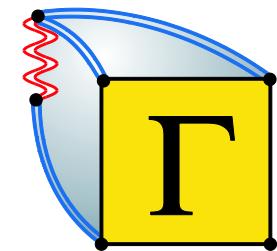


Dynamical vertex approximation (DΓA)

Karsten Held (TU Wien)

Jülich, September 18st, 2014

- Reducibility & parquet equation
- Methods: DΓA, DF, 1PI, DMF²RG
- Critical exponents 3D Hubbard model
- Fate of false Mott transition in 2D Hubbard model

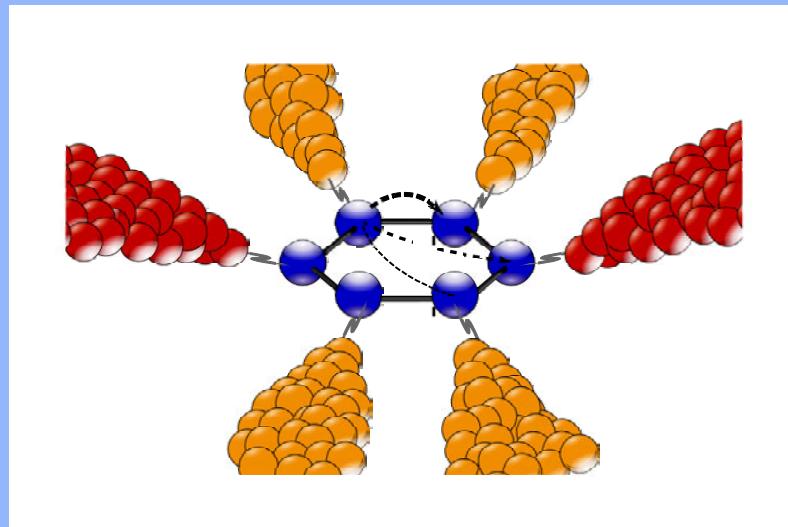


(a) quantum criticality

(b) high temperature superconductivity

(c) heterostructures

(d) nanostructures and molecules

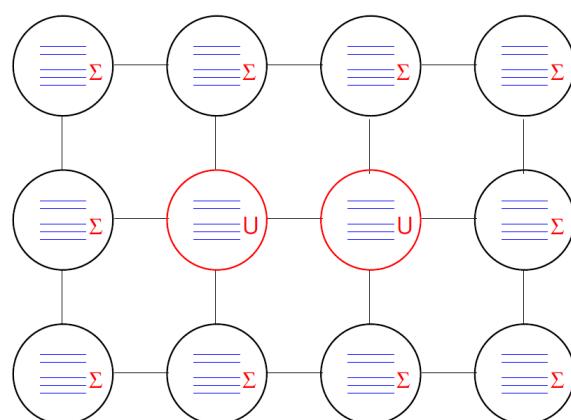


D Γ A: weak localization in nanostructures
quantum chemistry

Two routes beyond DMFT:

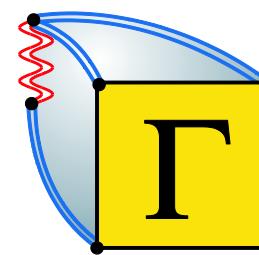
Correlations beyond DMFT:

d-, p-wave superconductivity, pseudogaps,
(para-)magnons, quantum criticality



Cluster extensions

diagrammatic extensions



D Γ A: Toschi, Katanin, KH'07

DF: Rubtsov et al.'08

... 1PI, DMF²RG

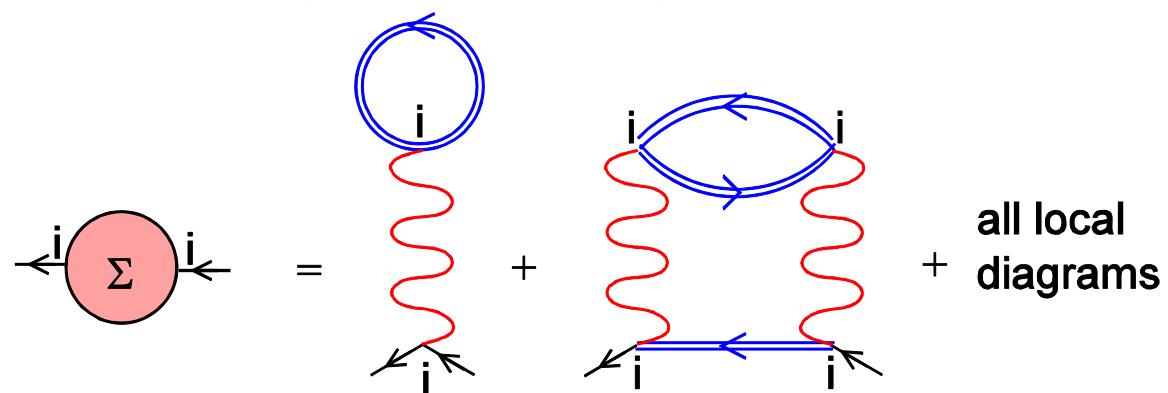
Hettler et al.'98, Lichtenstein, Katsnelson'00

Kotliar et al.'01, Potthoff et al.'03

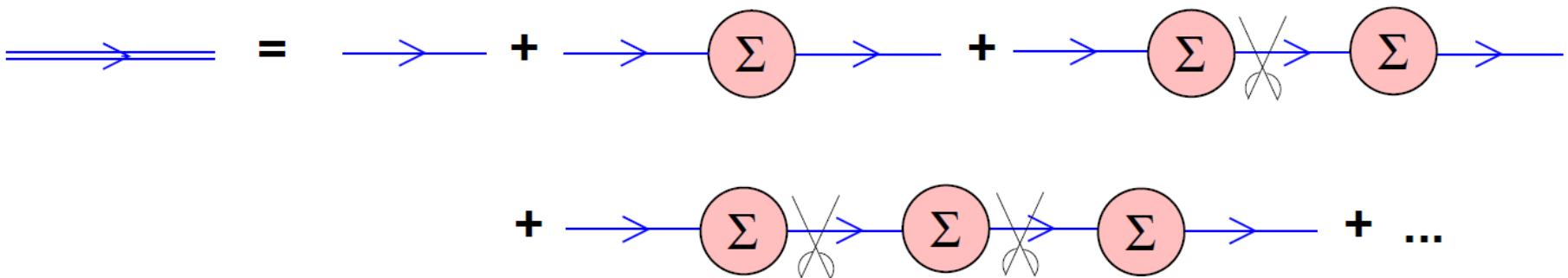
Resummation of Feynman diagrams in terms of locality *local n-particle fully irreducible vertex*

n=1: DMFT

- local 1-particle fully irreducible vertex Σ



Σ : one-particle irreducible one-particle vertex



Resummation of Feynman diagrams in terms of **locality** *local n-particle fully irreducible vertex*

n=1: DMFT

- local 1-particle fully irreducible vertex Σ

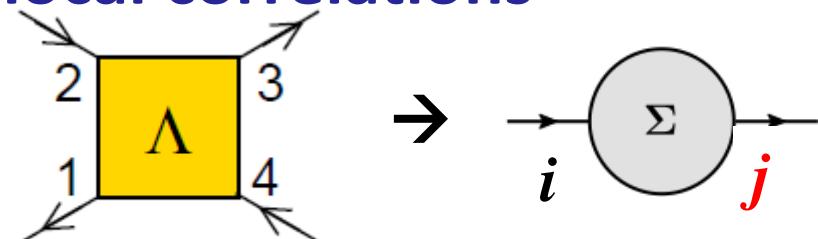
n=2: DΓA

- local 2-particle fully irreducible vertex Λ

→ non-local correlations

...

n→∞: exact



Two-particle irreducibility

$$\chi = - \frac{2}{3} + \frac{2}{3} - F$$

3 ways (channels) to cut two lines

$$F = \Lambda + \Phi_{ph} + \Phi_{\bar{ph}} + \Phi_{pp}$$

full vertex

parquet equation

Two-particle irreducibility

Parquet equation:

$$F = \Lambda + \Phi_{ph} + \Phi_{\bar{p}\bar{h}} + \Phi_{pp}$$

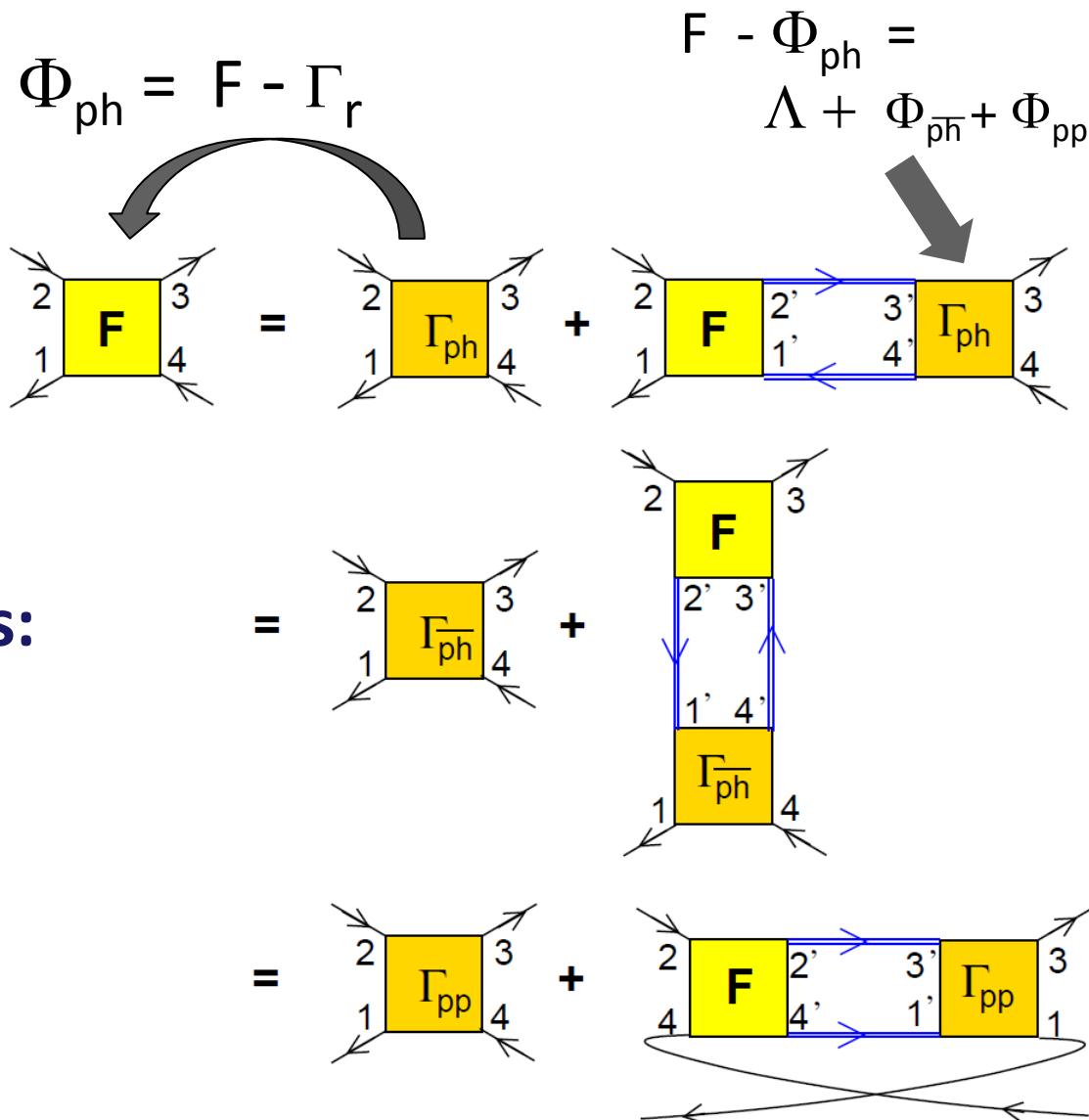
$$F = \Phi_r + \Gamma_r$$

$r \in \{ph, \bar{p}\bar{h}, pp\}$

defines Γ_r

Bethe Salpeter equations:

$$F = \Gamma_r + \int F G G \Gamma_r$$



Parquet equations

$$\begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Phi_{ph}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} = \boxed{F} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Lambda} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} + \boxed{F} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Phi_{ph}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} + \boxed{F} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Phi_{pp}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array}$$

$$\begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Phi_{ph}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} = \boxed{F} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Lambda} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} + \boxed{F} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Phi_{ph}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} + \boxed{F} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Phi_{pp}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array}$$

$$\begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Phi_{pp}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} = \boxed{F} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Lambda} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} + \boxed{F} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Phi_{ph}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} + \boxed{F} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{\Phi_{ph}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array}$$

Λ given

$\rightarrow F, \Phi_{ph}, \Phi_{\bar{ph}}, \Phi_{pp}, \Sigma, G$

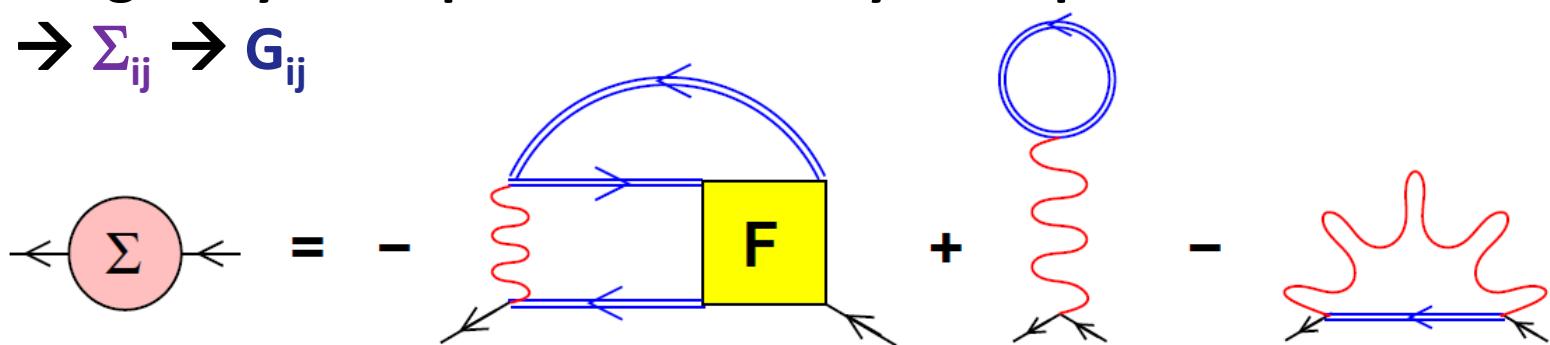
$\Lambda = U$ (parquet approx.)

$\Lambda = \Lambda_{loc}$ ($D\Gamma A$)

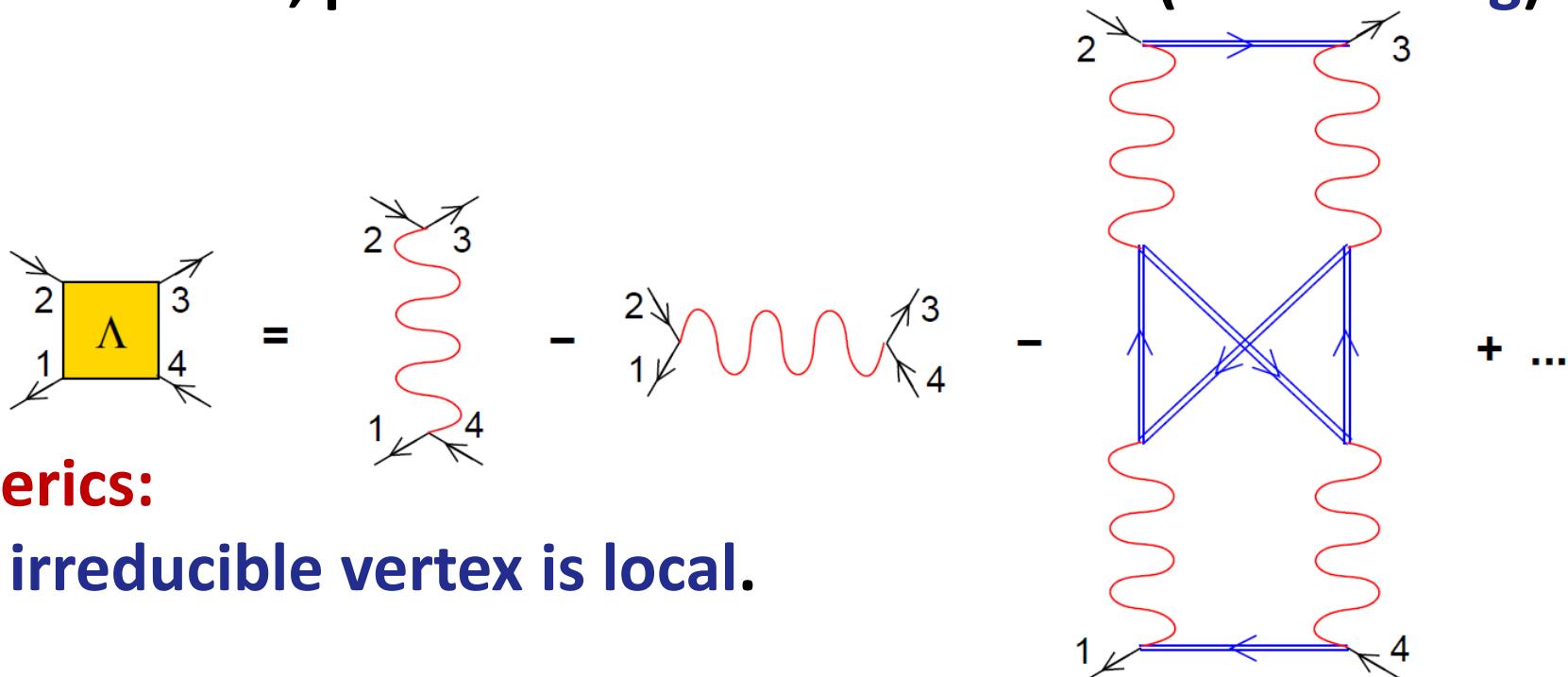
$$\begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \boxed{F} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} = \boxed{\Lambda} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} + \boxed{\Phi_{ph}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} + \boxed{\Phi_{\bar{ph}}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} + \boxed{\Phi_{pp}} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array}$$

Schwinger-Dyson eq. of motion & Dyson equation

$F_{ijkl} \rightarrow \Sigma_{ij} \rightarrow G_{ij}$



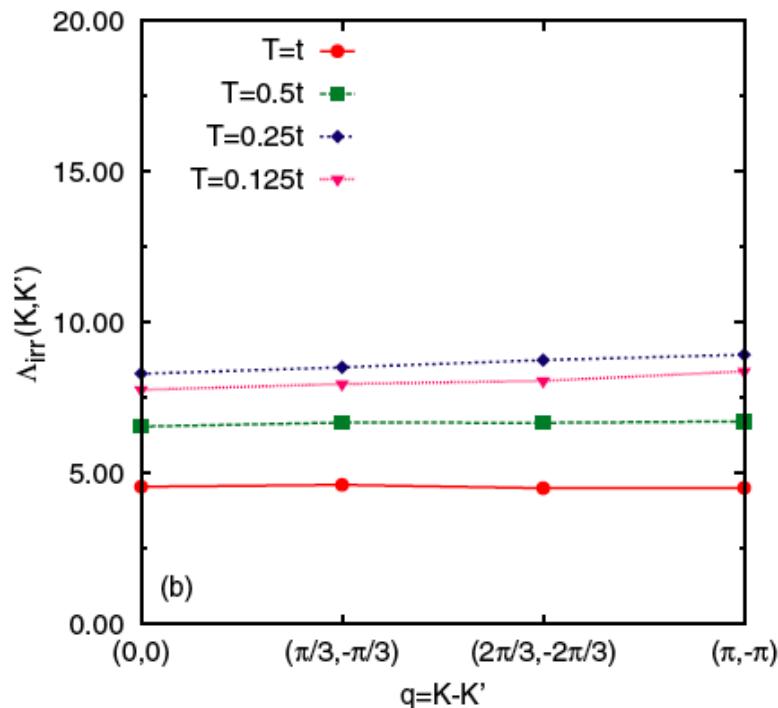
- 1) Physics:** our understanding either based on
1-particle [qp renormalization, Mott transition etc...] or
2-particle level [(para)magnons, (quantum)criticality]
- 2) Diagrammatics:**
includes DMFT, plus non-local contributions (**few missing**)



- 3) Numerics:**
fully irreducible vertex is local.

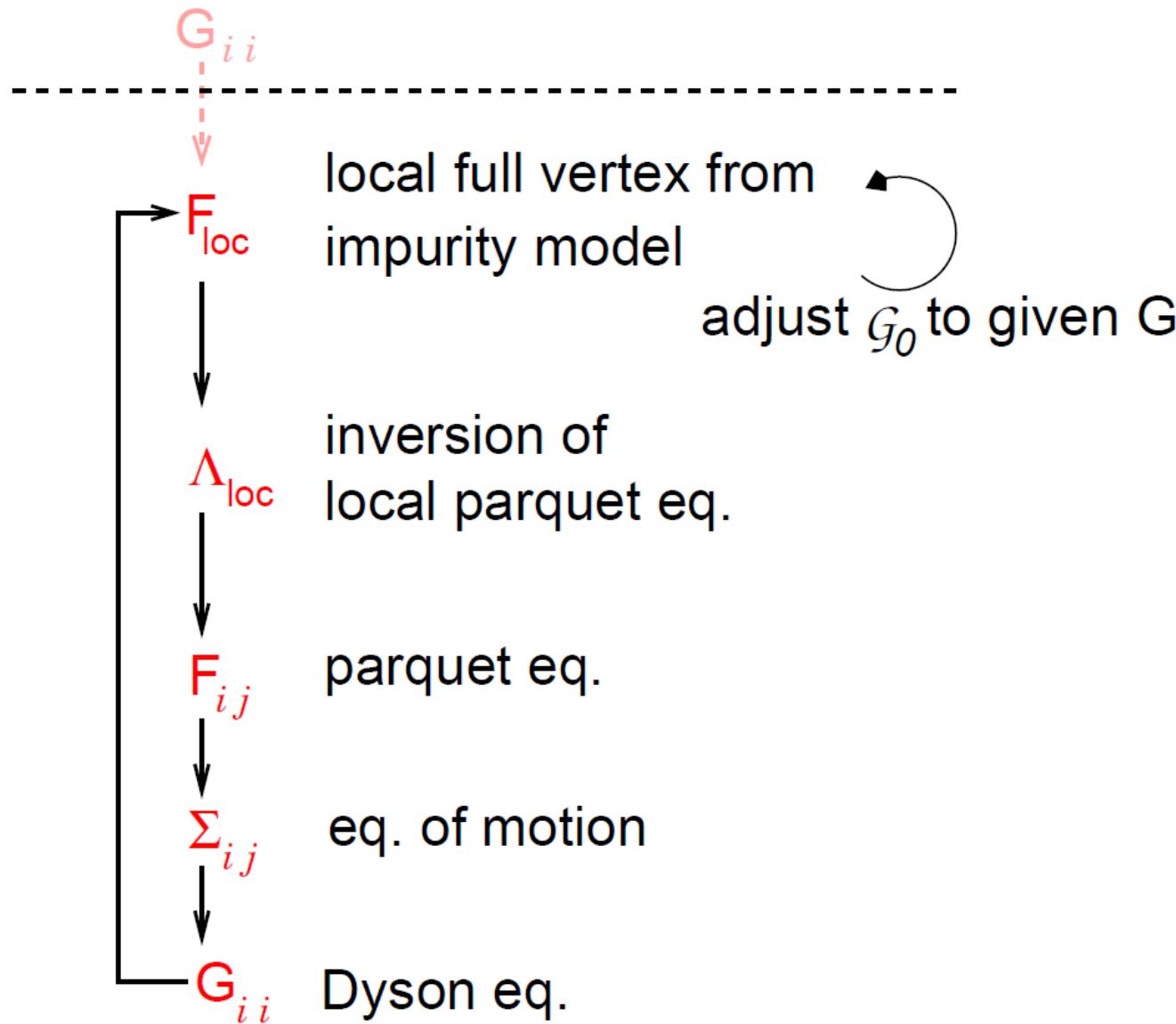
3) Fully irreducible vertex is local!

The fully irreducible vertex Λ_{irr} is **local** (k -independent) !



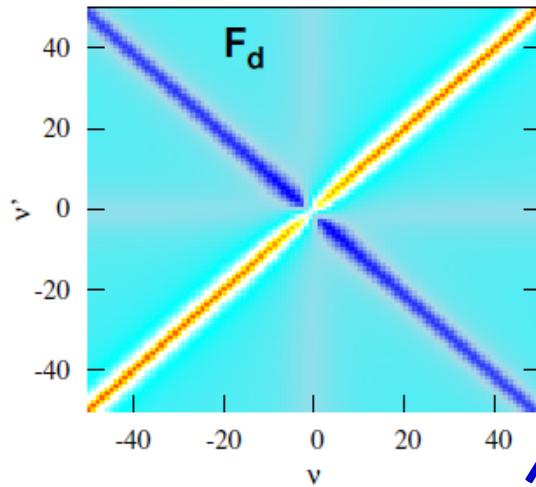
DCA, 2D-Hubbard model, $U=4t$, $n=0.85$,
 $v=v'= \pi/\beta$, $\omega=0$
Maier et al. PRL 96 (2006)

DΓA algorithm



Calculation of local irreducible vertex

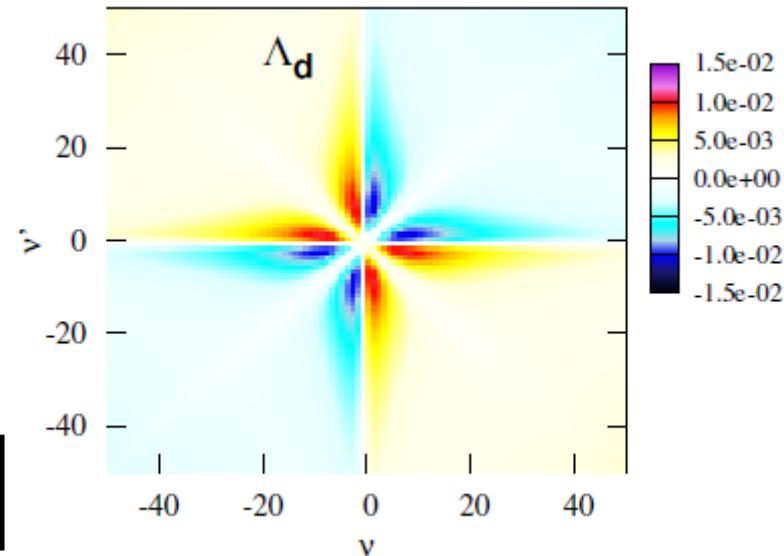
full vertex F (from SIAM, QMC or ED)



inversion of parquet eq.

Rohringer, Valli, Toschi '12

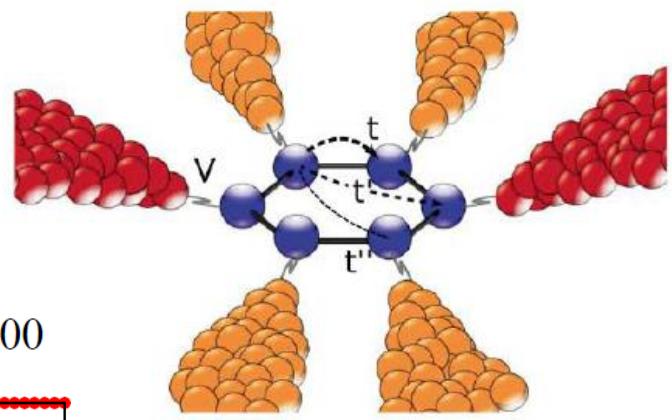
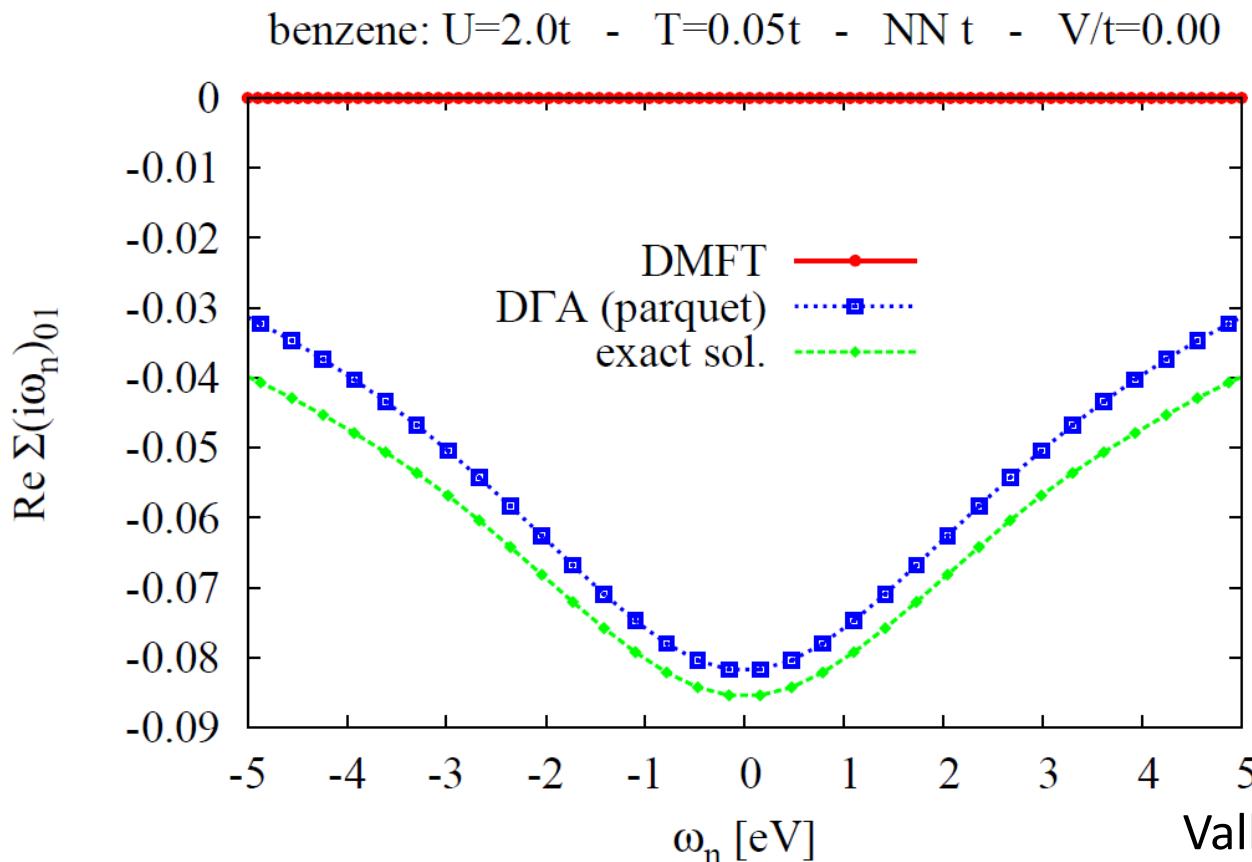
fully irreducible vertex Λ



Λ ω -dep. but k -indep.!

DΓA with parquet equation

Benzene Hubbard ring:
comparison non-local Σ
to exact solution



Other 2-particle vertex approaches:

Dual Fermions (DF):

Rubtsov, Lichtenstein, Katsnelson '08

functional integral with action of dual Fermions

$$S[f, f^*] = \sum_{\omega k \sigma} g_\omega^{-2} ((\Delta_\omega - \epsilon_k)^{-1} + g_\omega) f_{\omega k \sigma}^* f_{\omega k \sigma} + \sum_i V_i$$

$$V[f_i, f_i^*] = -\gamma_{1234}^{(4)} f_1^* f_2 f_3^* f_4 + \gamma_{123456}^{(6)} \cancel{f_1^* f_2 f_3^* f_4 f_5^* f_6} + \dots$$

1-particle reducible vertex γ

Rohringer et al. PRB '13

1PI: functional integral using 1-particle irreducible vertex

$$\begin{aligned} Z[\eta, \eta^+] = & \int d[\tilde{c}, \tilde{c}^+] \exp \left\{ \tilde{c}_{k\sigma}^+ [\zeta_\nu^{-1} - G_{0k}^{-1}]^{-1} \tilde{c}_{k\sigma} \right. \\ & \left. + (\eta_{k\sigma}^+ + \tilde{c}_{k\sigma}^+) \phi_{k\sigma} + \phi_{k,\sigma}^+ (\eta_{k\sigma} + \tilde{c}_{k\sigma}) - \Gamma_{\text{DMFT}}[\phi, \phi^+] \right\} \end{aligned}$$

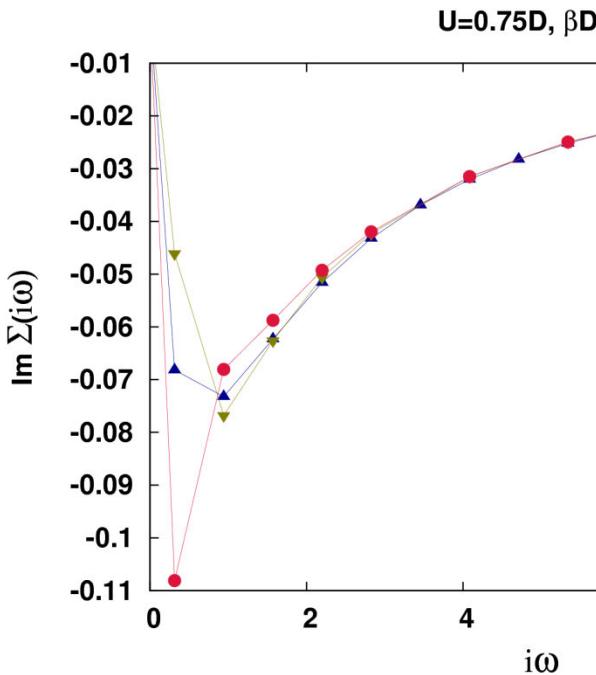
1PI unifies positive features of DF and DΓA

Other 2-particle approaches: DMF²RG

Taranto et al.PRL'04

DMFT 1PI vertex² Γ → fRG

2d Hubbard model



$$\frac{d}{d\Lambda} \Sigma^\Lambda = S^\Lambda \Gamma^{(4)\Lambda}$$

taken from
Metzner et al.'12

$$\frac{d}{d\Lambda} \Gamma^{(4)\Lambda} = \Gamma^{(4)\Lambda} G^\Lambda \Gamma^{(4)\Lambda} + \cancel{\Gamma^{(6)\Lambda}}$$

Overview 2-particle vertex approaches:

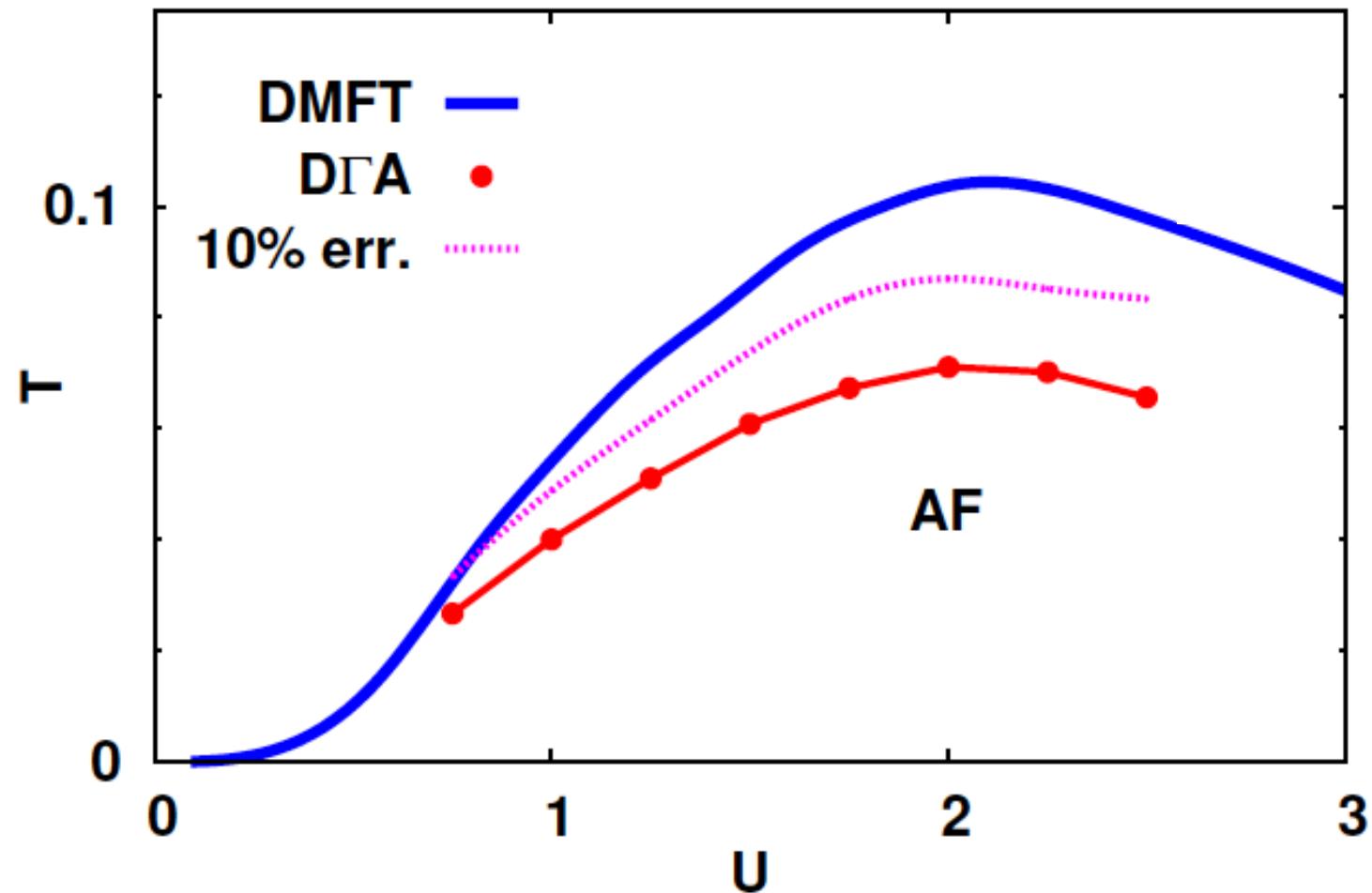
| Method | local 2-particle vertex | Feynman diags |
|---------------------|---|-------------------------------|
| DF | 1-particle reducible vertex, here F_{loc} | 2nd order, ladder, parquet |
| 1PI | 1-particle irreducible vertex F_{loc} | ladder |
| DMF ² RG | 1-particle irreducible vertex F_{loc} | RG flow |
| ladder DΓA | 2-particle irreducible vertex in $r \Gamma_r$ | ladder |
| full DΓA | 2-particle fully irreducible vertex Λ | parquet |

D Γ A results: 3D and 2D Hubbard model

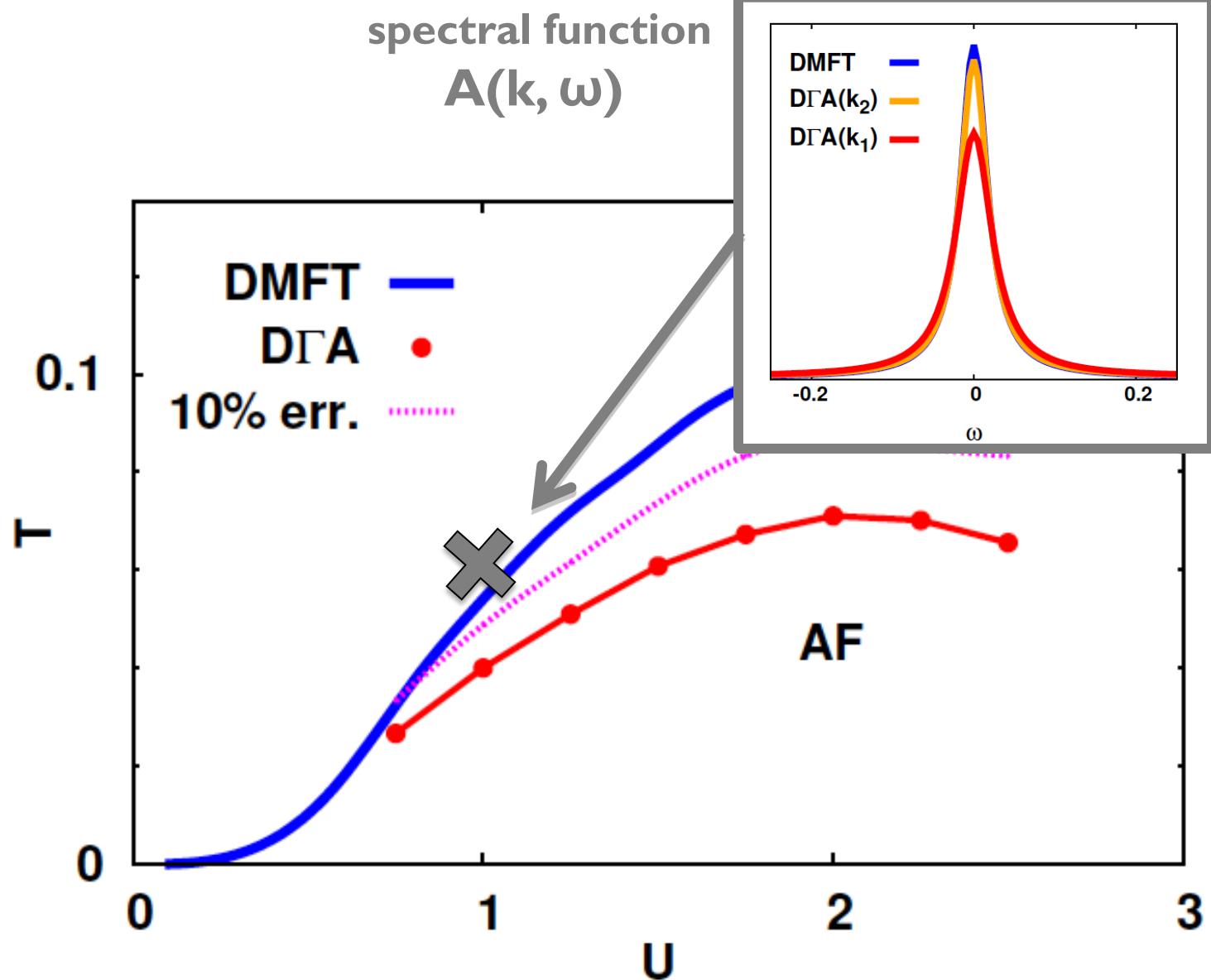
D Γ A results in 3 dimensions

Phase diagram: Hubbard model in $d=3$ (cubic lattice $D=1$, half-filling)

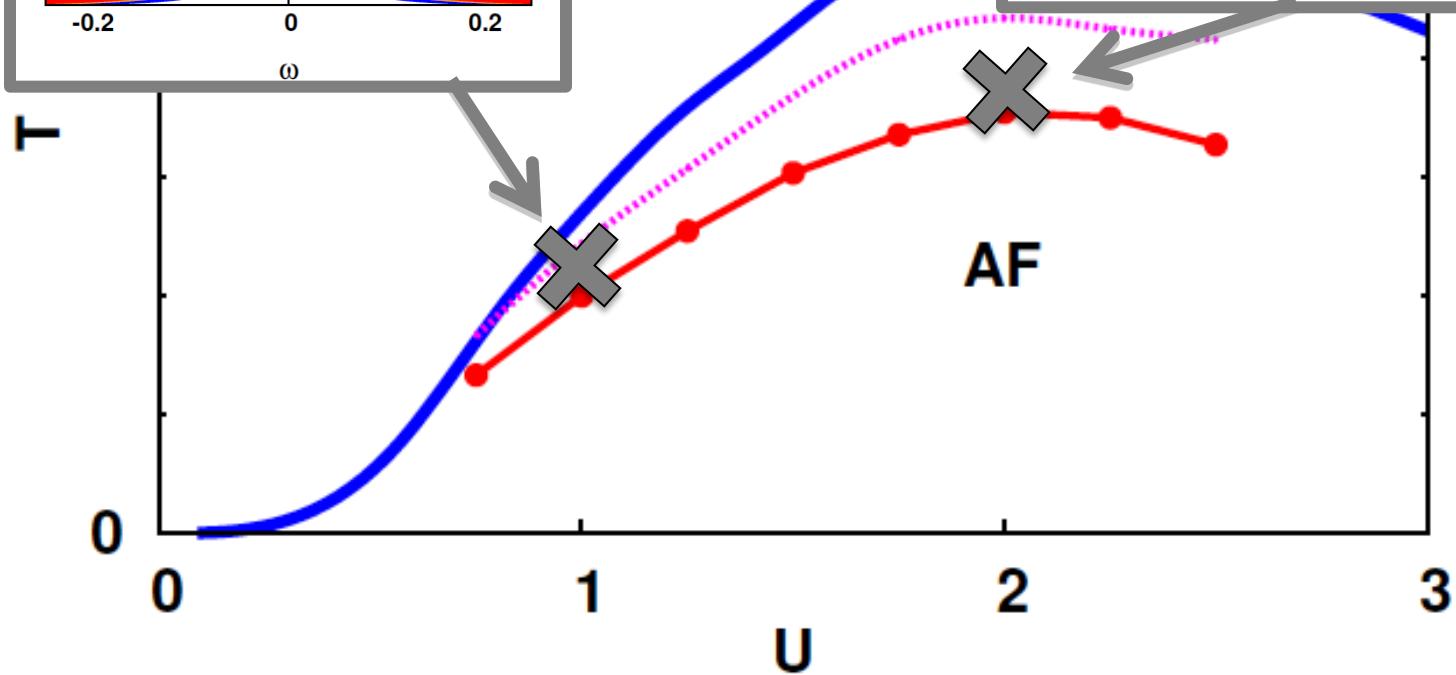
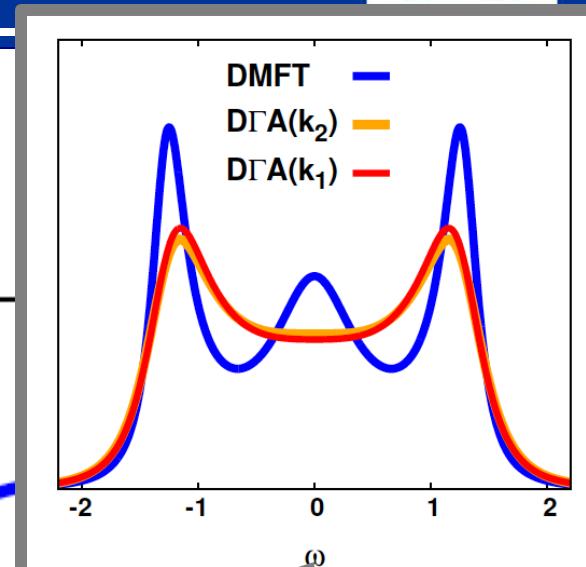
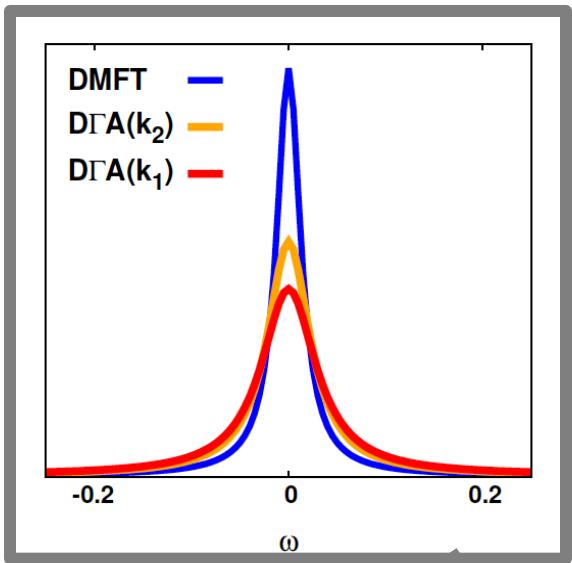
G. Rohringer et al., PRL (2011)



DΓA results in 3 dimensions

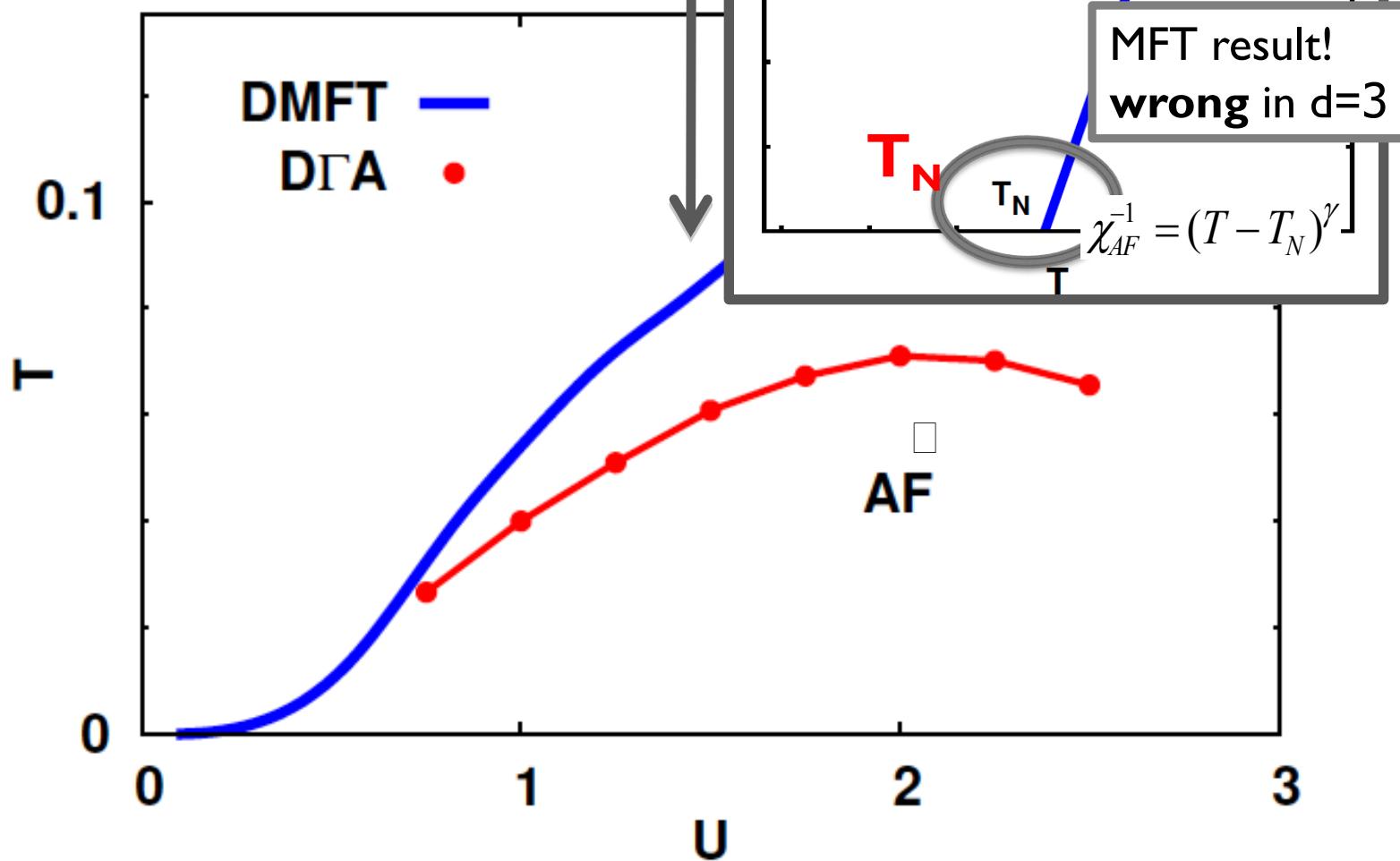


D Γ A results in 3 dimensions



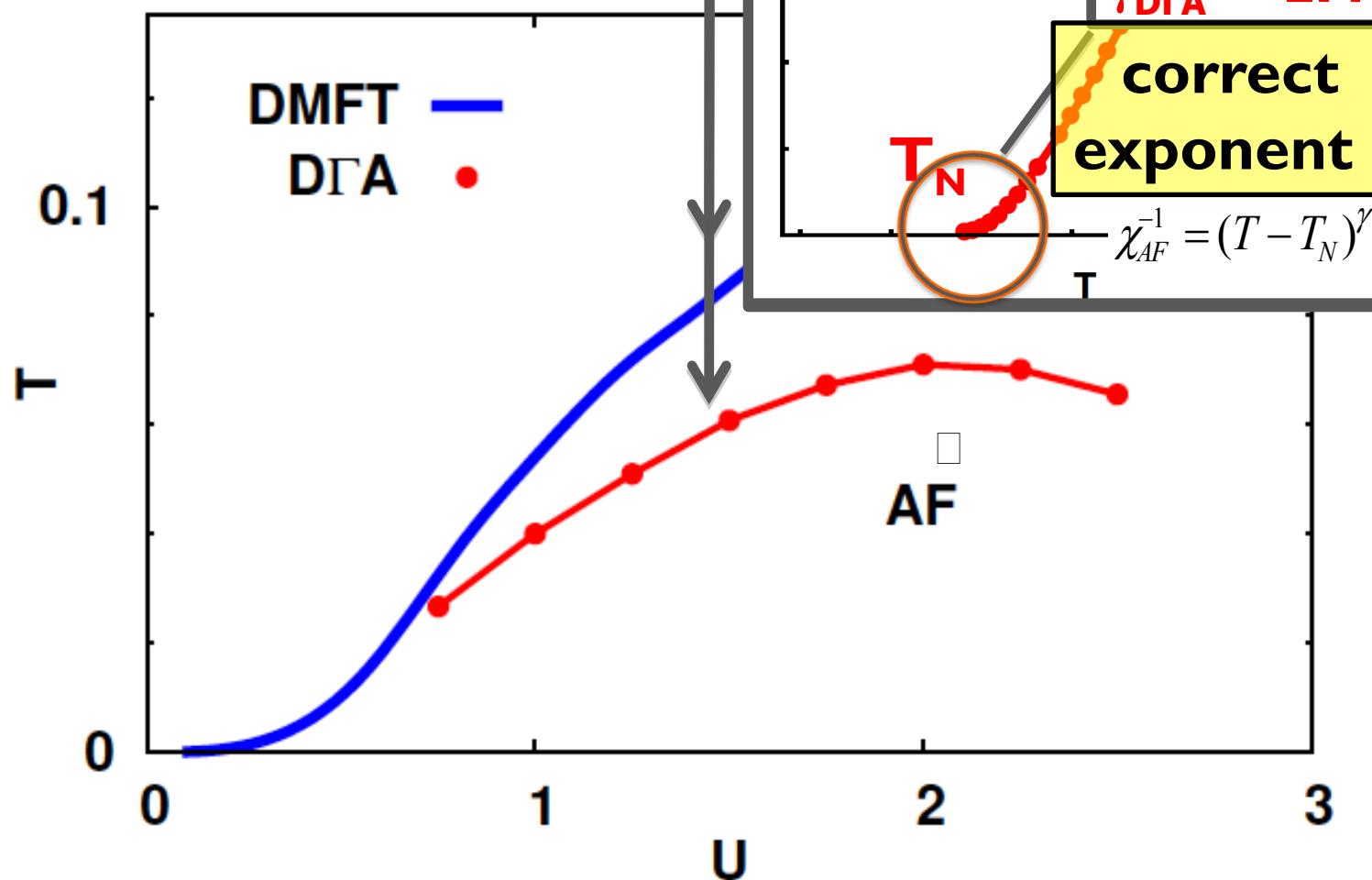
DΓA results in 3 dimensions

$$\chi_{AF}^{-1} = \chi_S^{-1}(q = (\pi, \pi, \pi))$$

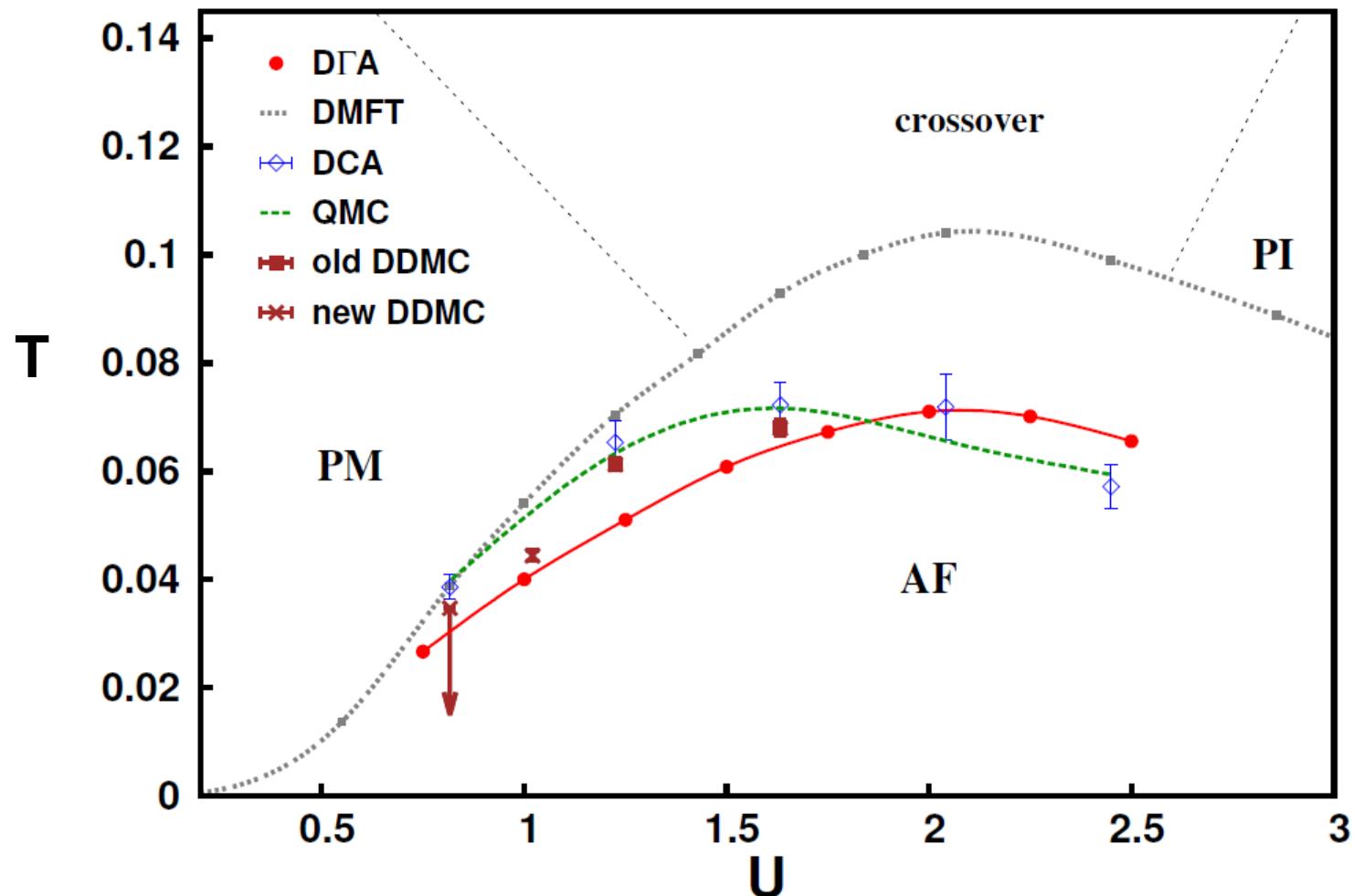


DΓA results in 3 dimensions

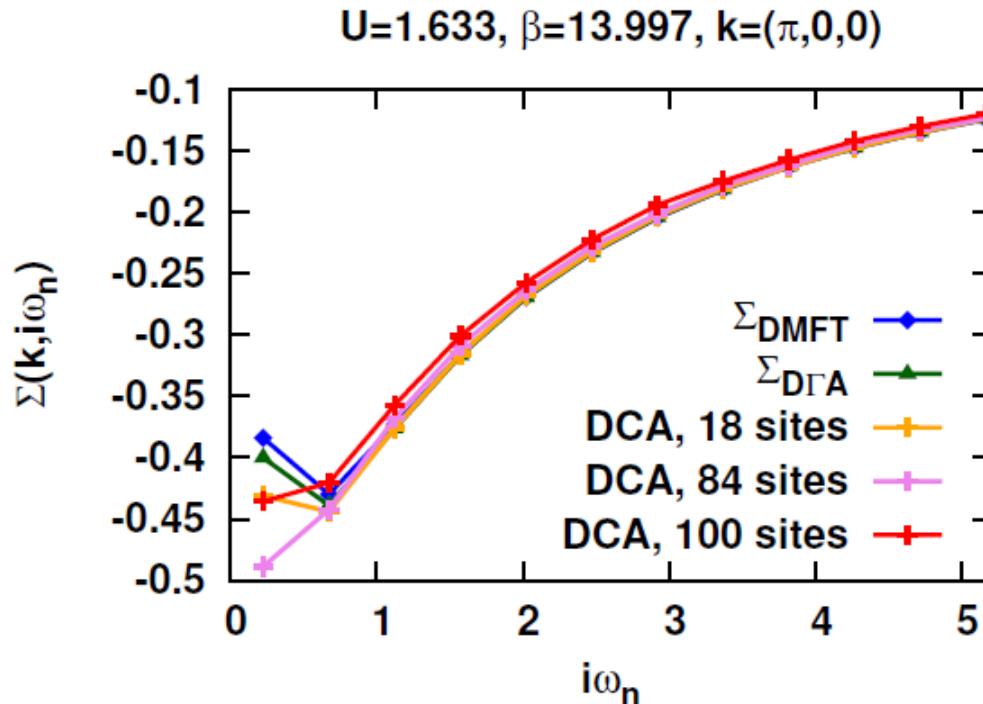
$$\chi_{AF}^{-1} = \chi_S^{-1}(q = (\pi, \pi, \pi))$$



Comparison with DCA, QMC



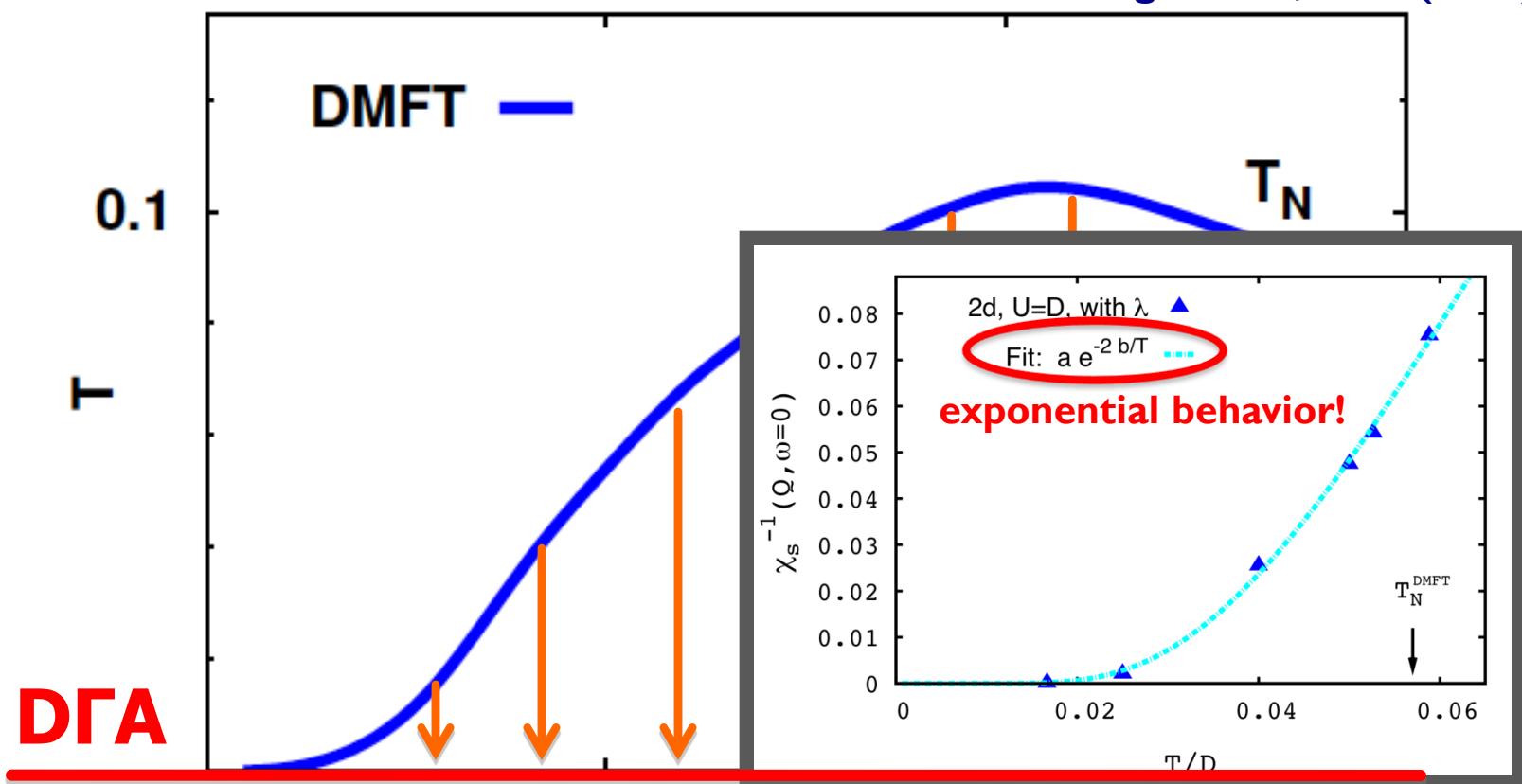
Comparison of self energy



DΓA vs. DCA by Gull et al.
3d Hubbard ($D=1$)
Rohringer et al'13

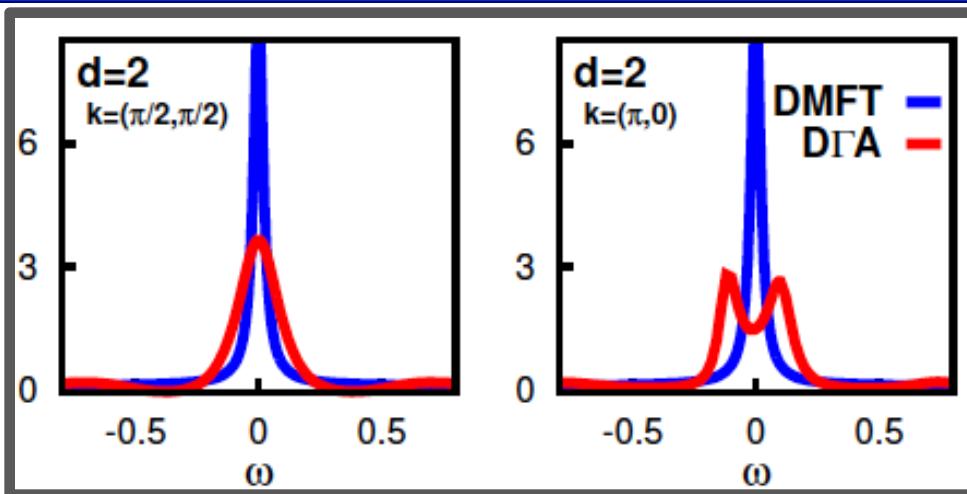
DΓA results in 2 dimensions

Phase diagram:
Hubbard model in $d=2$
G. Rohringer et al., PRL (2011)

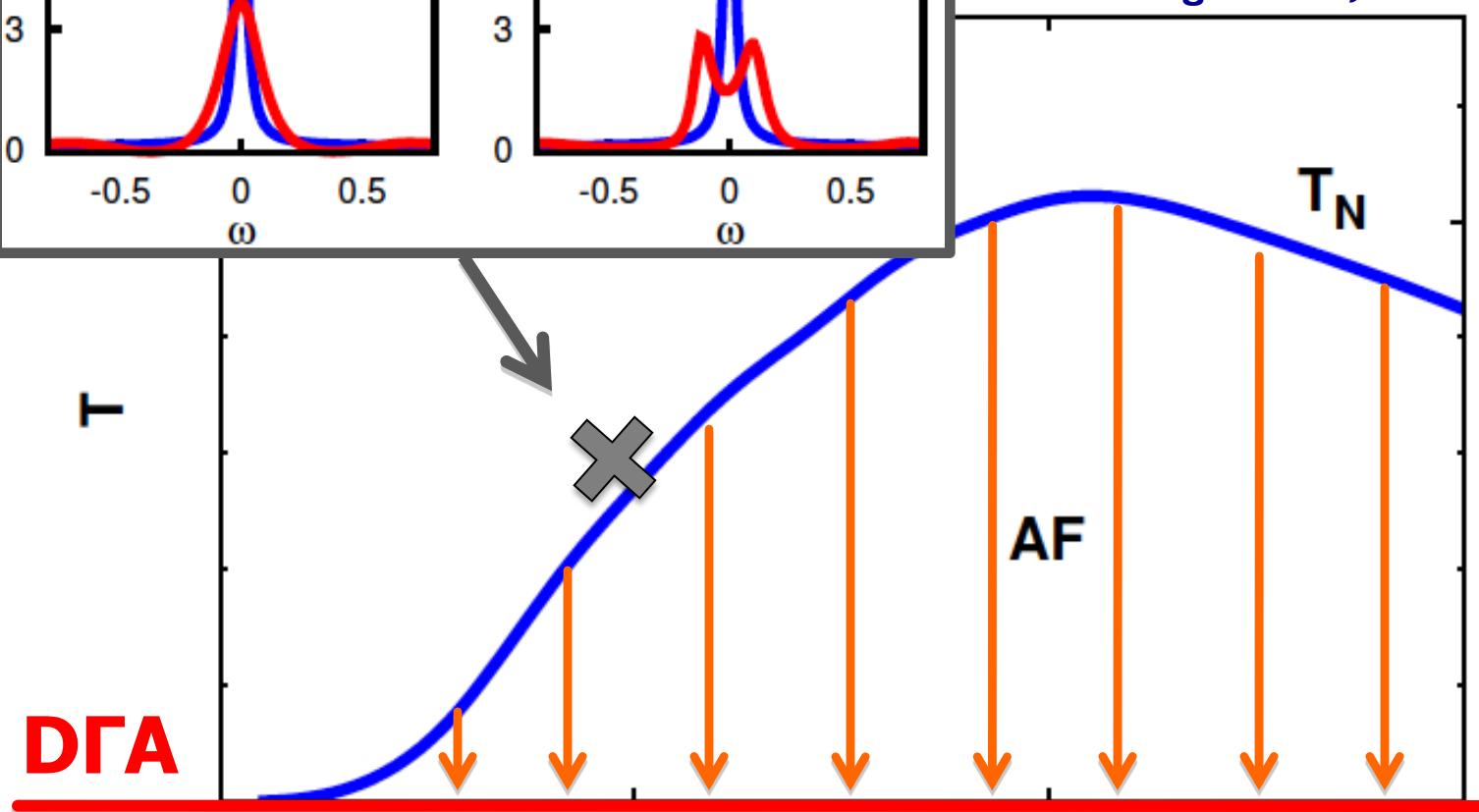


$T_N = 0 \rightarrow$ Mermin-Wagner Theorem in $d = 2!$

DΓA results in 2 dimensions



Phase diagram:
Hubbard model in **d=2**
G. Rohringer et al., PRL (2011)

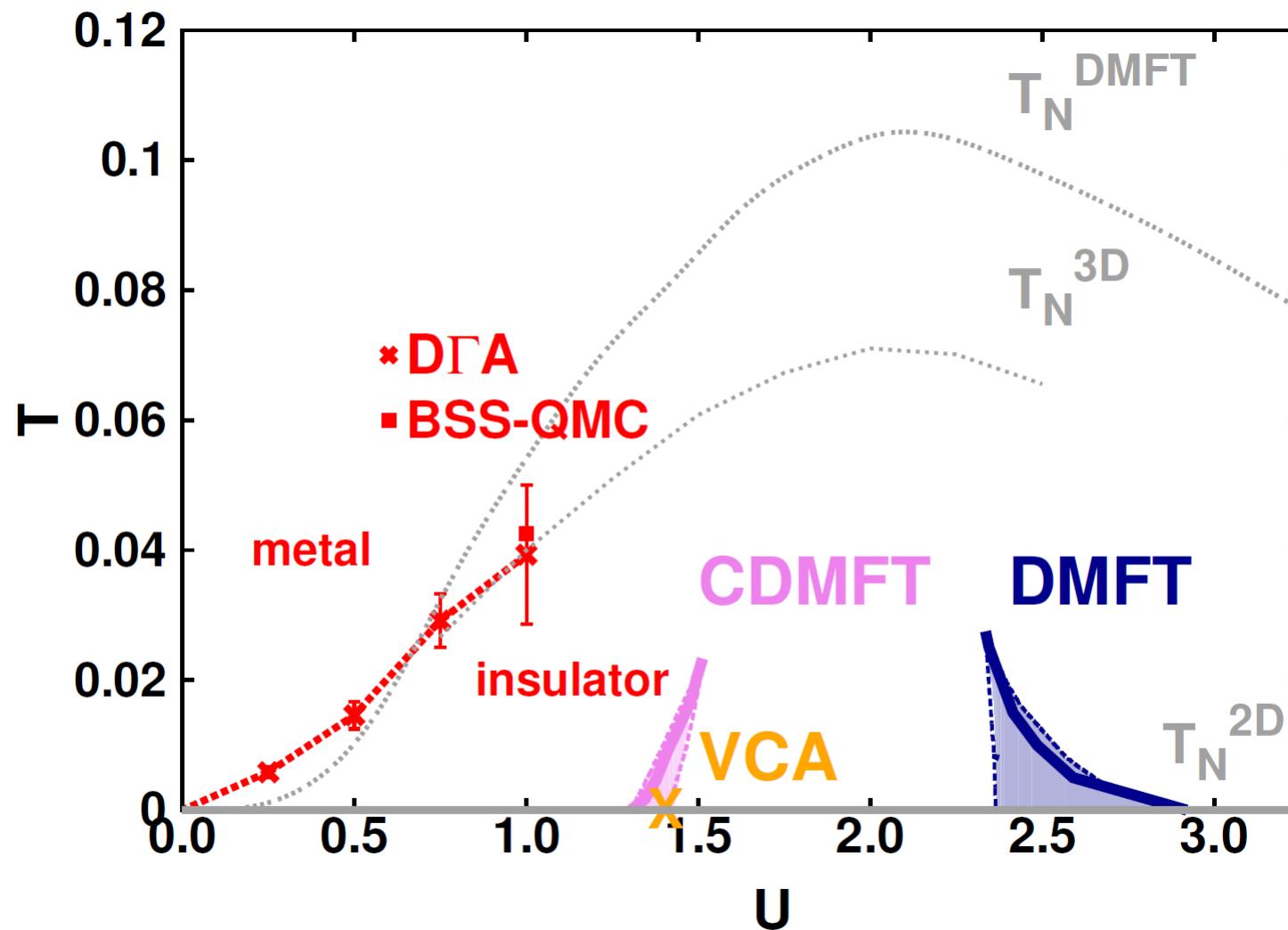


$T_N = 0 \rightarrow$ Mermin-Wagner Theorem in $d = 2!$

Fade of false Mott transition in 2D

T. Schäfer et al. arXiv:1405.7250

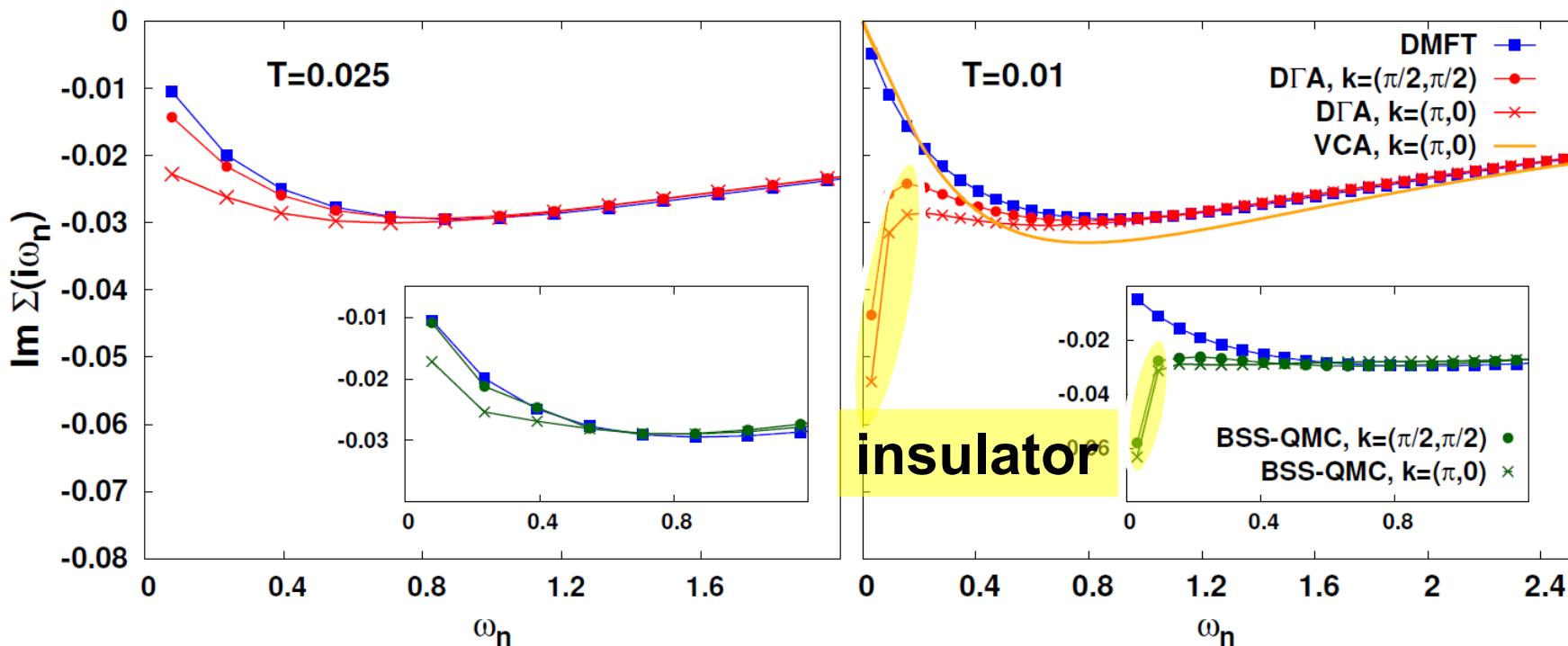
2D Hubbard model on square lattice ($4t \equiv 1$)



Fade of false Mott transition in 2D

T. Schäfer et al. arXiv:1405.7250

2D Hubbard model on square lattice (**U=0.5 * 4t**)

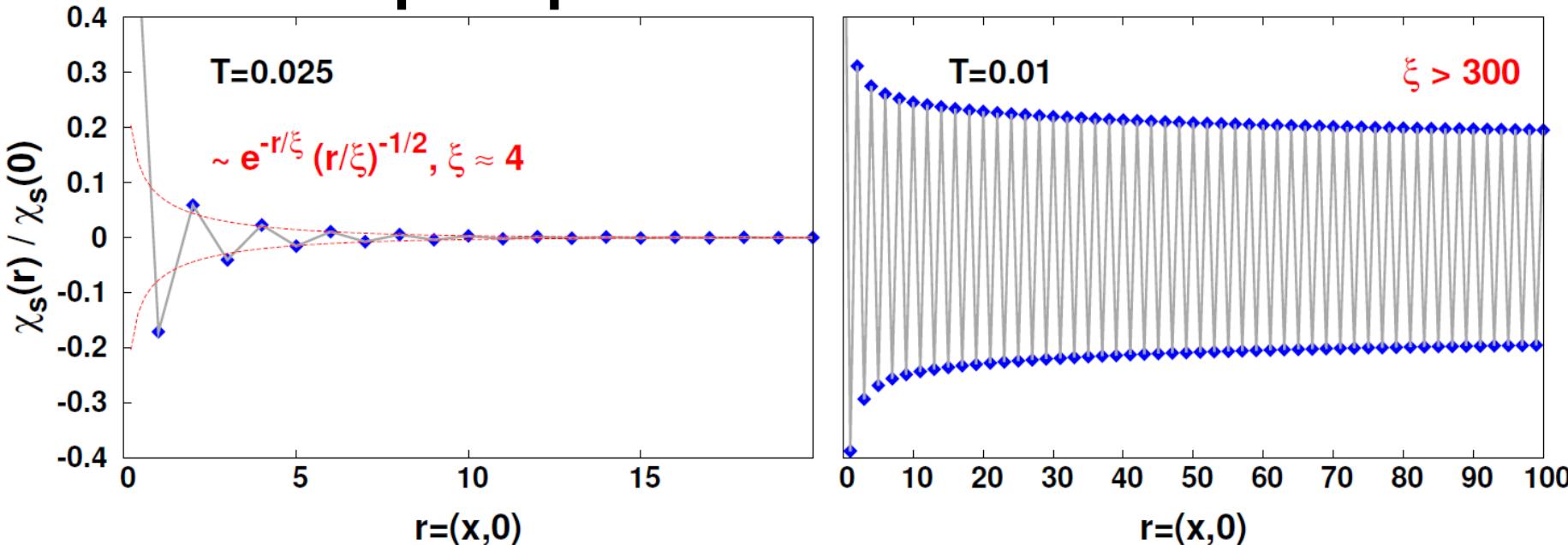


Fade of false Mott transition in 2D

T. Schäfer et al. arXiv:1405.7250

2D Hubbard model on square lattice (**U=0.5 * 4t**)

spin-spin correlation function

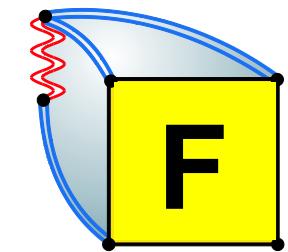


Origin of gap: **long-range antiferromagnetic correlations of Slater type (H_U energy gain)**

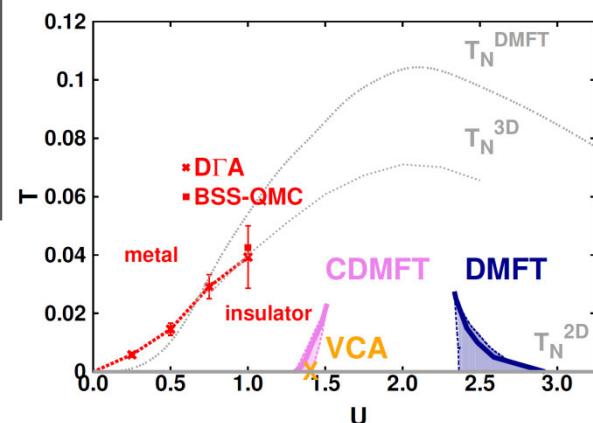
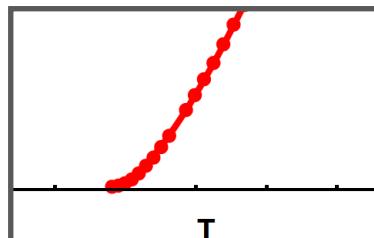
- ✗ no effective Heisenberg model (Anderson'87)
- ✓ similar to TPSC (Tremblay et al. '07)

Conclusion

- **DF:** reducible local vertex
- **1PI & DMF²RG:** 1-particle irreducible local vertex
- **DΓA:** fully, 2-particle irreducible local vertex Λ
 → full vertex F via parquet equations



- **DΓA:** critical exponents
- **DΓA:** metal-insulator transition
 → $U_c=0$



- A. Katanin (Ekaterinburg)
- A. Toschi, G. Rohringer, A. Valli, C. Tarato, T. Schäfer
(TU Wien)
- S. Andergassen (Uni Wien), W. Metzner (Stuttgart)

Further reading:

PRB 75, 045118 (2007), PRB 80, 075104 (2009),
Ann. Phys. 523, 698 (2011), PRL 107, 256402 (2011)
PRB 88, 115112 (2013), arXiv:1405.7250

PhD, postdoc
position
available

Vienna Computational Materials
Laboratory



European Research Council