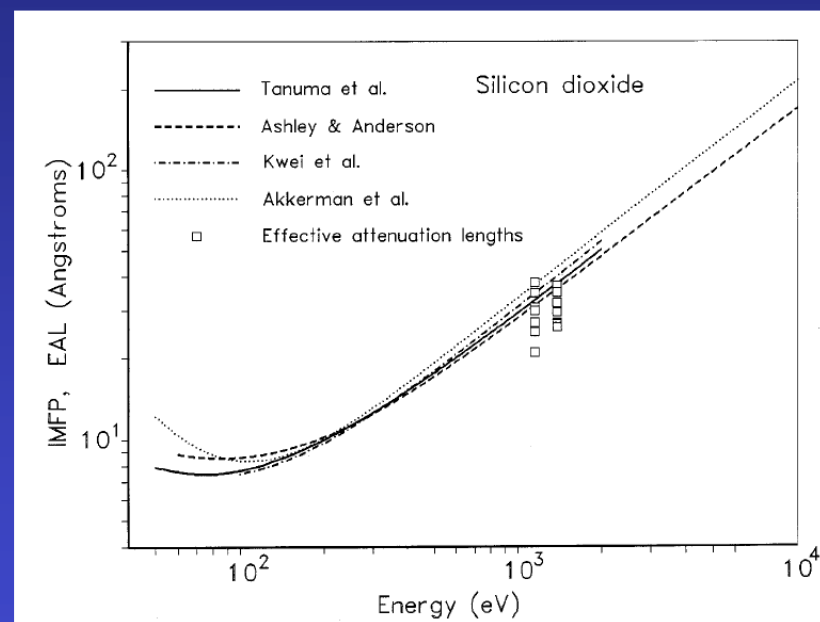
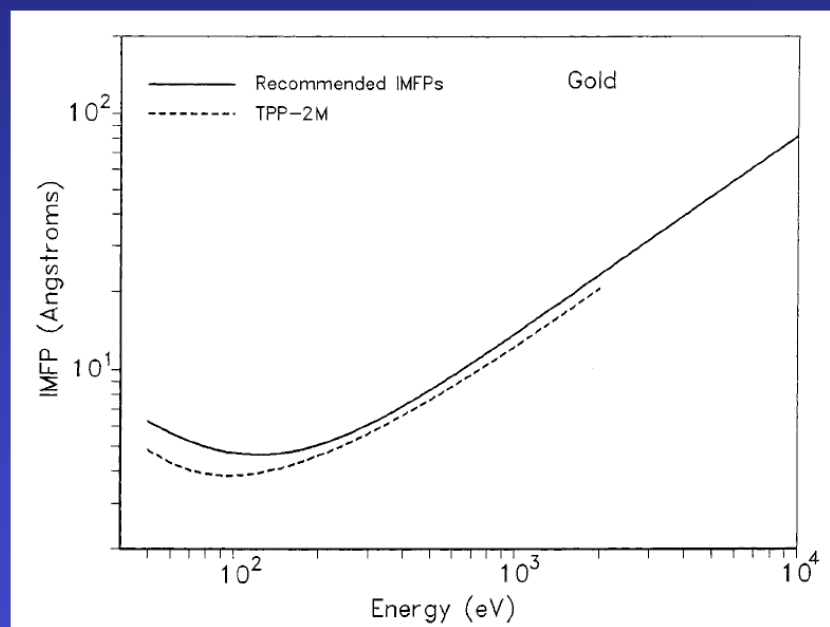
Problematic for high temperature phases: loss of oxygen at surface

How to increase the probing depth of photoemission ?

SURFACE AND INTERFACE ANALYSIS
Surf. Interface Anal. 29, 108–114 (2000)

Evaluation of electron inelastic mean free paths for selected elements and compounds†

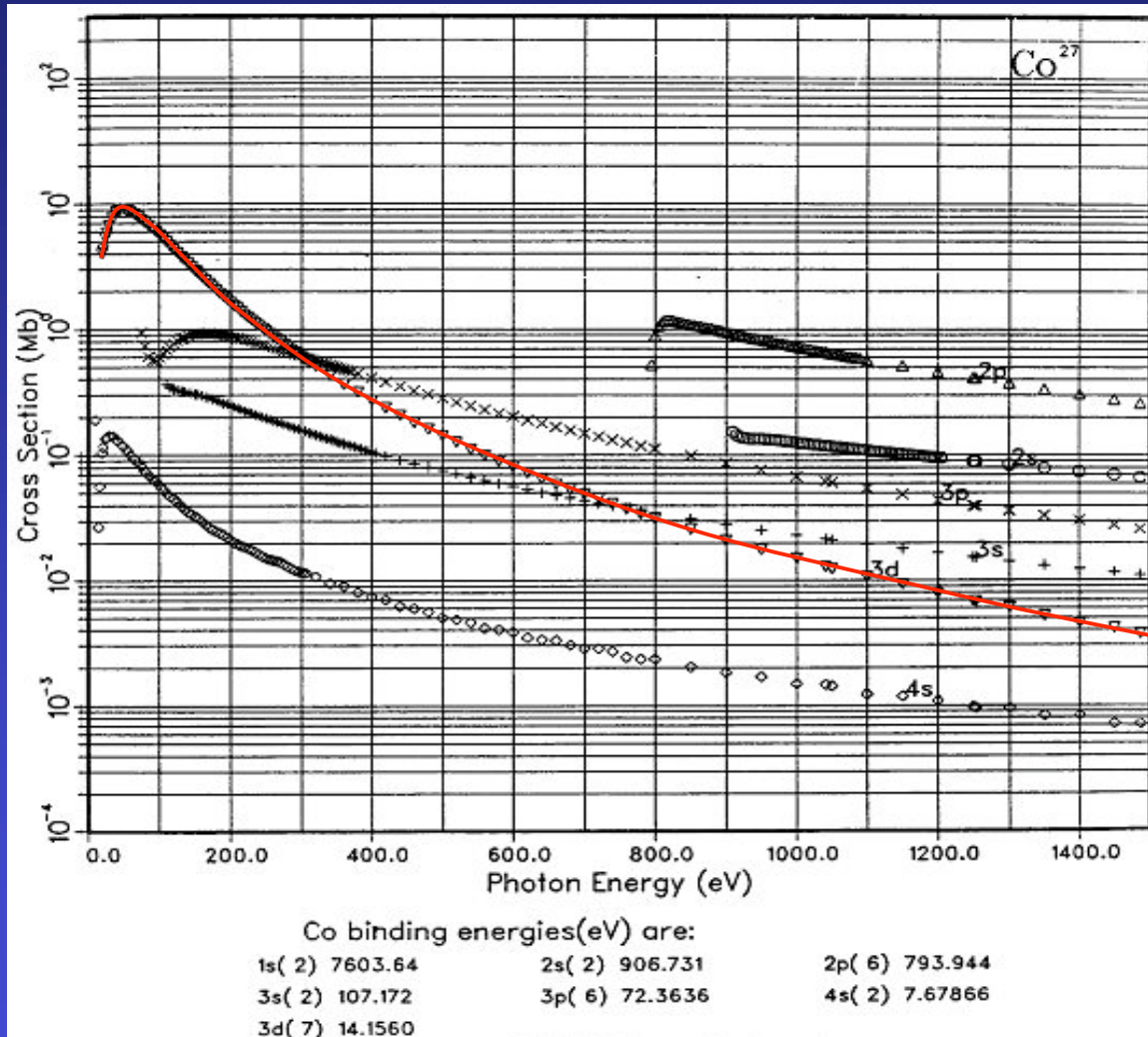
C. J. Powell^{1*} and A. Jablonski²



electron energies 10 keV → probing depths 80-200 Ångstrom !!

→ hard-x-ray photoelectron spectroscopy = HAXPES

Photo-ionization cross-section



- cross-sections go down with increasing $h\nu$
- transition metal 3d
20 eV \rightarrow 10 keV : 10^{-6}
- needs bright sources !!

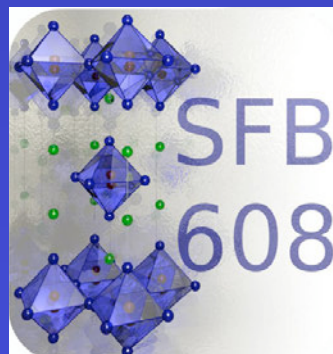
Hard-x ray photoelectron spectroscopy (HAXPES):

MPI-CPfS-Dresden
Univ. Cologne

NSRRC-Taiwan

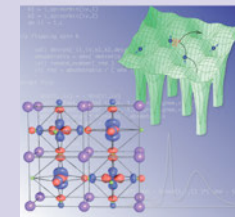


• PhD project Jonas Weinen

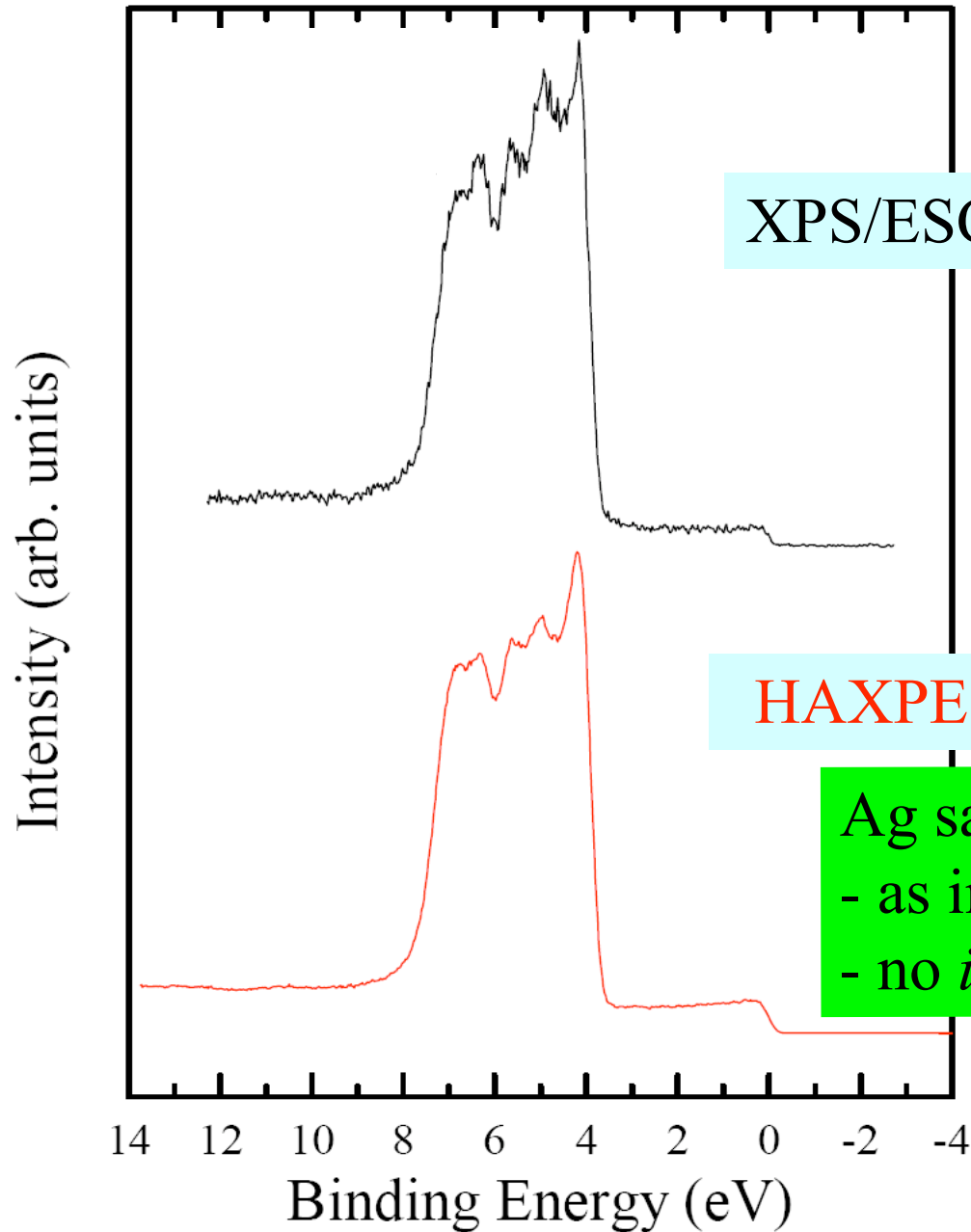


MAX-PLANCK-GESELLSCHAFT

FOR 1346:
Dynamical Mean-Field Approach
with Predictive Power
for Strongly Correlated Materials



Ag valence band



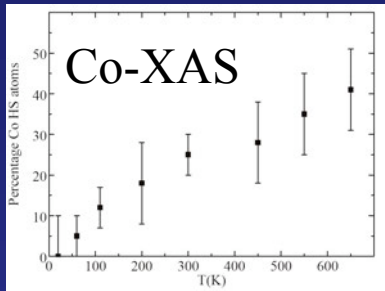
XPS/ESCA: $h\nu = 1486 \text{ eV}$

HAXPES: $h\nu = 7600 \text{ eV}$

Ag sample

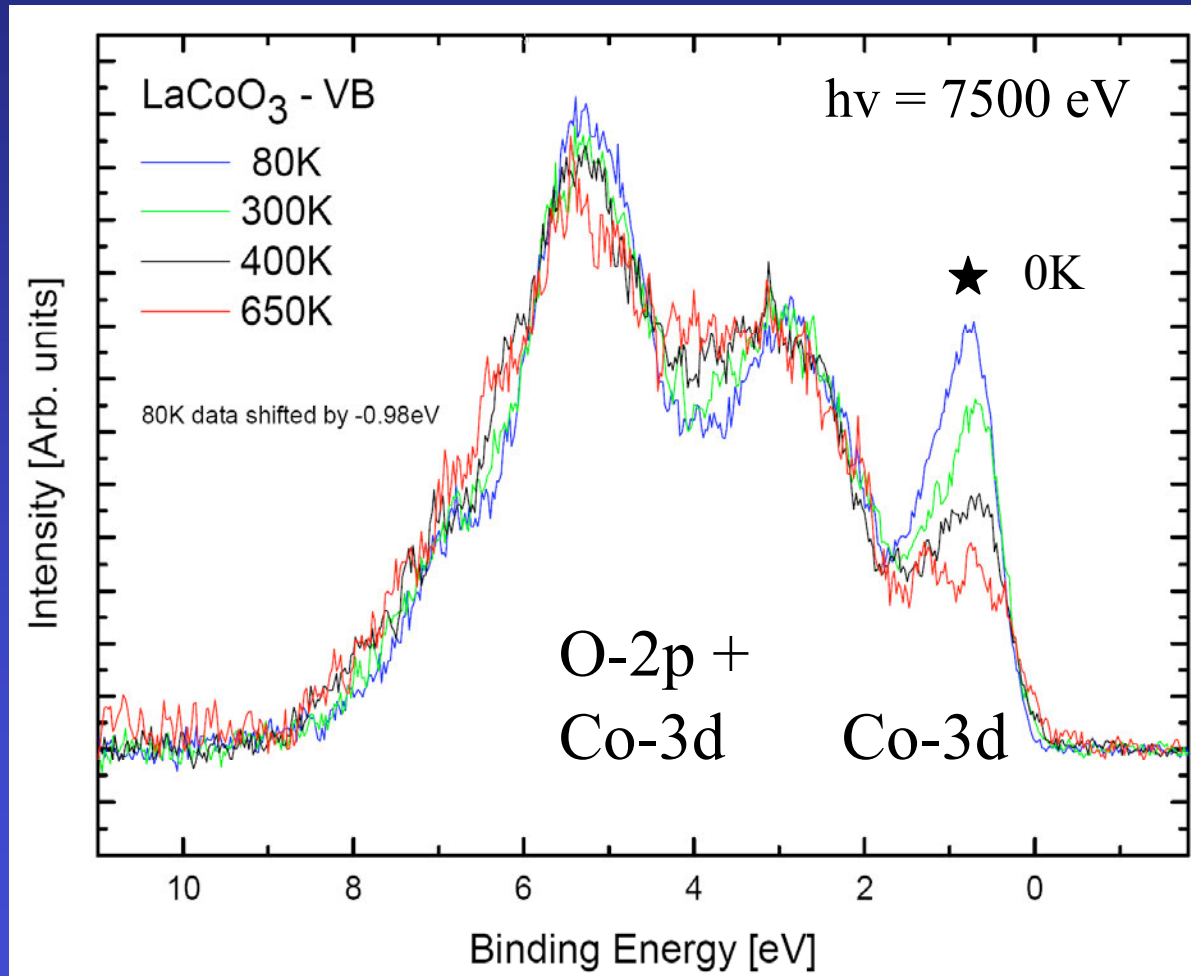
- as introduced from air

- no *in-situ* treatments



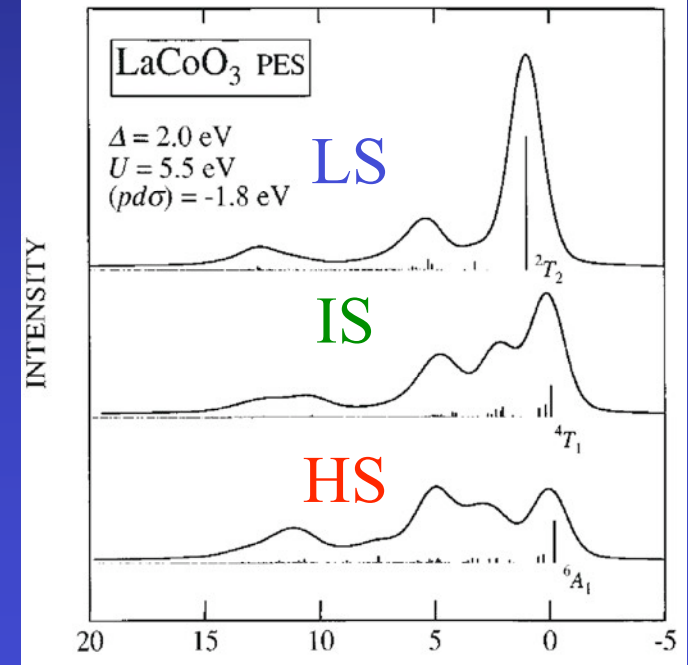
HAXPES: valence band LaCoO_3

Stefano Agrestini, Jonas Weinen

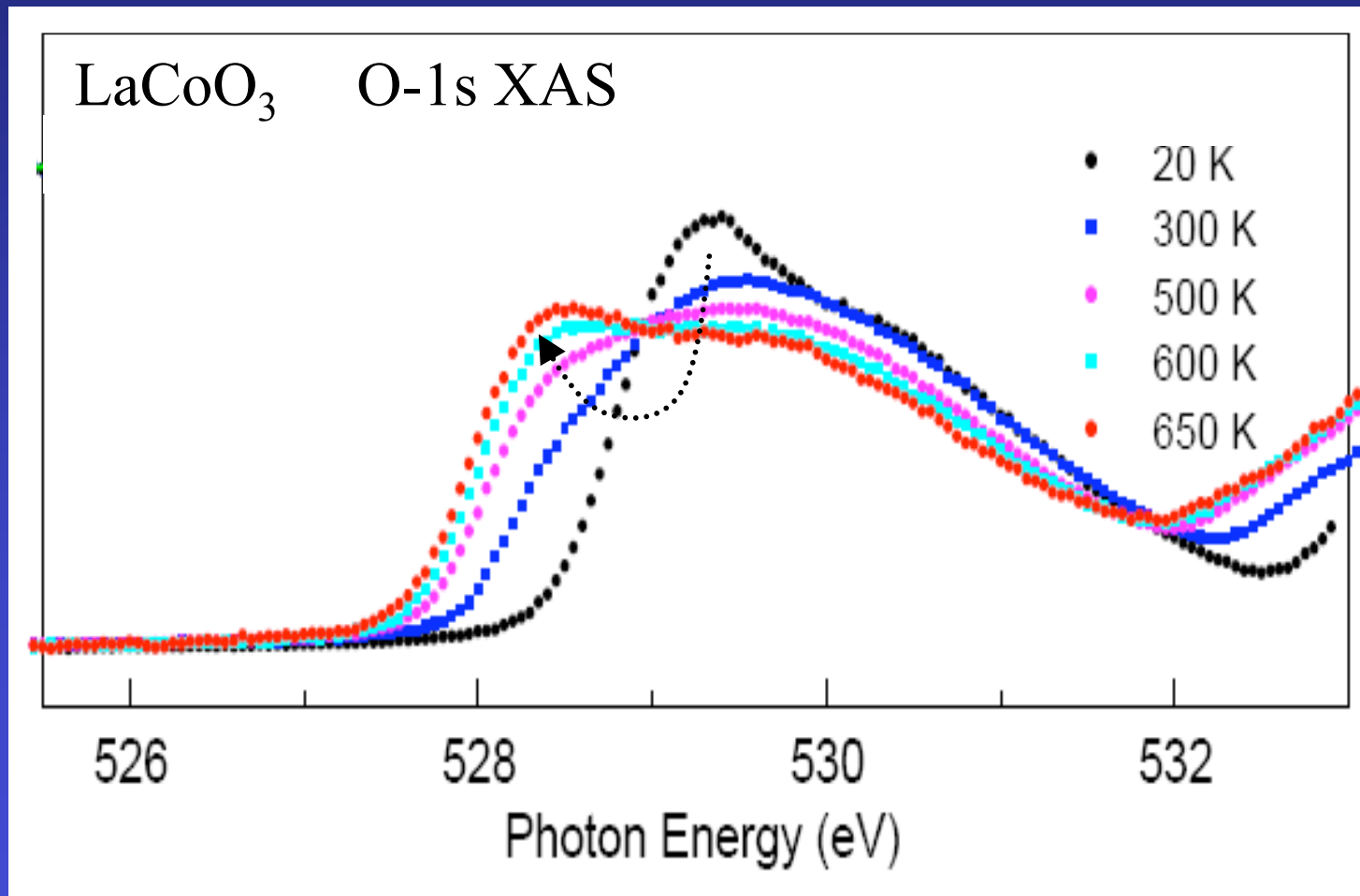
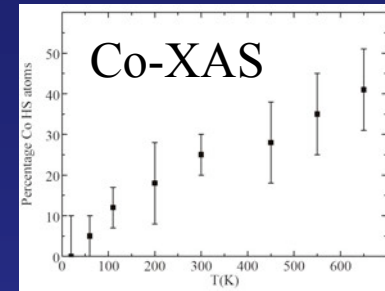


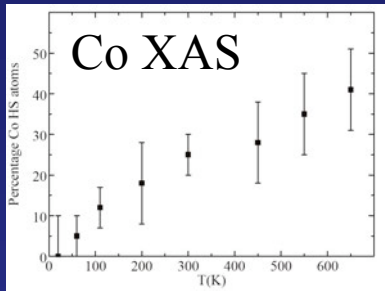
Cluster model calculations:

Saitoh *et al.*, PRB **55**, 4257 (1997)

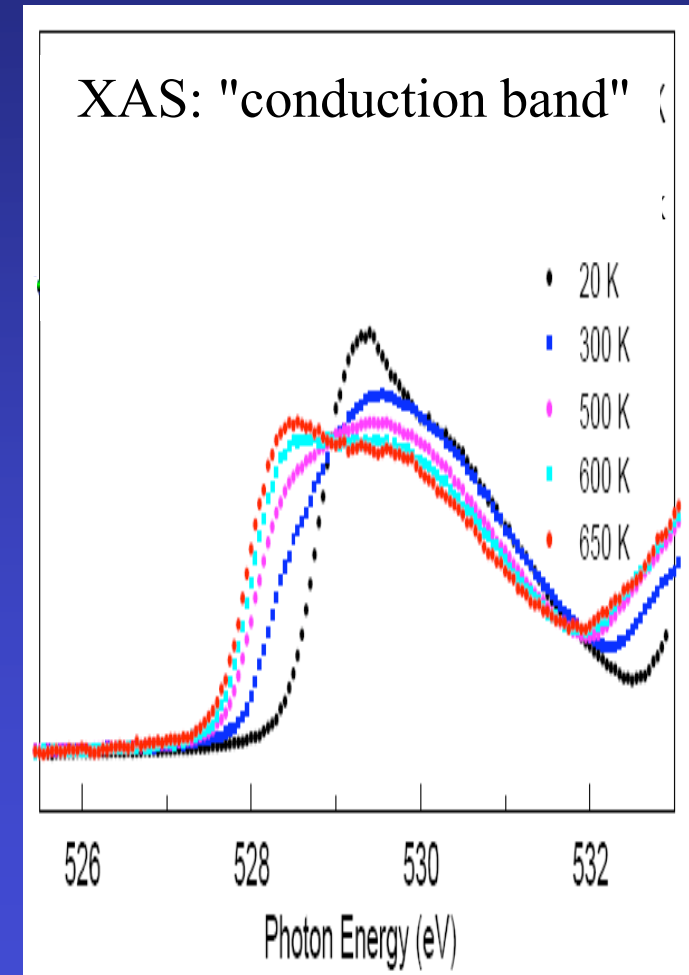
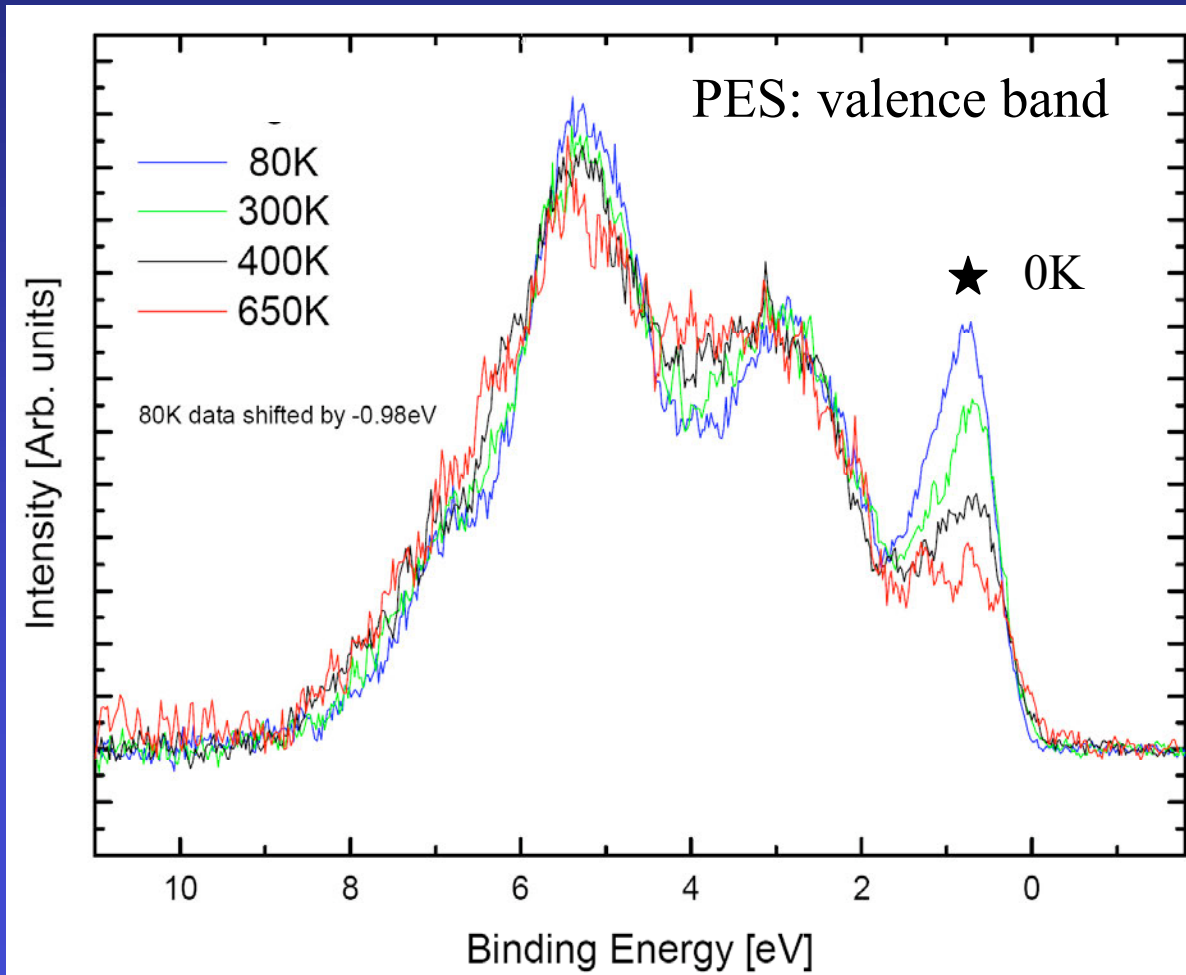


conduction band LaCoO_3 :
O-1s XAS
(IPES would have been better)



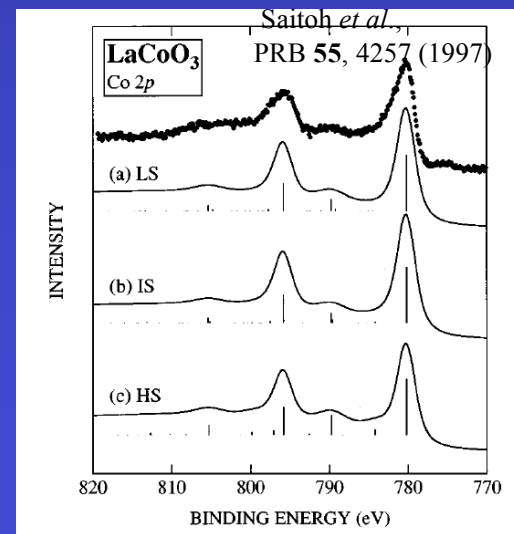
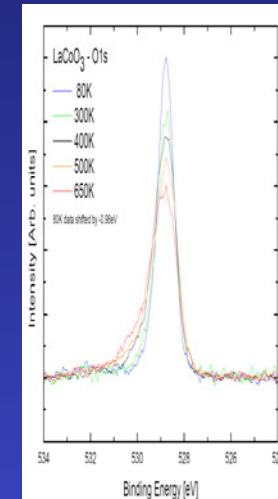
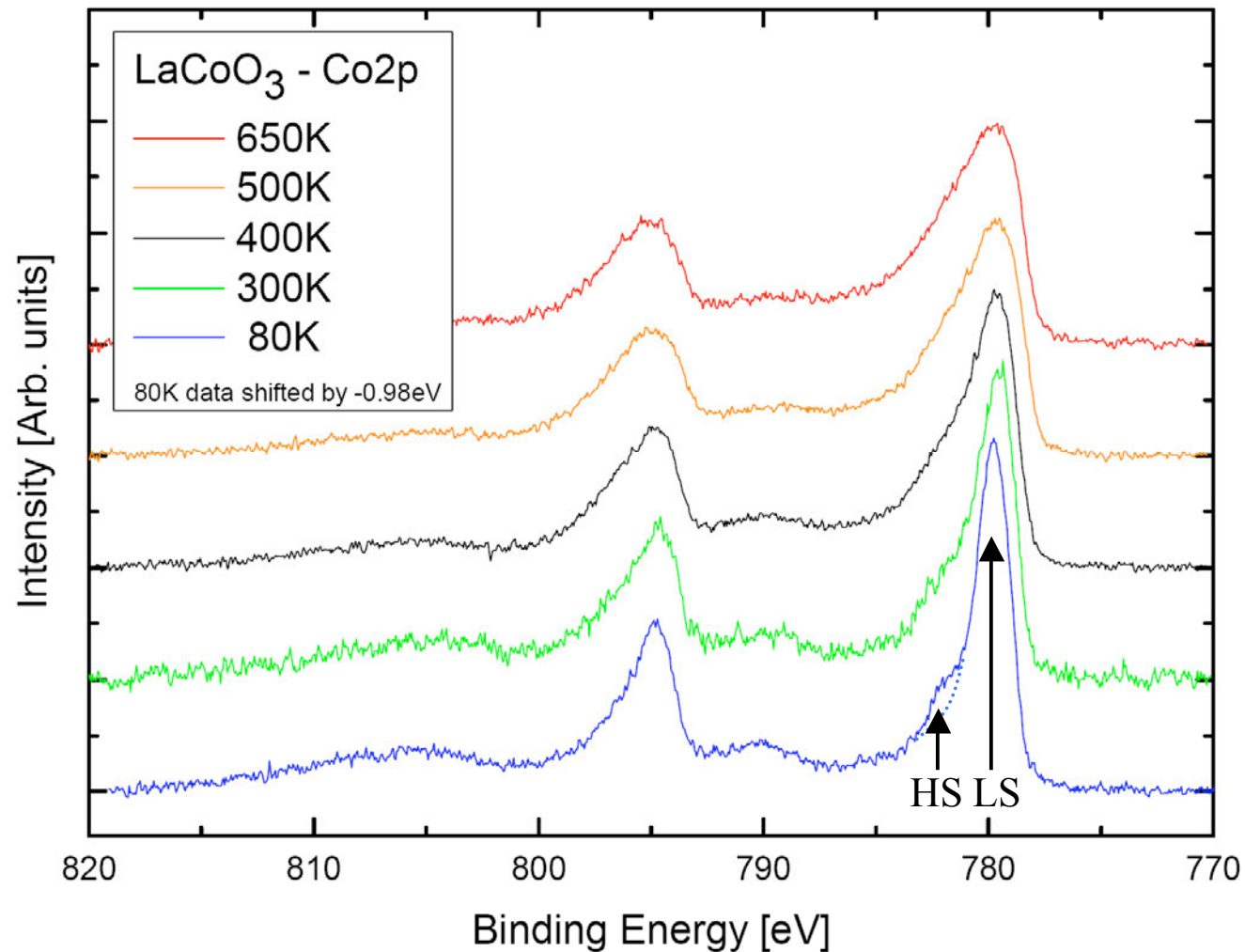
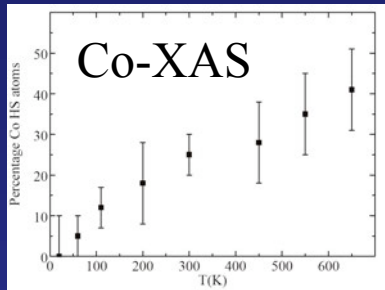


- Band gap reduces by 1 eV from 0 to 650 K
- LaCoO_3 is still a bad metal across the "MIT"



HAXPES: core levels LaCoO_3

- LaCoO_3 is an inhomogeneous spin-state system



- Cobaltates: a material class with very rich physics *metal-insulator-transitions, colossal-magneto-resistance, thermoelectric power, superconductivity, multiferroicity, unusual magnetic properties.*

- valence state : 2+ 3+ 4+
- spin state: HS HS/IS/LS HS/IS/LS
- spin state transitions and spin-blockade mechanism
- local coordination: oct/pyramidal/trig/prismatic-trig

113 cobaltates:

LaCoO₃ Co: 3+ LS/HS
 LaMn_{0.5}Co_{0.5}O₃ Co: 2+ HS Mn: 4+ HS
 EuMn_{0.5}Co_{0.5}O₃ Co: 2+ HS Mn: 4+ HS

one-dimensional cobaltates:

Ca₃Co₂O₃ Co: trig. 3+ HS Co: oct. 3+ LS
 Ca₃CoRhO₆ Co: trig. 2+ HS Rh: oct. 4+ LS
 Ca₃FeRhO₆ Fe : trig. 3+ HS Rh: oct. 3+ LS
 Ca₃CoMnO₆ Co: trig. 2+ HS Mn: oct. 4+ HS

layered cobaltates:

GdBaCo₂O_{5.5} Co: 3+ HS/LS Co_{pyr}: 3+ HS
 La_{1.5}Sr_{0.5}CoO₄ Co: 3+ LS Co: 2+ HS

Na_xCoO₂ Co: 3+ LS Co: 4+ LS
 Na_xCoO₂·nH₂O Co: 3+ LS Co: 4+ LS

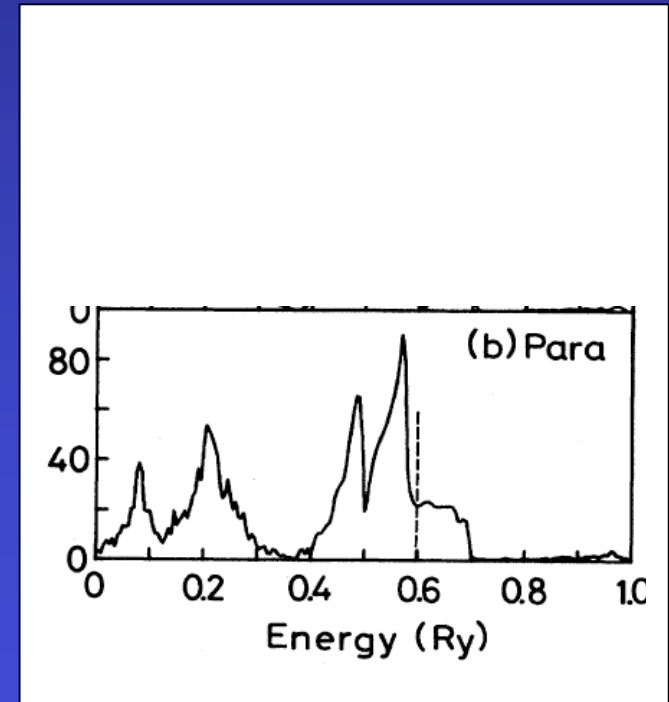
II. Multiplet interactions and band formation

Example : NiO

Benchmark system in solid state physics

- partially filled 3d-shell
- antiferromagnetic insulator

Standard band theory: NiO metallic



Example : NiO

Benchmark system in solid state physics

- partially filled 3d-shell
- antiferromagnetic insulator

Standard band theory: NiO metallic

Early attempt to fix: Slater insulator

PHYSICAL REVIEW B VOLUME 28, NUMBER 11 1 DECEMBER 1983
Band theory of the magnetic interaction in MnO, MnS, and NiO

VOLUME 52, NUMBER 20 PHYSICAL REVIEW LETTERS 14 MAY 1984
Transition-Metal Monoxides: Band or Mott Insulators

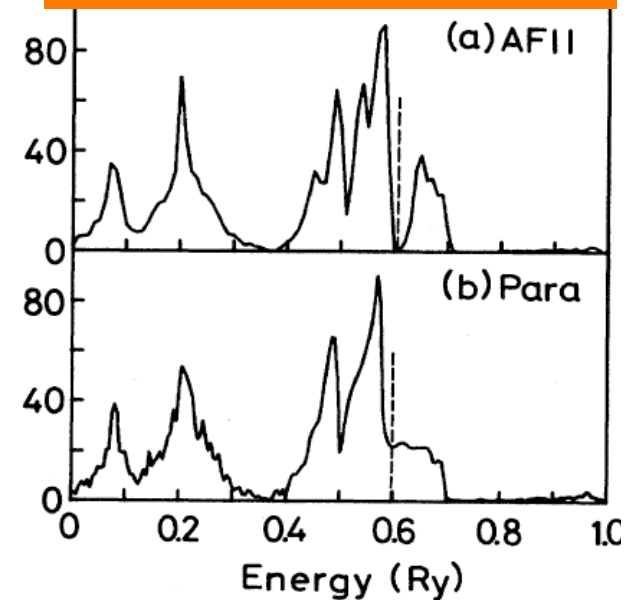
PHYSICAL REVIEW B VOLUME 30, NUMBER 8 15 OCTOBER 1984
Band theory of insulating transition-metal monoxides: Band-structure calculations

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Institute for Solid State Physics, University of Tokyo, Roppongi, Minato-ku, Tokyo 106, Japan

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IBM Thomas J. Watson Research Center, P.O. Box 218, Yorktown Heights, New York 10598

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(Received 11 April 1984)

Gap is very small: 0.2 eV



Example : NiO

Benchmark system in solid state physics

- partially filled 3d-shell
- antiferromagnetic insulator

PHYSICAL REVIEW B VOLUME 30, NUMBER 2 15 JULY 1984

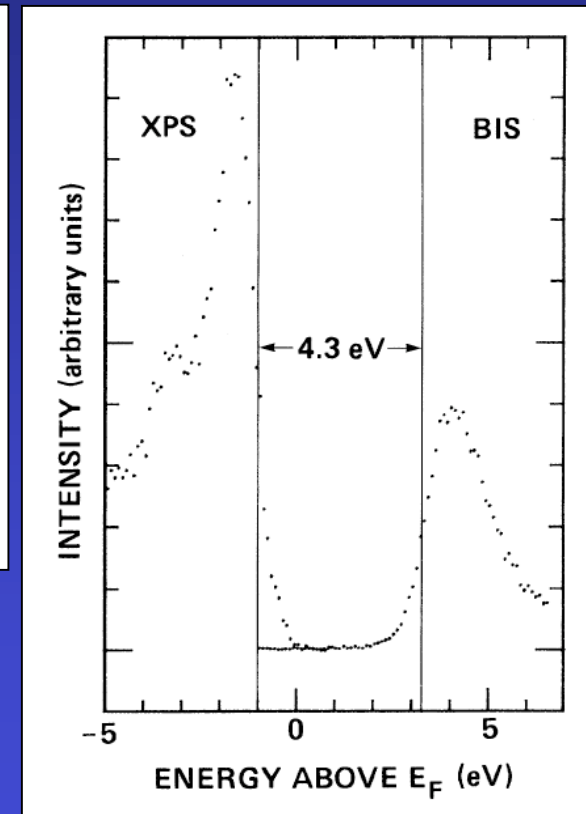
Valence-band photoemission and optical absorption in nickel compounds

Atsushi Fujimori* and Fujio Minami
National Institute for Research in Inorganic Materials, Sakura-mura, Niihari-gun, Ibaraki 305, Japan
(Received 21 February 1984)

VOLUME 53, NUMBER 24 PHYSICAL REVIEW LETTERS 10 DECEMBER 1984

Magnitude and Origin of the Band Gap in NiO

G. A. Sawatzky^(a) and J. W. Allen
Xerox Palo Alto Research Center, Palo Alto, California 94304
(Received 5 July 1984)



PES- IPES experiment: gap is large → 4.3 eV

importance of Hubbard U

Example : NiO

Benchmark system in solid state physics

- partially filled 3d-shell
- antiferromagnetic insulator

VOLUME 65, NUMBER 9

PHYSICAL REVIEW LETTERS

27 AUGUST 1990

Transition-Metal Oxides in the Self-Interaction-Corrected Density-Functional Formalism

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(Received 14 May 1990)

TABLE I. The LSD, SIC-LSD, and experimental energy band gaps (in eV) and spin magnetic moments (in μ_B) for VO, CrO, MnO, FeO, CoO, NiO, and CuO. The moments in parenthesis include the orbital contribution. The last column shows the energy difference ΔE (in eV) between the LSD and the SIC-LSD calculation.

Compound	Band gap			Magnetic moment			ΔE
	LSD	SIC-LSD	Expt.	LSD	SIC-LSD	Expt.	
VO	0.0	0.0	0.0	0.0	0.0	0.0	
CrO	0.0	1.01		2.99	3.49 (3.44)		1.0
MnO	0.8	3.98	3.6–3.8 ^a	4.39	4.49 (4.49)	4.79, ^b 4.58 ^c	3.9
FeO	0.0	3.07		3.42	3.54 (4.55)	3.32 ^d	4.6
CoO	0.0	2.81	2.4 ^e	2.33	2.53 (3.72)	3.35, ^f 3.8, ^d 3.8 ^g	6.6
NiO	0.2	2.54	4.3, ^h 4.0 ⁱ	1.04	1.53 (1.80)	1.77, ^b 1.64, ^j 1.90 ^c	9.5
CuO	0.0	1.43	1.37 ^k	0	0.65 (0.78)	0.65 ^l	14.7

Ab-initio:
SIC fixes the gap

Example : NiO

Benchmark system in solid state physics

- partially filled 3d-shell
- antiferromagnetic insulator

PHYSICAL REVIEW B

VOLUME 44, NUMBER 3

15 JULY 1991-I

Band theory and Mott insulators: Hubbard U instead of Stoner I

Vladimir I. Anisimov,* Jan Zaanen,[†] and Ole K. Andersen

PHYSICAL REVIEW B

VOLUME 48, NUMBER 23

15 DECEMBER 1993-I

Density-functional theory and NiO photoemission spectra

V. I. Anisimov, I. V. Solovyev, and M. A. Korotin
M. T. Czyżyk and G. A. Sawatzky

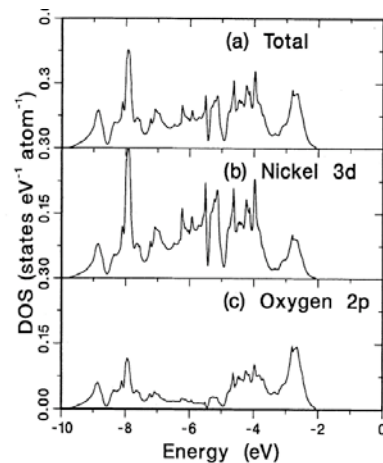


FIG. 1. The density of states (DOS) for the NiO valence band in the LDA+ U calculation ("unmodified"). (a) The total DOS, (b) the Ni 3d partial DOS, and (c) the O 2p partial DOS.

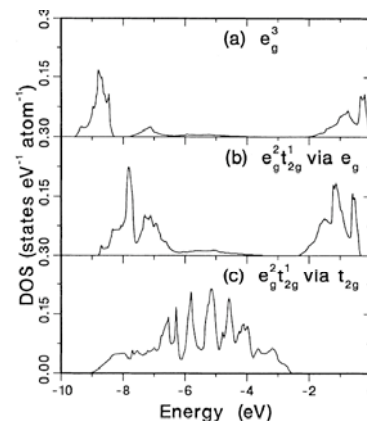


FIG. 2. The partial Ni 3d DOS from the LDA+ U calculation with "modified" hopping parameters for (a) configuration e_g^3 ; (b) the configuration $e_g^2t_{2g}^1$, hybridizing the e_g channel; and (c) the configuration $e_g^2t_{2g}^1$, hybridizing the t_{2g} channel.

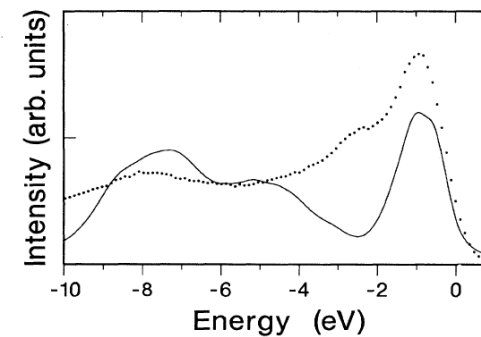


FIG. 3. The experimental (dots) and the calculated (solid line) photoemission spectra for NiO.

LDA+ U not "dynamic" enough for excitation spectra

Aspects of the Correlation Effects, Antiferromagnetic Order, and Translational Symmetry of the Electronic Structure of NiO and CoO

Z.-X. Shen,⁽¹⁾ C. K. Shih,⁽²⁾ O. Jepsen,⁽³⁾ W. E. Spicer,⁽¹⁾ I. Lindau,⁽¹⁾ and J. W. Allen⁽⁴⁾

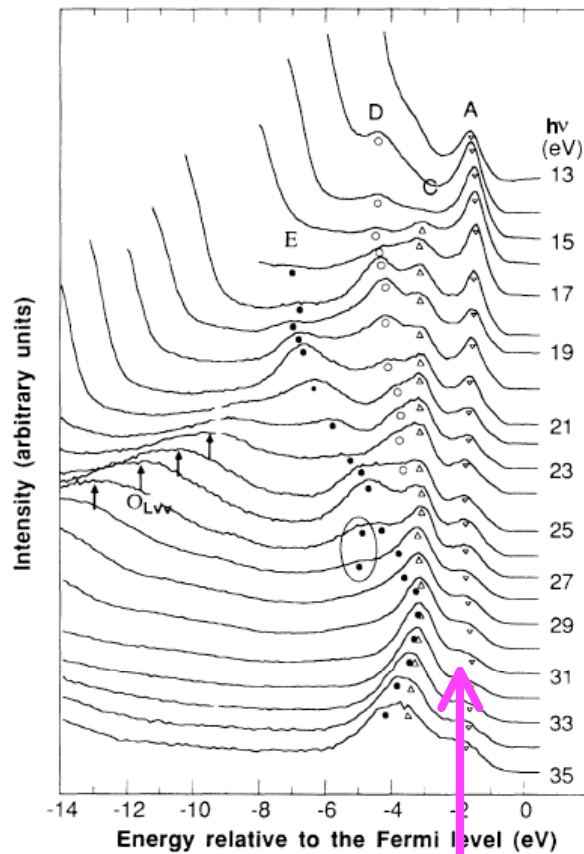
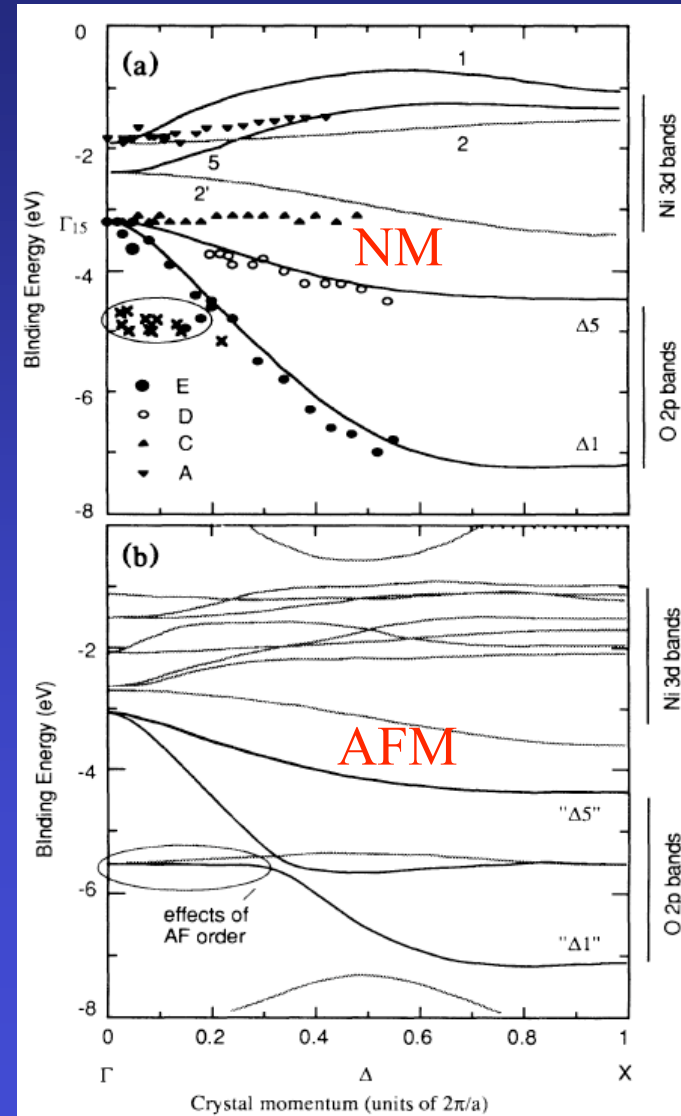
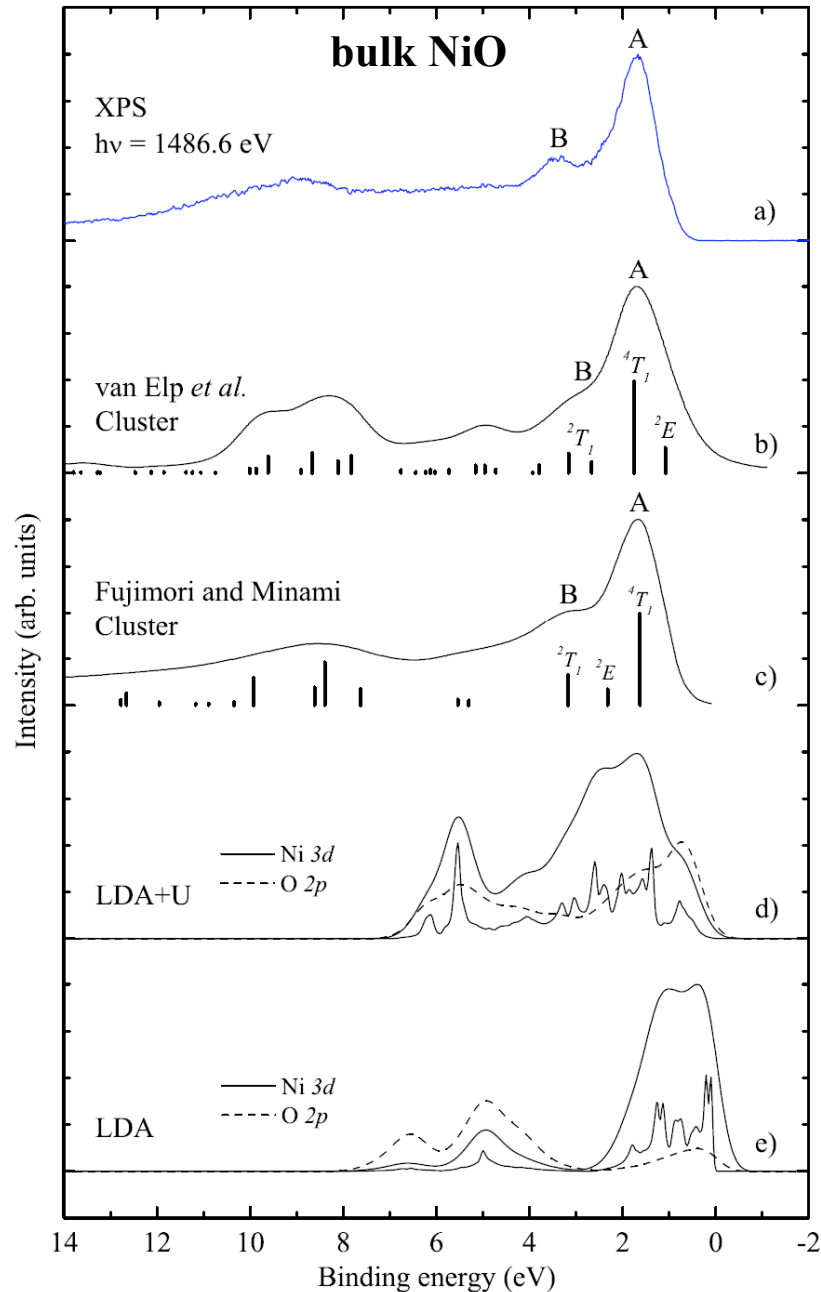


FIG. 1. Normal-emission data from the NiO [001] surface recorded at photon energies between 13 and 35 eV.

hardly any dispersion in the 3d band



Haupricht, Weinen *et al.*



peak B due to surface effects ??

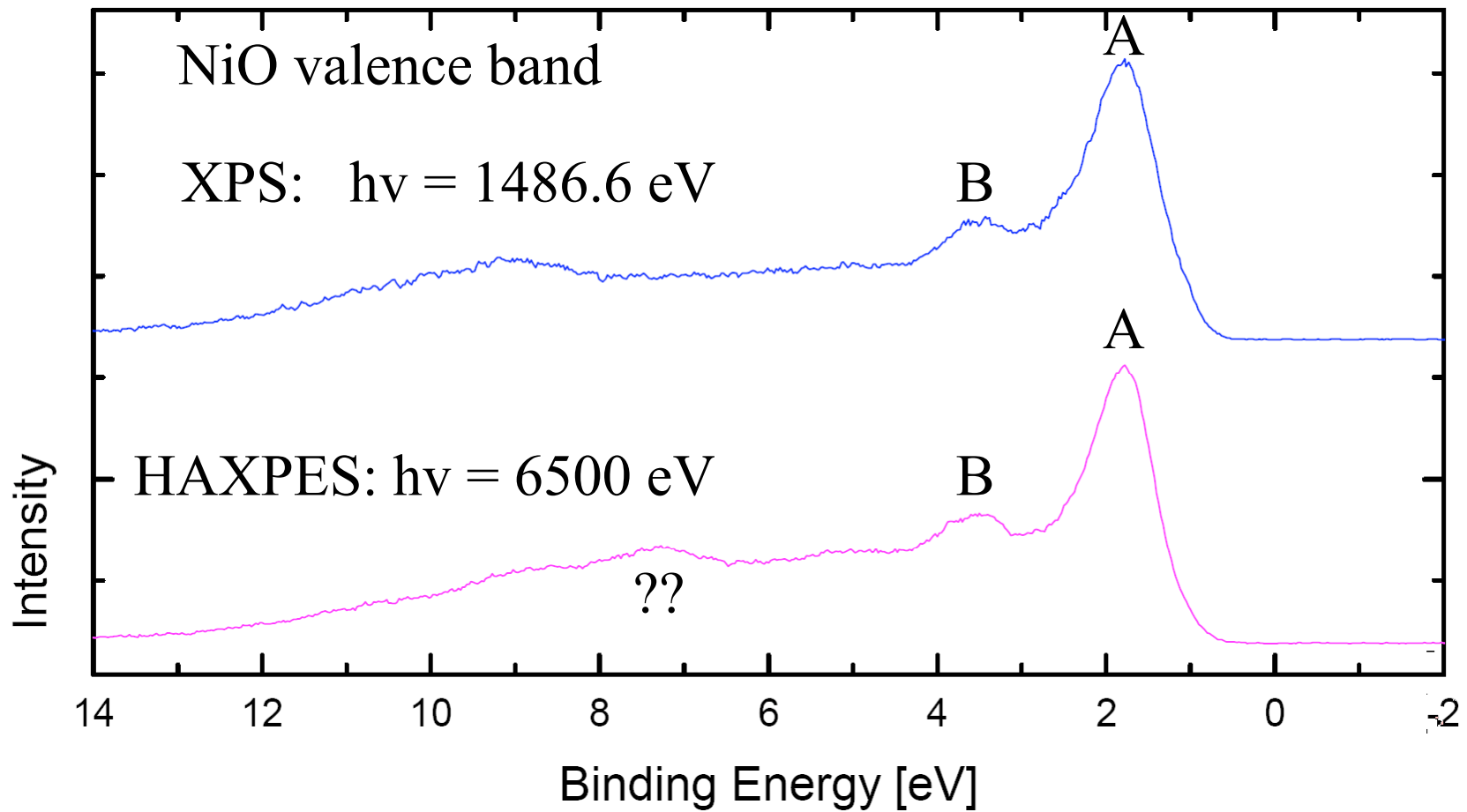
experimental spectrum

cluster: even smaller peak B

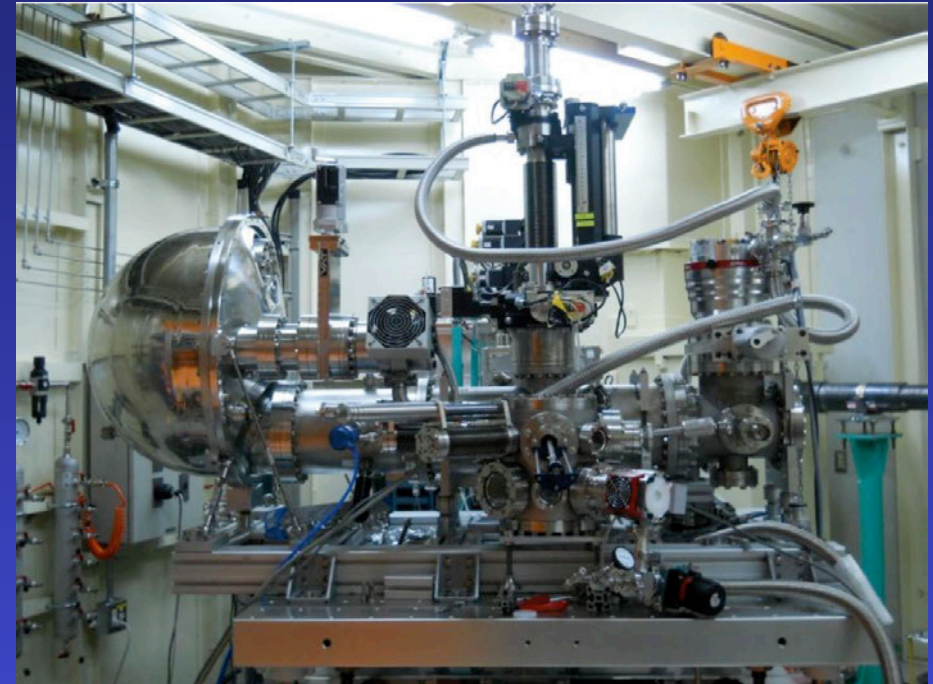
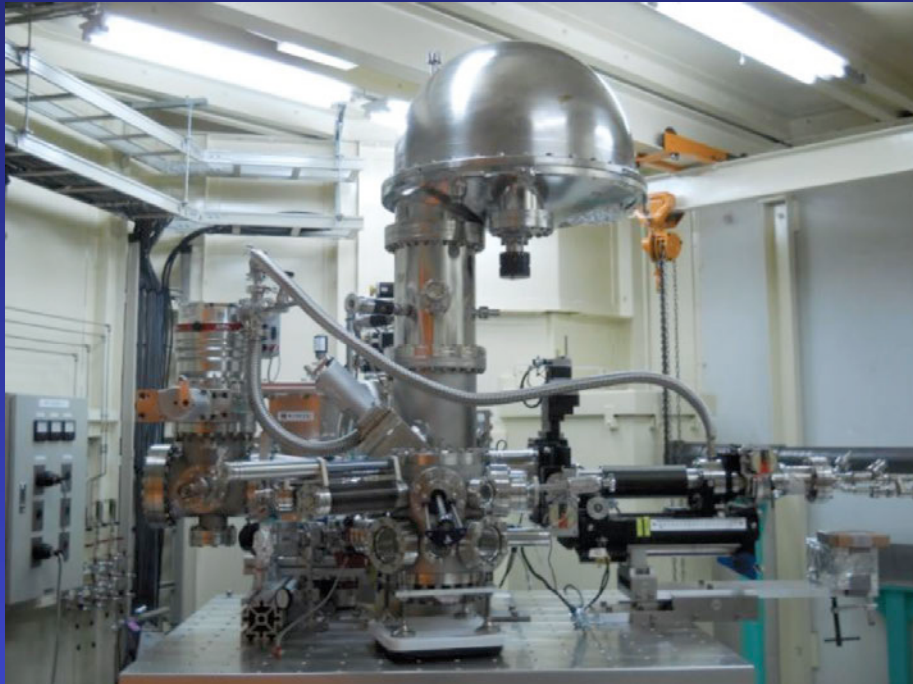
cluster: small peak B

LDA+U: wrong lineshape

LDA: metal



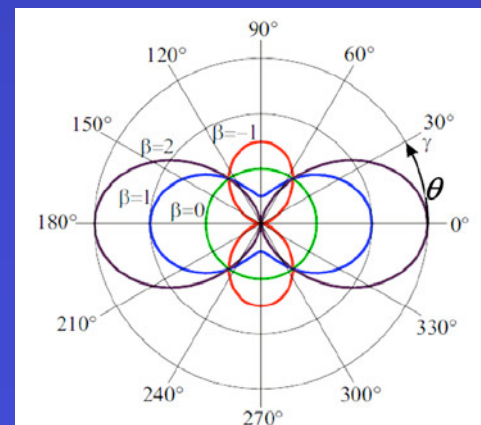
"vertical" set-up $\xleftarrow[\text{polarization}]{\text{light}}$ „horizontal" set-up

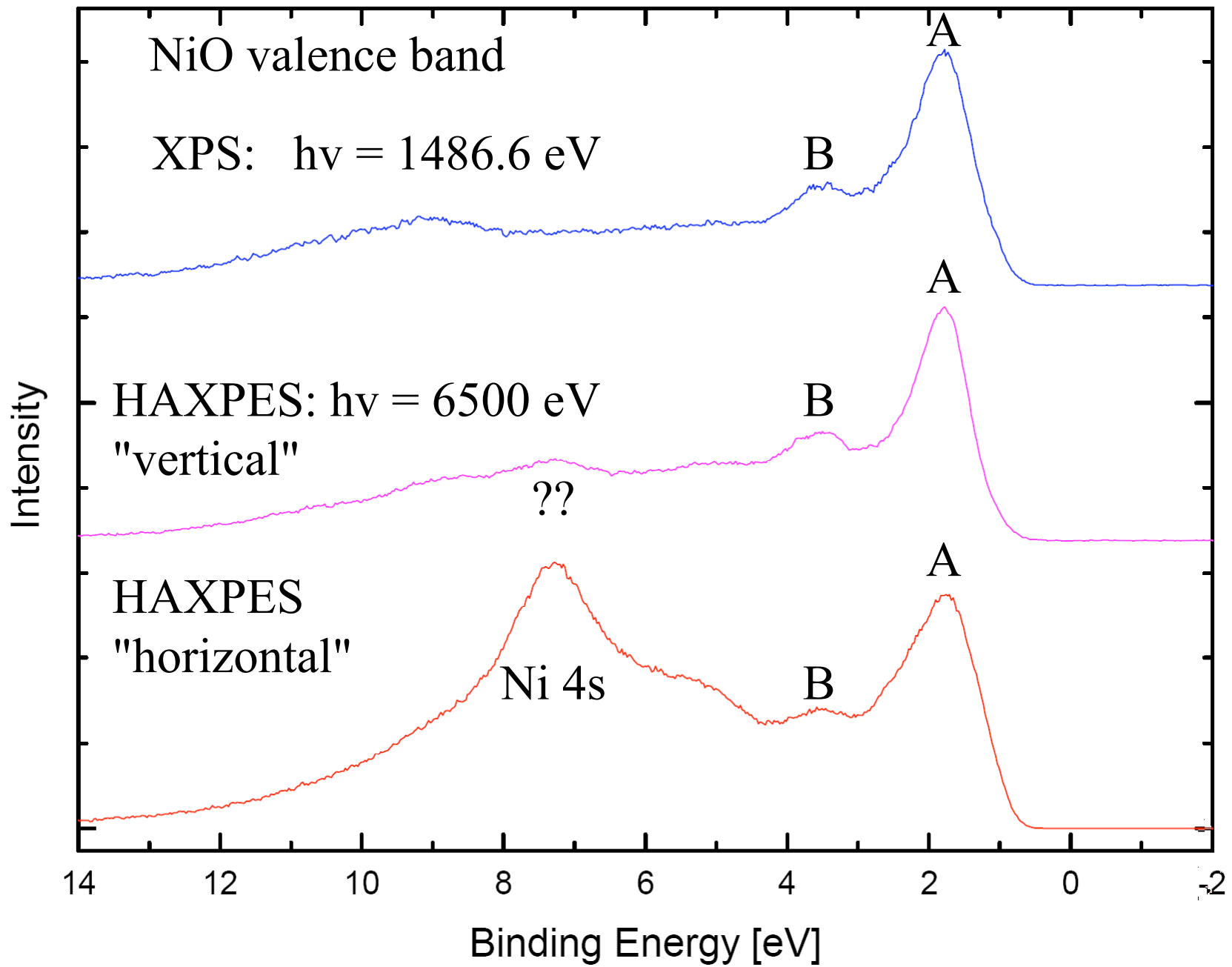


$$\frac{d\sigma_i}{d\Omega} = \frac{\sigma_i}{4\pi} [1 + \beta P_2(\cos\theta) + \dots]$$

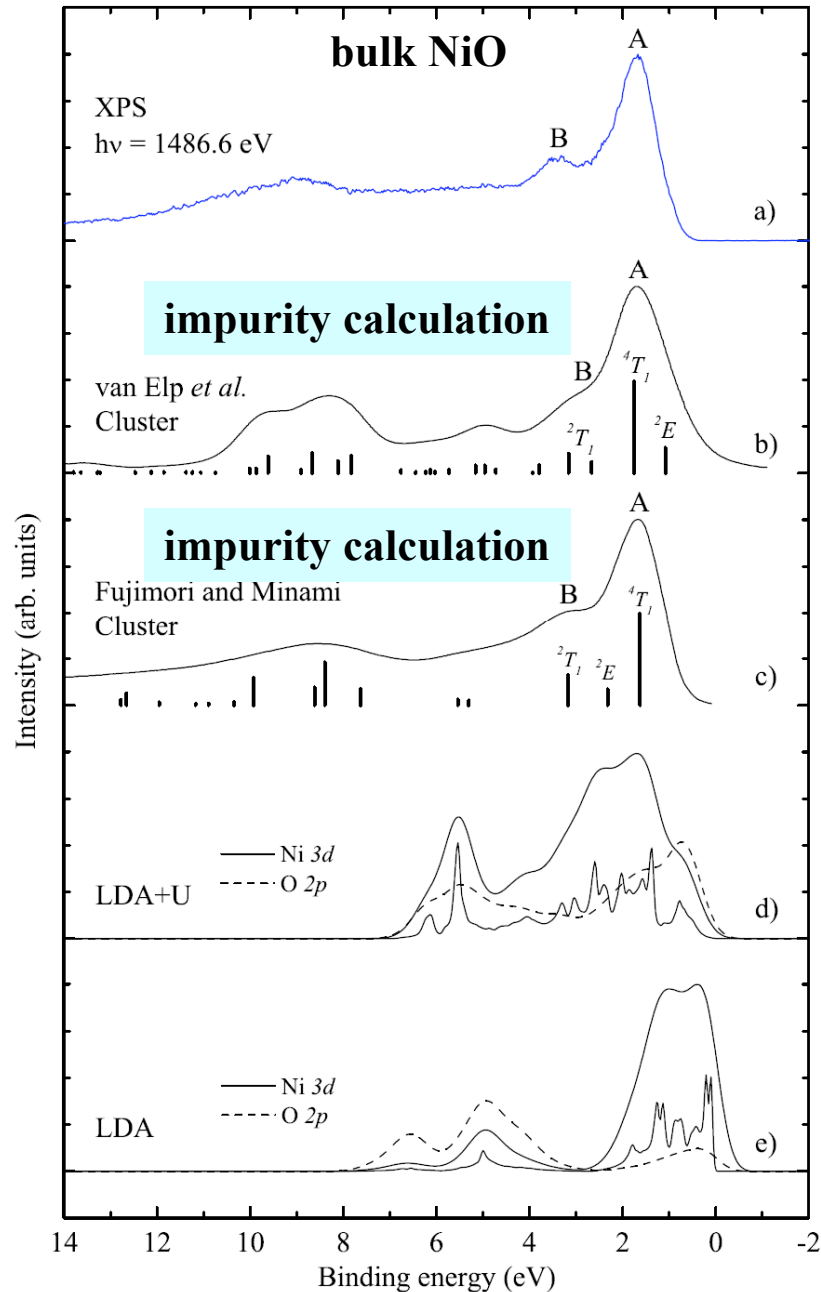
β -parameters @ $h\nu = 5-10$ keV

Cu 3d	0.48 - 0.32
Cu 4s	1.985
Zn 3d	0.50 - 0.33
Zn 4s	1.987 - 1.986

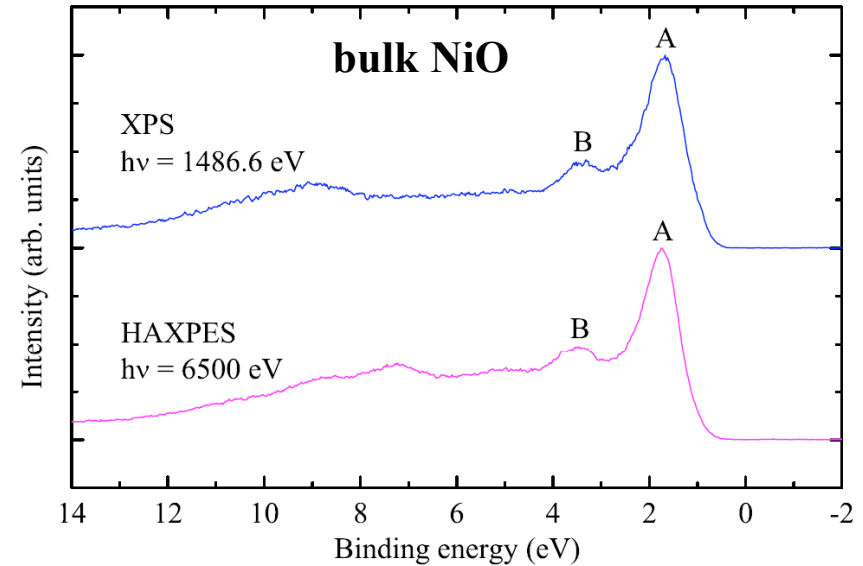




Hauptrecht, Weinen *et al.*



Hauptrecht, Weinen *et al.*



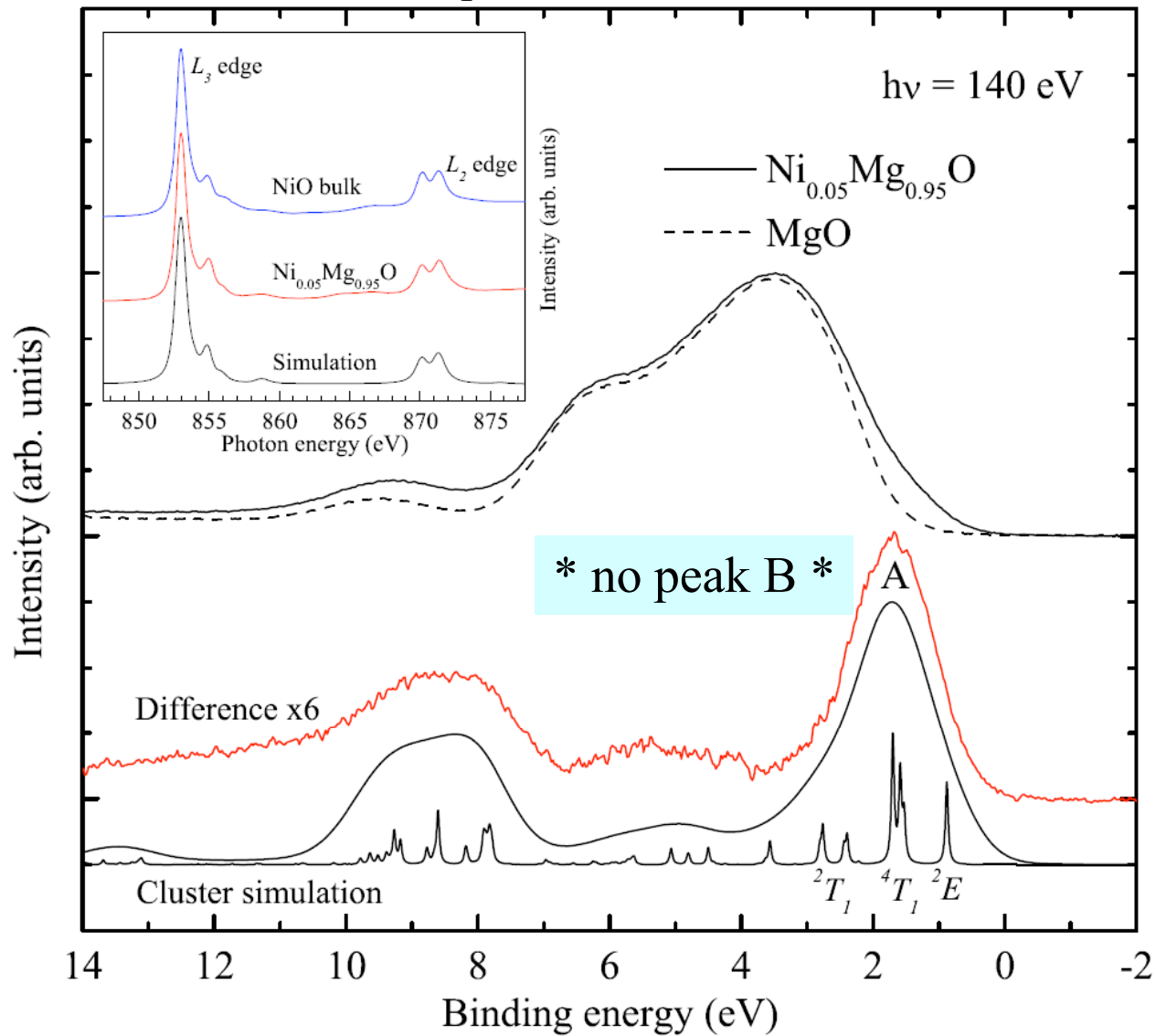
HAXPES is more bulk sensitive than XPS
80 Angstrom - probing depth - 20 Angstrom

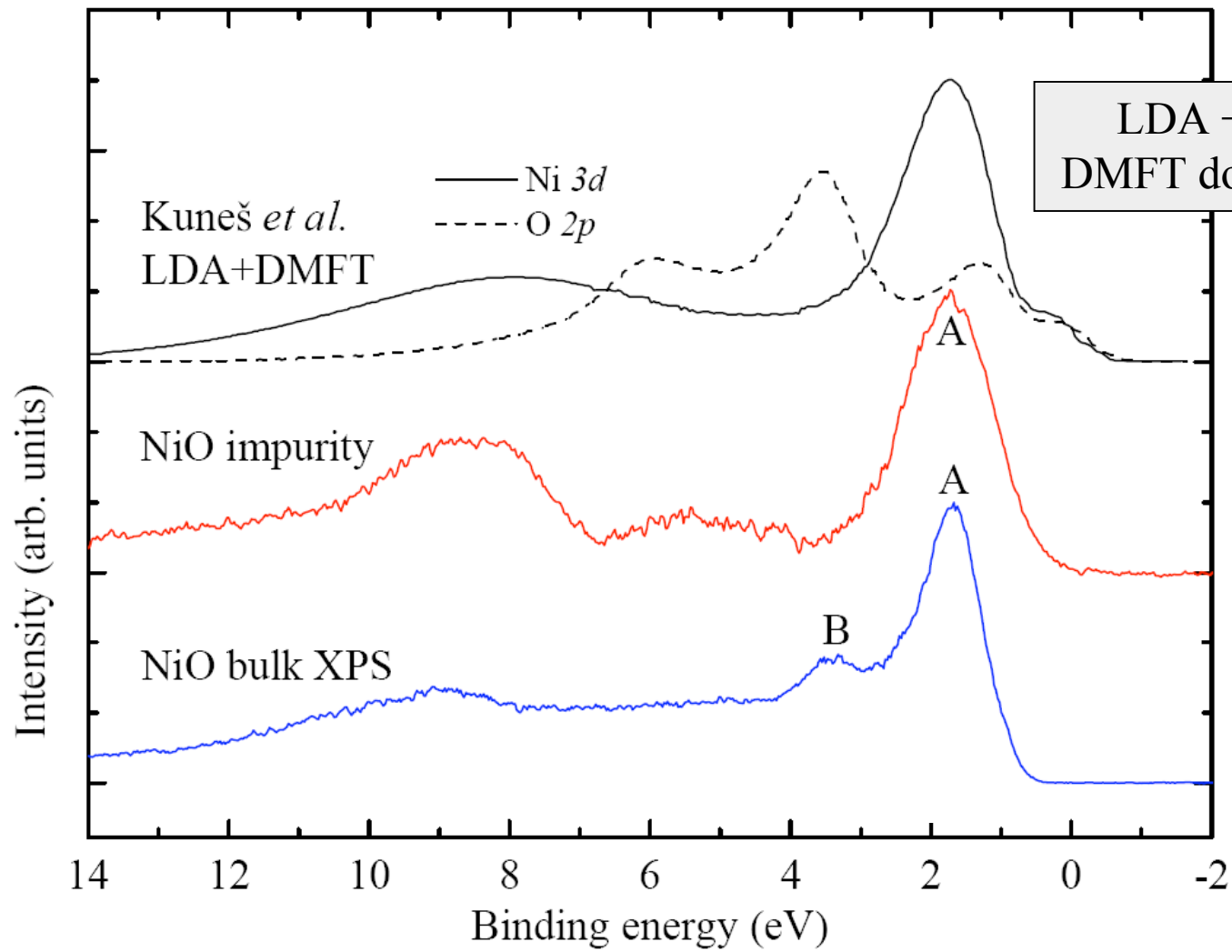
peak B belongs to
bulk spectrum !!

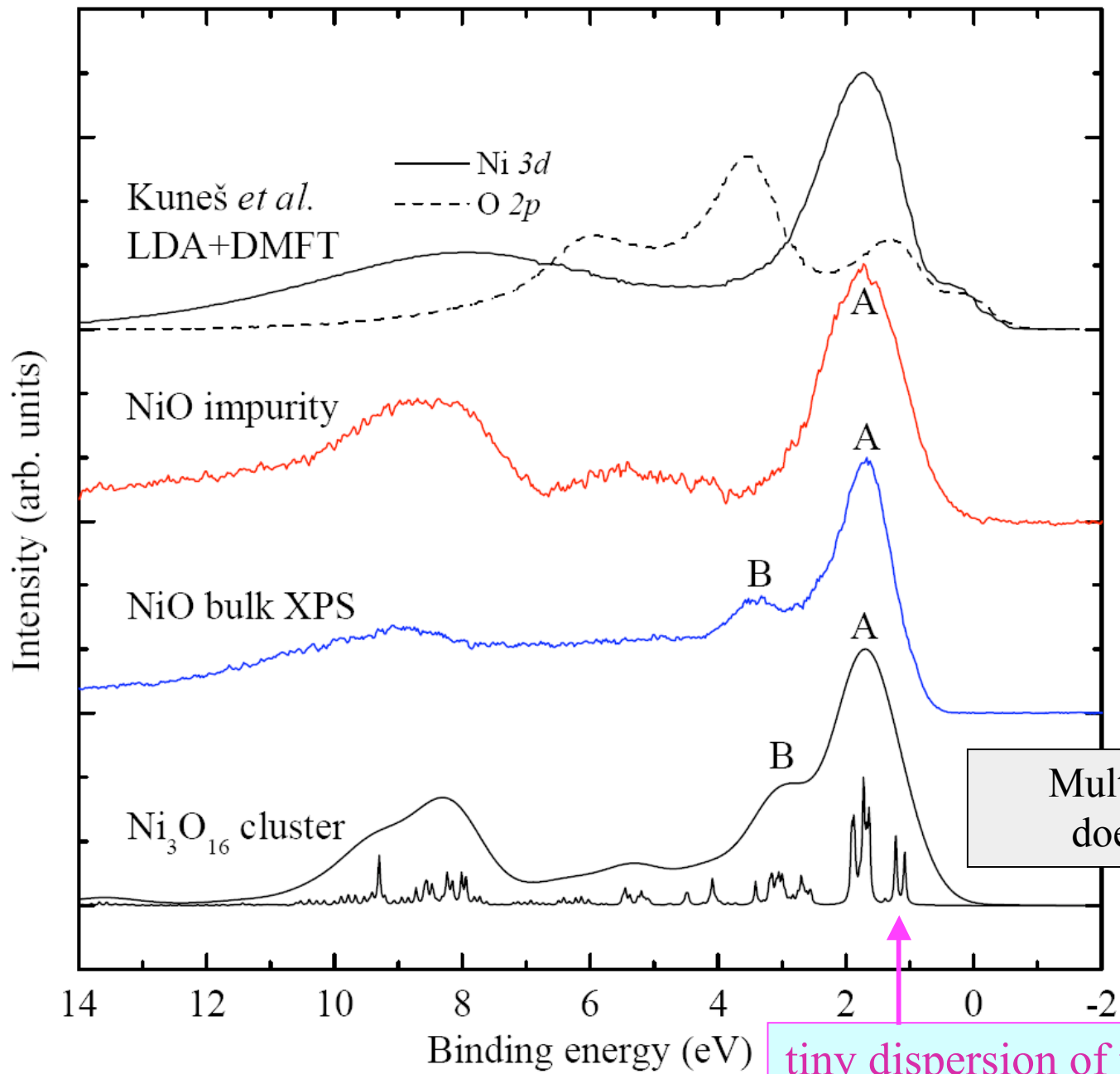
How does the NiO impurity spectrum
look like ?? → Ni:MgO

NiO impurity spectrum → Ni:MgO

Haupricht, Weinen *et al.*

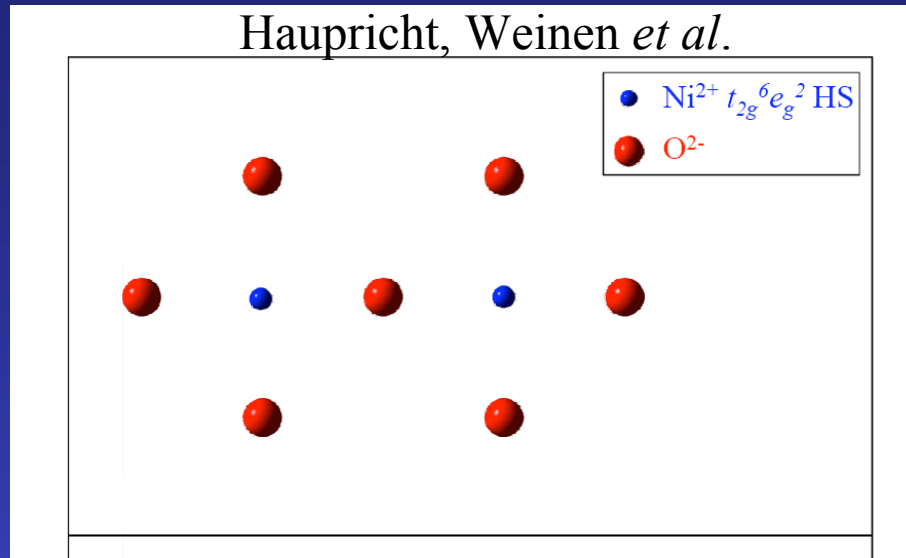






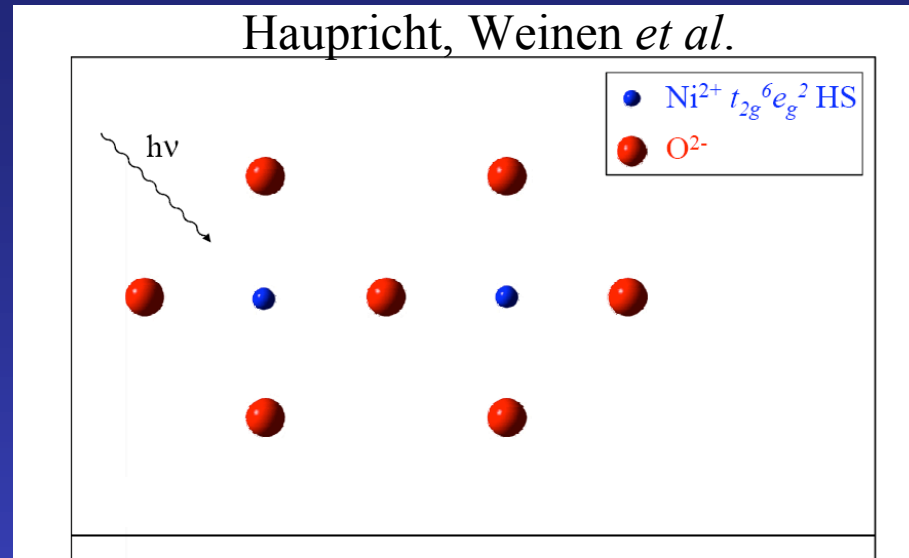
What is the origin of peak B ??

before PES



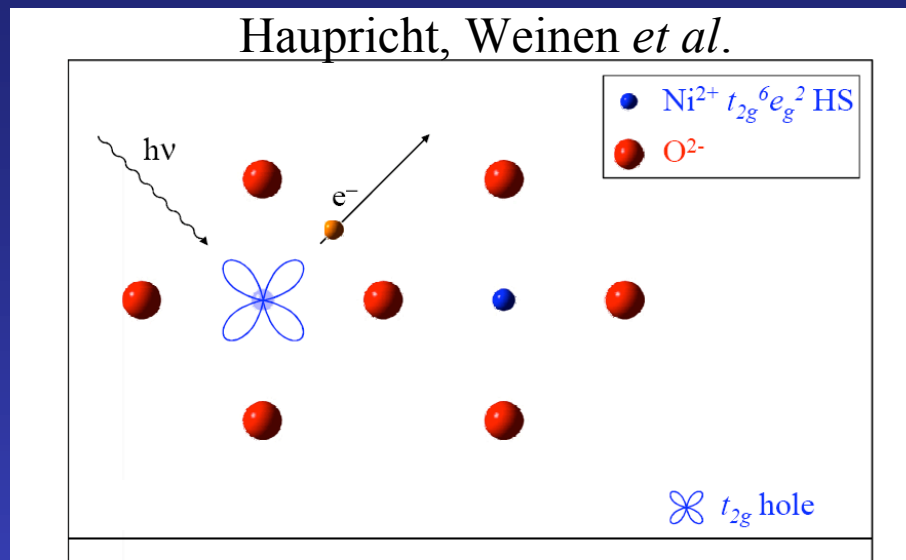
What is the origin of peak B ??

PES



What is the origin of peak B ??

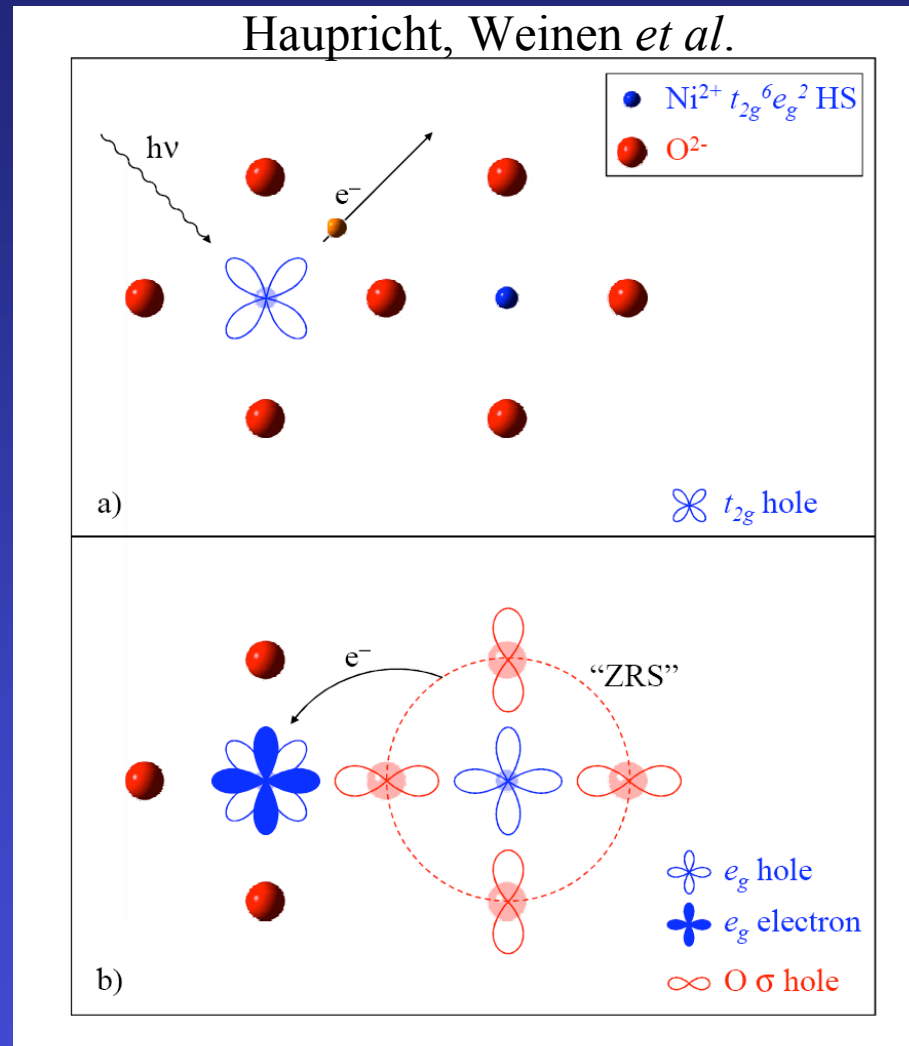
PES



What is the origin of peak B ??

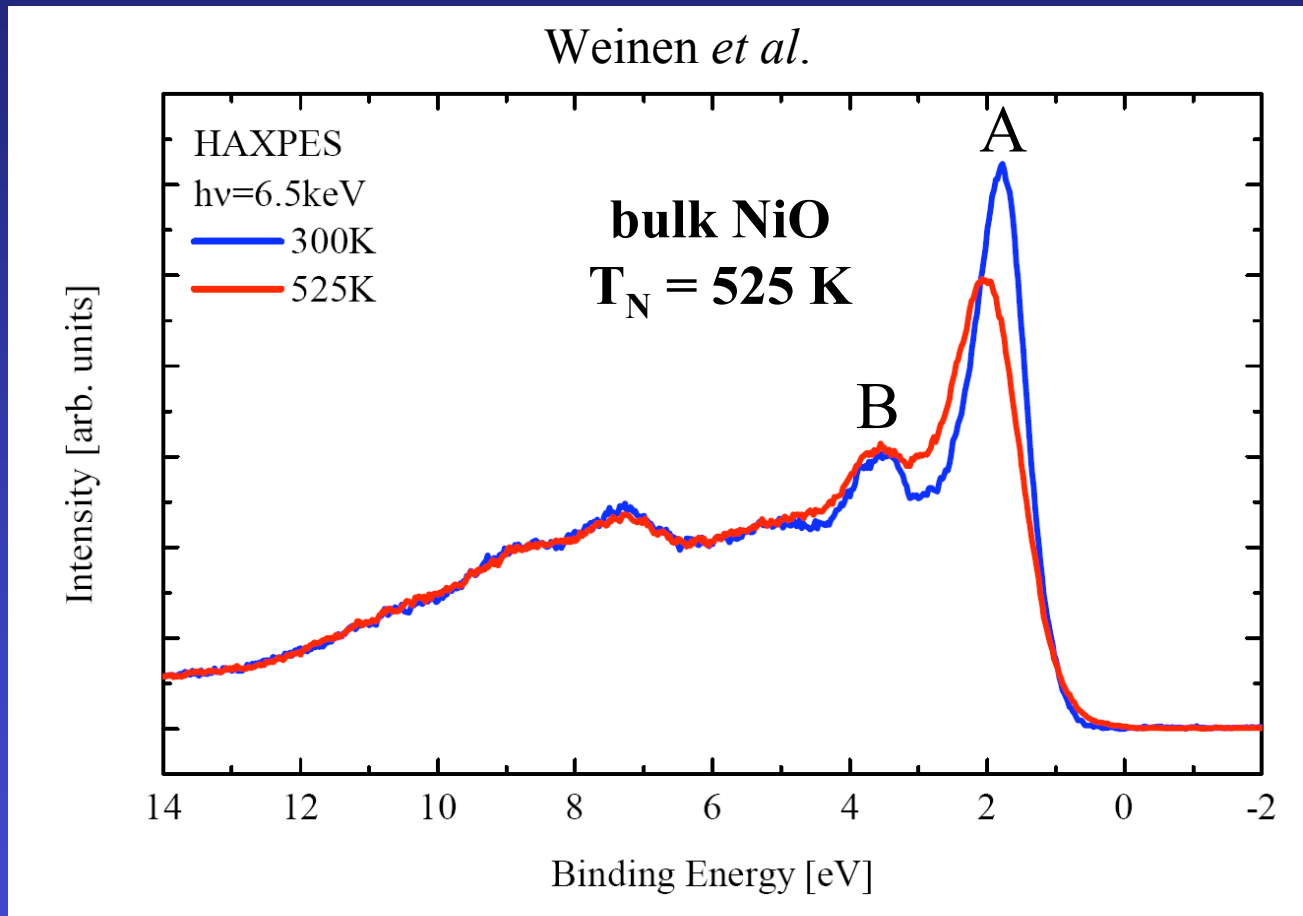
PES

“after” PES



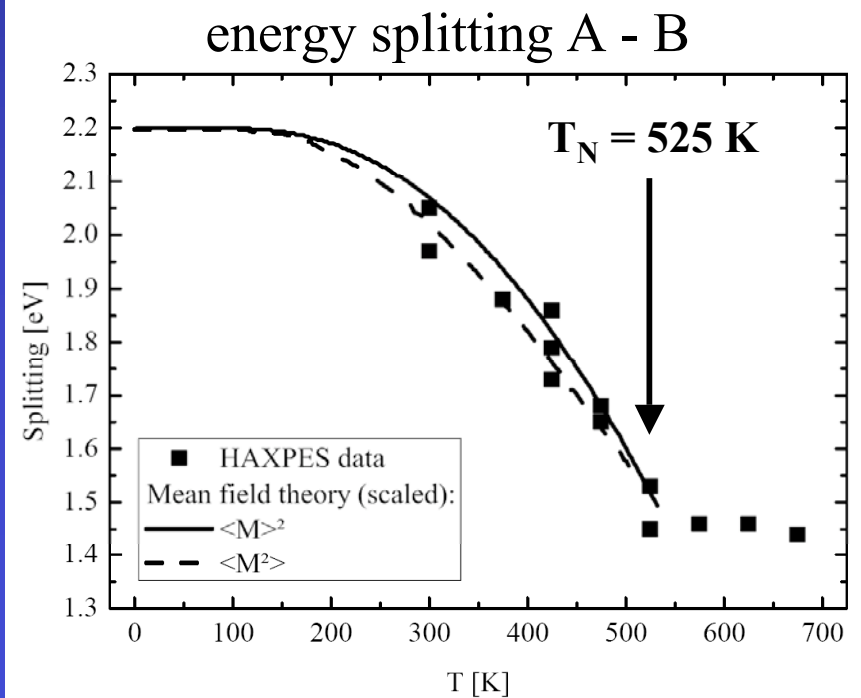
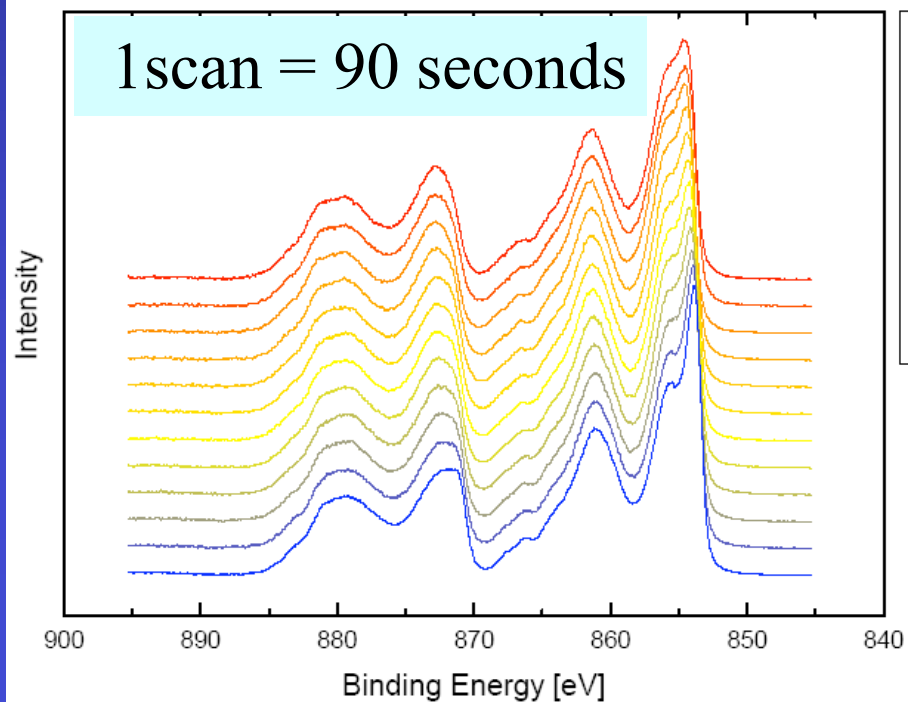
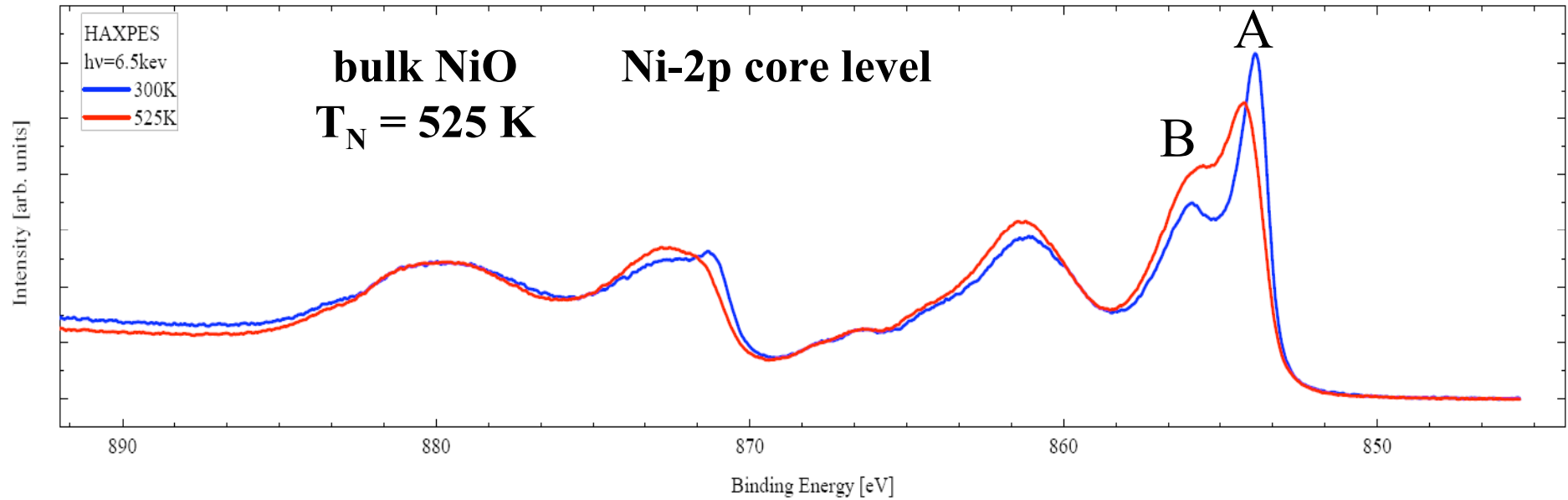
Screening from a neighboring Ni cluster,
creating a *many-body* hole state ('ZRS')

Testing non-local screening model: temperature dependence ?



Screening is reduced above T_N :
indeed, screening needs AF nearest neighbor spin alignment !

Weinen *et al.*



Concluding remarks:

LaCoO₃: a prototype spin state transition material

- full atomic multiplet theory is necessary
- influence of lattice is important

NiO: simple AFI system but very complex electronic structure

- correlations, multiplets, band formation, screening
- charge transfer insulator, compensated-spin 1st ionization state

Soft-x-ray absorption spectroscopy (XAS/MCD):

- extremely sensitive to valence, spin and orbital state

Photoelectron spectroscopy:

- very sensitive to short-range spin-spin correlations