

DFT+DMFT for Oxide Heterostructures

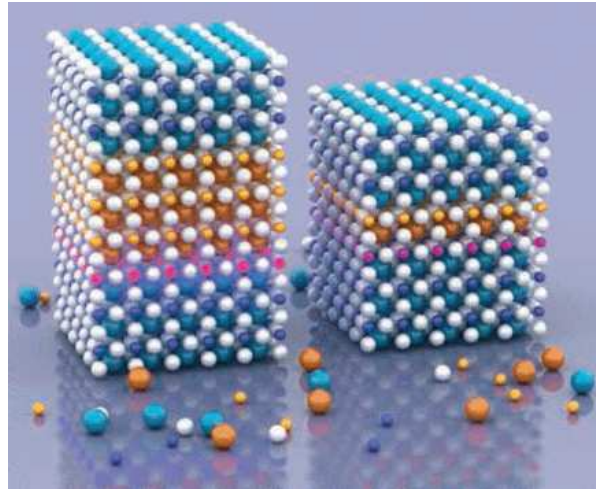
Frank Lechermann

Theoretische Physik III, Ruhr-Universität Bochum, Germany

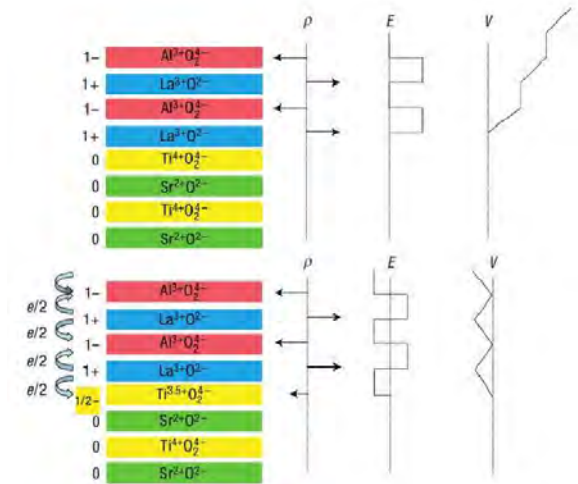
**Autumn School on Correlated Electrons:
Dynamical Mean-Field Theory of Correlated Electrons
4-7 October 2022, Forschungszentrum Jülich**

Oxide Heterostructures : (selected) Phenomenology

- superlattice/thin film build from
 - SrTiO_3 (band insulator)
 - LaAlO_3 (band insulator)
 - LaTiO_3 (Mott insulator)
 - \vdots
- possible quasi-twodimensional electron system (2DES)
- instabilities : ferromagnetism, superconductivity, . . .



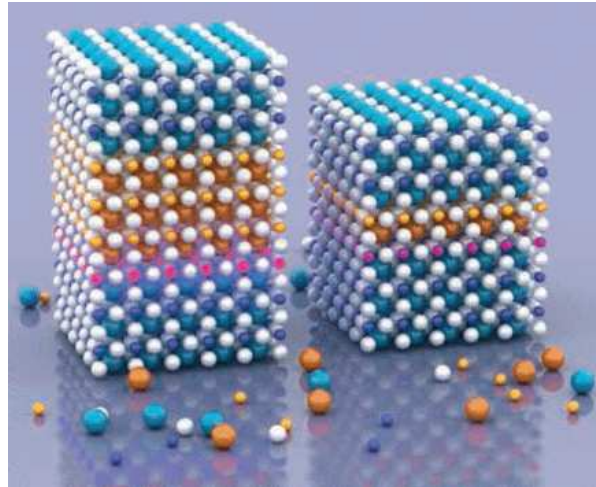
e.g. [Hwang et al., Nature Mat. 11, 103 (2012)]



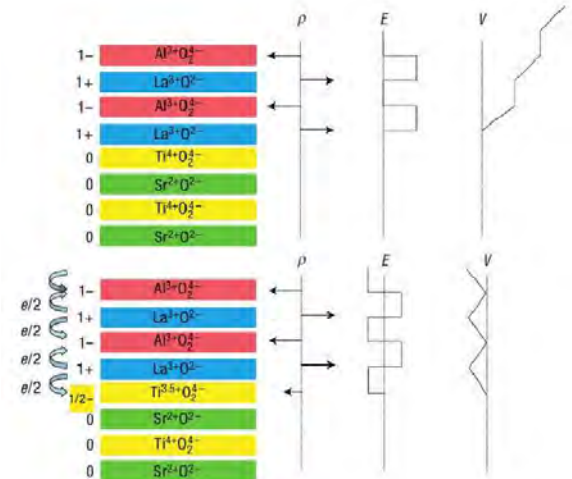
[Nakagawa et al., Nature Mat. 5, 204 (2006)]

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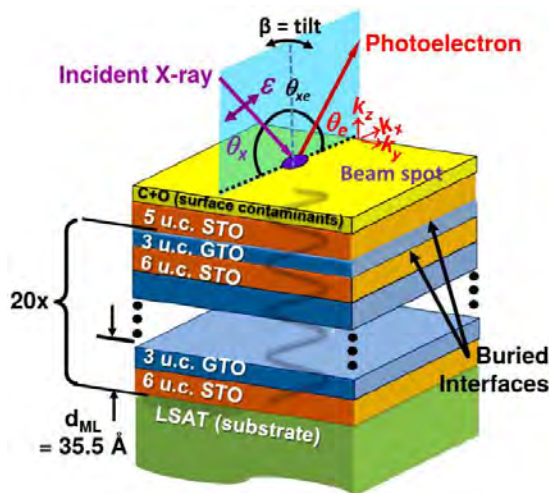


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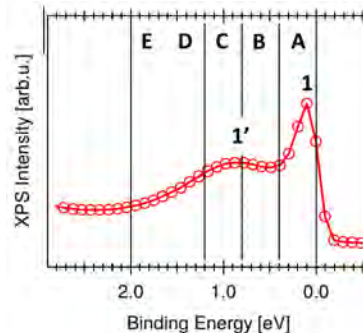


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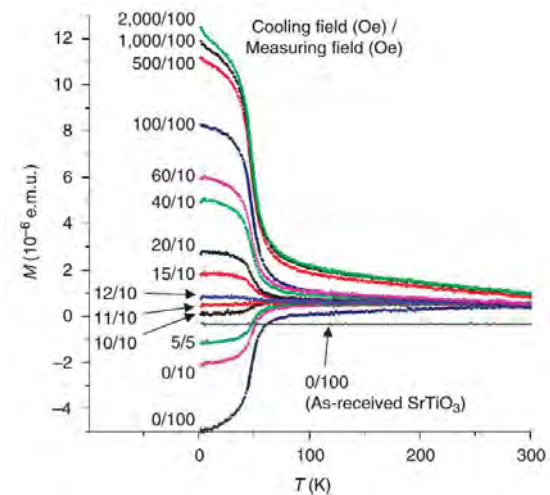
band-Mott insulators e.g. $\text{SrO}/\text{GdTiO}_3$



[Nemšák et al., PRB 93, 245103 (2016)]



band-band insulators e.g. $\text{LaAlO}_3/\text{SrTiO}_3$



[Ariando et al., Nat. Commun. 2, 188 (2015)]

Oxide Heterostructures : Theoretical Work

- model-Hamiltonian dynamical mean-field theory (DMFT) for Mott/band-insulator interfaces
[Okamoto and Millis, Nature 428, 630 (2004)]
- density functional theory (DFT) for band-insulating heterostructures
[Popovic and Satpathy, PRL 94, 176805 (2005)]
- DFT plus static Hubbard correlations (DFT+U) for band- and Mott heterostructures
[Pentcheva and Pickett, PRL 99, 016802 (2007)] [Pavlenko et al., PRB 86, 064431 (2012)]
- model-Hamiltonian field-theoretical approach to $\text{LaAlO}_3/\text{SrTiO}_3$
[Joshua et al., Nat. Commun. 3, 1129 (2012)]
- (realistic) Anderson-Hamiltonian exact-diagonalization/Lanczos method for isolated-defect SrTiO_3
[Lin and Demkov, PRL 111, 217601 (2014)]

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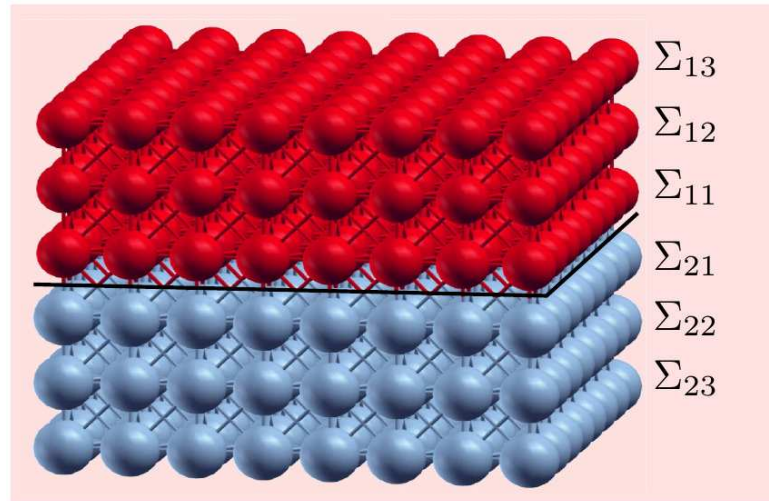
this talk

- ★ realistic many-body approach via DFT+DMFT
- ★ correlation physics by design

★ other recent DFT+DMFT approaches to oxide heterostructures

[K. Held, G. Sangiovanni, P. Hansmann, J. Tomczak] [C. Ederer] [C. Marianetti, A. J. Millis], [E. Dagotto], ...

Oxide Heterostructures : Why at all?



issues and features

- charge transfers and electrostatic boundary conditions
- spin and orbital reconstructions
- compressive and tensile strain
- source/sink for additional (point) defects
- two-dimensional characteristics

prospects

- unique, well-defined doping scenario
- tuning (proximity) options building up on the bulk
- (control of) layer-selective physics
- possibly novel (emergent) electronic phases
- oxide membranes and merging oxides with other materials

Outline

- theory: from DFT to DFT+DMFT
- interfacing band insulator with band insulator
- interfacing band insulator with Mott insulator
- natural-heterostructure systems: delafossites

Density Functional Theory (DFT) : Effective Single-Particle Approach

electronic density functional theory (DFT) with approximate exchange-correlation

- DFT is ground-state many-body theory with inhomogeneous charge density $n(\mathbf{r})$ ($\Leftrightarrow \Psi(\{\mathbf{r}_i \sigma_i\})$) as central quantity, formally exact

$$E[n] = T[n] + \int d\mathbf{r} n(\mathbf{r}) v_{\text{ext}}(\mathbf{r}) + E_{\text{H}}[n] + \mathbf{E}_{\text{xc}}[n]$$

- approximate exchange-correlation functional \mathbf{E}_{xc} in effective single-particle Kohn-Sham mapping to render the approach practicable

local density approximation (LDA) : homogeneous electron gas is reference
generalized-gradient approximation (GGA) : ∇n corrections to LDA

- **works very well for many realistic systems !**
→ many molecules, simple metals, semiconductors (not band gap), . . .



W. Kohn, Nobel Prize 1998

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- **PROBLEM: insufficient for materials with strong correlations among the electrons !**

- ⚡ some systems with odd electron number in unit cell are isolating
- ⚡ satellite structures in one-particle spectral function
- ⚡ strong changes of spectral function with doping
- ⚡ metal-to-insulator transition without symmetry change
- ⚡ coherence temperatures and generally subtle behavior with temperature
- ⋮



W. Kohn, Nobel Prize 1998

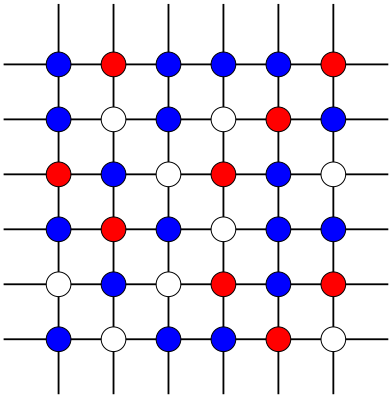
Correlated Electrons : Quasiparticles and Coherence, intuitively



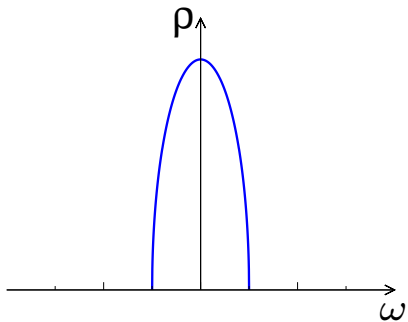
Correlated Electrons : Mott-Hubbard picture

Hubbard model:
$$H = -t \sum_{\langle ij \rangle \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}, \quad \text{for half filling } n = 1$$

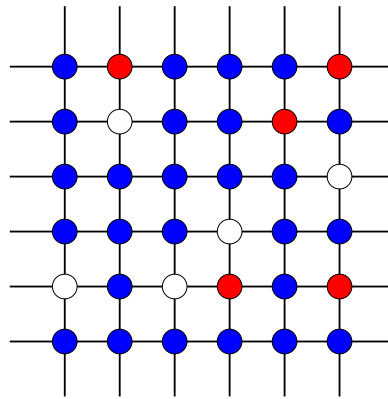
nonint. limit ($U = 0$)



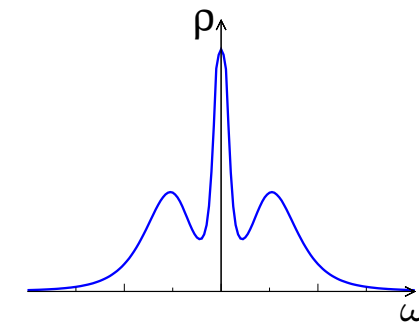
- $H = \sum_{\mathbf{k}\sigma} \varepsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma}$
- ideal metal, Fermi gas



intermediate ($U \sim t$)

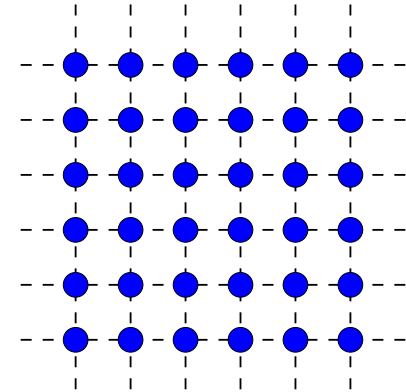


- correlated metal
- dominance of single occupations

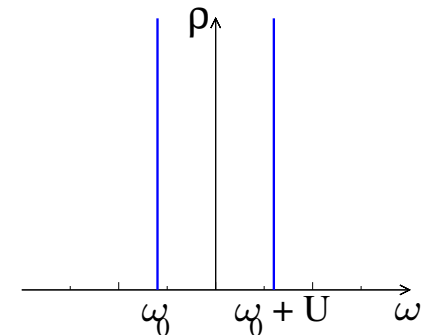


"Three-Peak Structure"

atomic limit ($t = 0$)

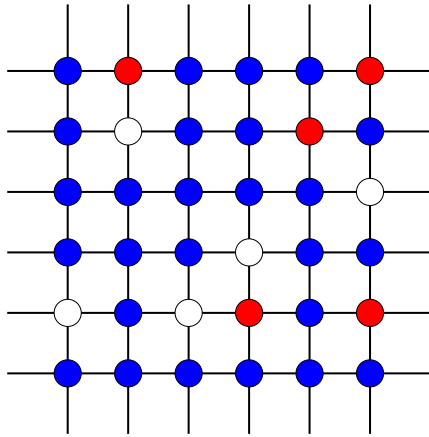


- decoupled lattice sites
- only single occupations in ground state

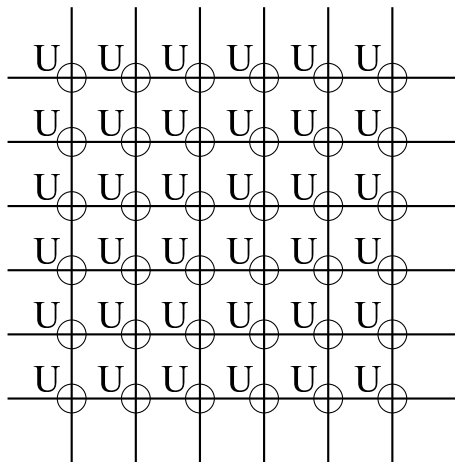
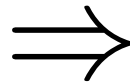


Dynamical Mean-Field Theory (DMFT) : Many-Body Approach

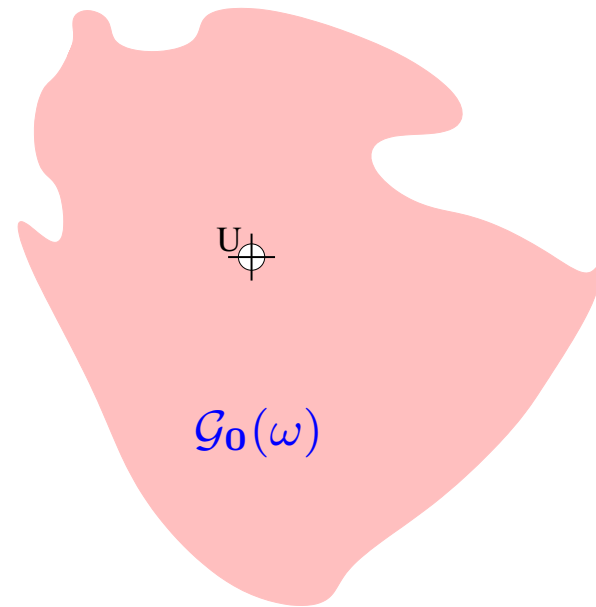
[Metzner, Vollhardt, PRL (1989)] [Georges, Kotliar, PRB (1992)]



$$H = -t \sum_{\langle ij \rangle \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$



$$G_{\text{loc}}(\omega) = \sum_{\mathbf{k}} [\omega + \mu - \varepsilon_{\mathbf{k}} - \Sigma(\mathbf{k}, \omega)]^{-1}$$



$$G_{\text{imp}}(\omega) = \left[\mathcal{G}_0^{-1}(\omega) - \Sigma_{\text{imp}}(\omega) \right]^{-1}$$

DMFT approximation:

$$G_{\text{loc}}(\omega) \stackrel{!}{=} G_{\text{imp}}(\omega)$$

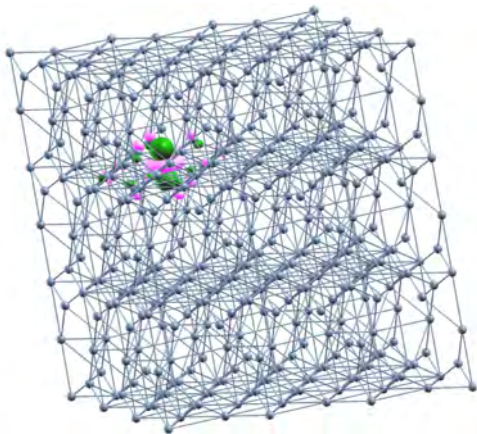
$$G_{\text{loc}}^{\text{DMFT}}(\omega) = \sum_{\mathbf{k}} [\omega + \mu - \varepsilon_{\mathbf{k}} - \Sigma_{\text{imp}}(\omega)]^{-1}$$

→ local correlations are fully taken into account non-pertubatively, explicit non-local correlations are neglected

Correlated Materials : different levels of correlation

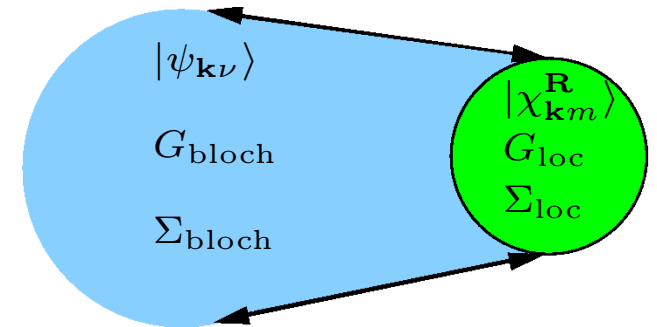
density functional theory (DFT)

- Hohenberg-Kohn theorems, Kohn-Sham construction, exchange-correlation functional, . . .
- mapping interacting electrons onto non-interacting electrons in an complicated effective potential
- very good description of the realistic, single-particle based chemical bonding (\rightarrow band structure for solids)

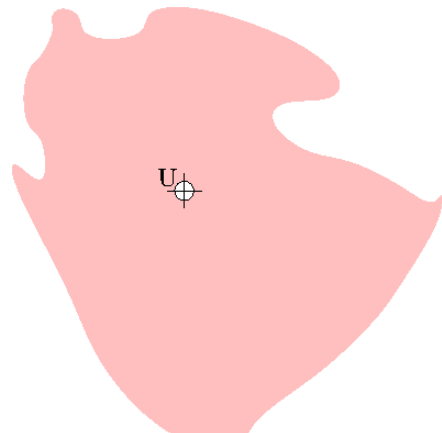
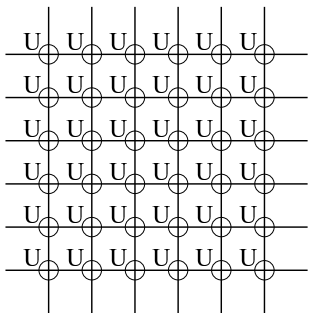


interface

- define correlated subspace by Wannier(-like) orbitals from correlated sites



dynamical mean-field theory (DMFT)



- many-body Green's functions, Hubbard model, quasiparticles vs. Hubbard bands, finite temperature, . . .
- mapping interacting electrons onto an interacting impurity in a self-consistent bath
- very good description of modeled interacting electrons on a lattice

Realistic Many-Body

[Anisimov, Poteryaev, Korotin et al., JPCM 9, 7359 (1997)]

[Held, Keller, Eyert et al., PRL 86, 5345 (2001)]

[Minár, Chioncel, Perlov, et al., PRB 72, 045125 (2005)]

[FL, Georges, Poteryaev et al., PRB 74, 125120 (2006)]

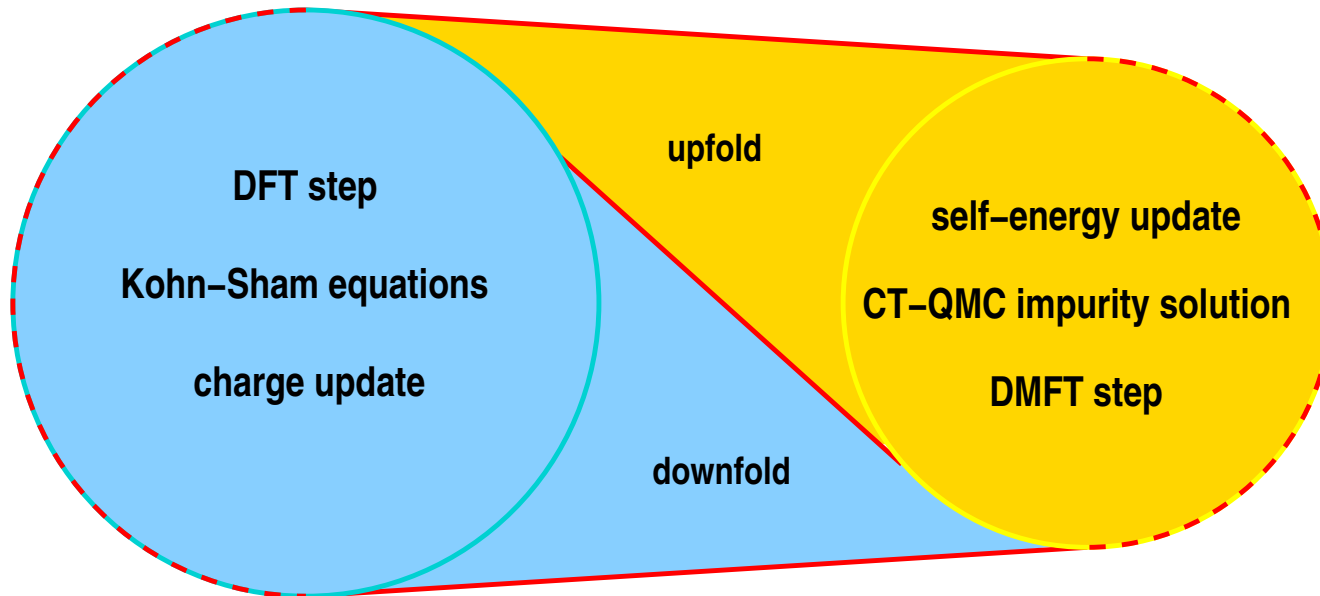
[Lichtenstein and Katsnelson, PRB 57, 6884 (1998)]

[Pavarini, Biermann, Poteryaev et al., PRL 92, 176403 (2004)]

[Kotliar, S. Y. Savrasov, K. Haule et al, RMP 78, 865 (2006)]

[Grieger, Piefke, Peil and FL, PRB 86, 155121 (2012)] ..., ..., ...

Density Functional Theory (DFT) + Dynamical Mean-Field Theory (DMFT)



(Daniel Grieger, Lewin Boehnke, Christoph Piefke, FL)

● DFT part : mixed-basis pseudopotential

[Meyer, Elsässer, Lechermann and Fähnle, MPI for Metals Research]

● DMFT impurity solver : hyb-CT-QMC (TRIQS code)

[Parcollet et al., Comput. Phys. Commun. 196, 398 (2015)]

[Seth et al, Comput. Phys. Commun. 200, 274 (2016)]

● charge self-consistency

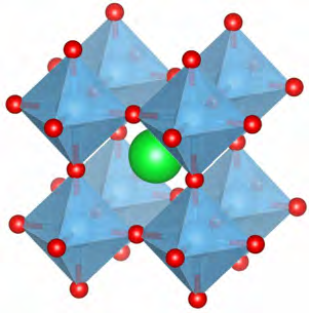
● 1-5 local correlated orbitals

● up to 200-atom unit cells

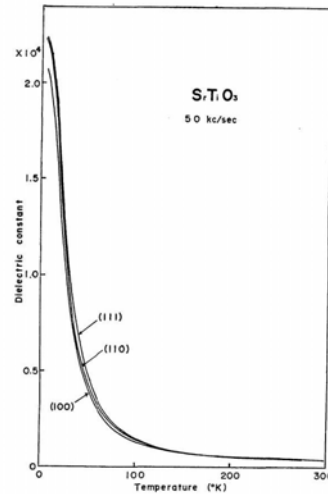
● OX-HE: manifest multi-site DMFT



Bulk SrTiO₃ : A Band Insulator prone to a Plethora of Physics

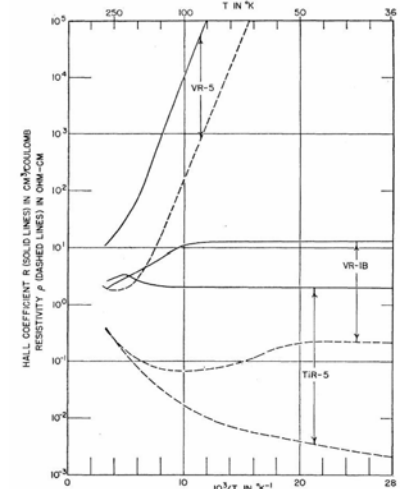


- cubic perovskite, transition to tetragonal phase at low temperature
- band insulator ($\Delta_g \sim 3.2$ eV)



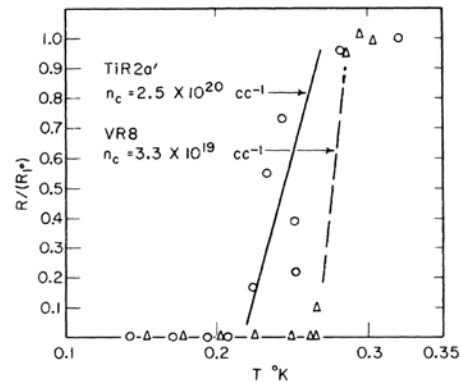
paraelectric

[Sawaguchi et al., JPSN 17, 1666 (1962)]



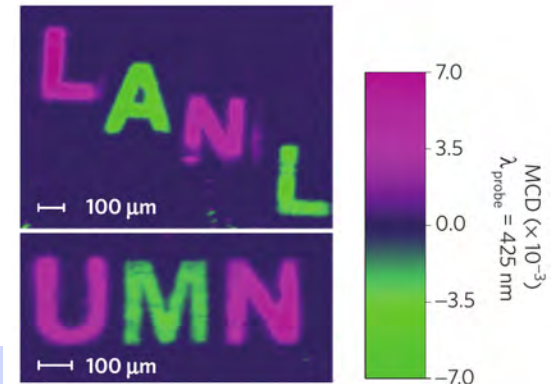
semiconductor

[Frederikse et al., PR 134, A442 (1964)]



superconductor

[Schooley et al., PRL 12, 474 (1964)]



magnet

[Rice et al. Nat. Mater. 13, 481 (2014)]

Emergent Physics at the $\text{LaAlO}_3/\text{SrTiO}_3$ Interface

LETTERS

PUBLISHED ONLINE: 4 SEPTEMBER 2011 | DOI: 10.1038/NPHYS2080

nature
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Coexistence of magnetic order and two-dimensional superconductivity at $\text{LaAlO}_3/\text{SrTiO}_3$ interfaces

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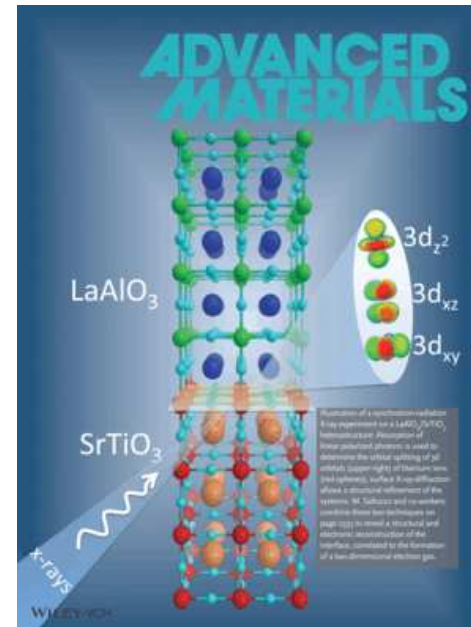
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Titanium d_{xy} ferromagnetism at the $\text{LaAlO}_3/\text{SrTiO}_3$ interface

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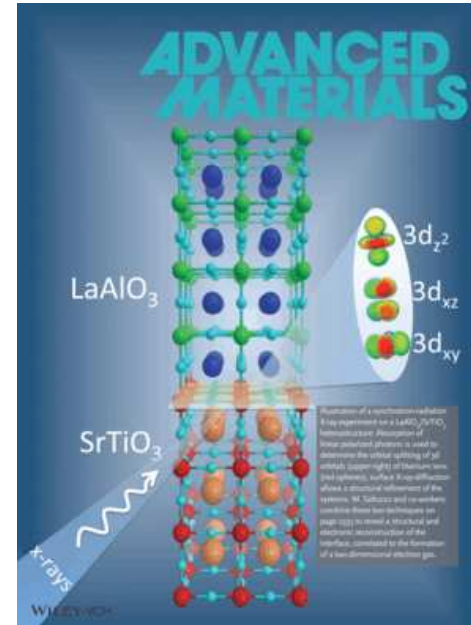
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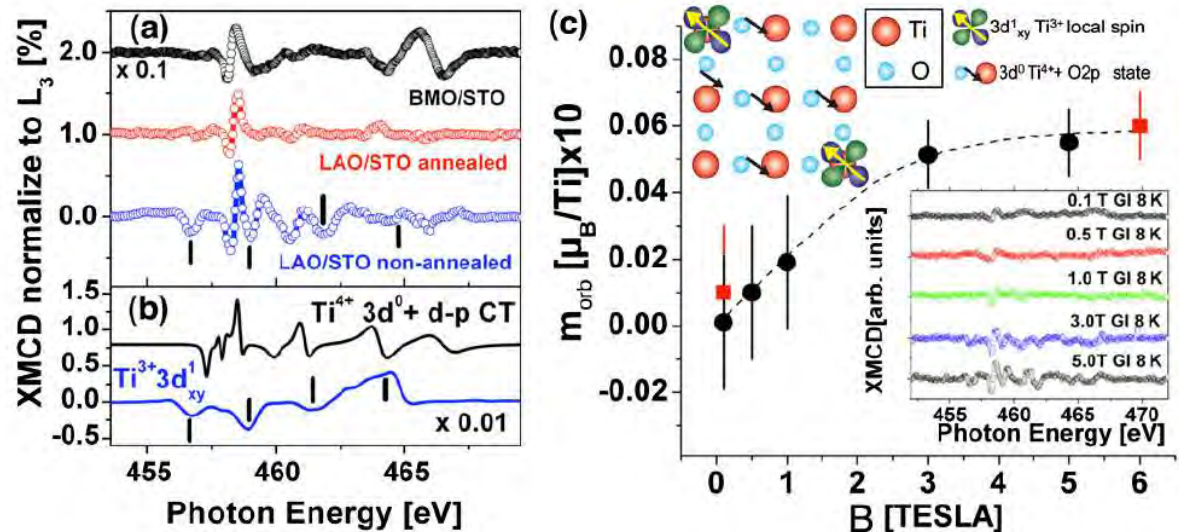
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RIXS and XMCD measurements
relevance of oxygen vacancies
at the interface

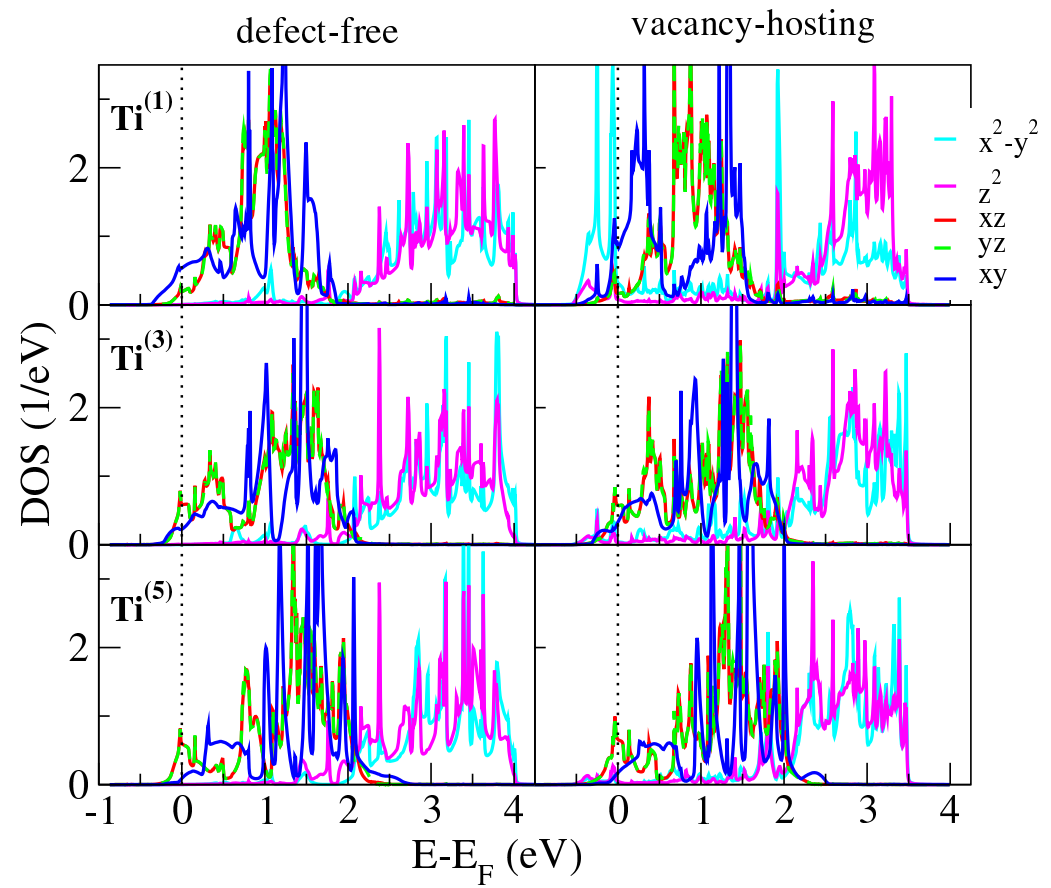
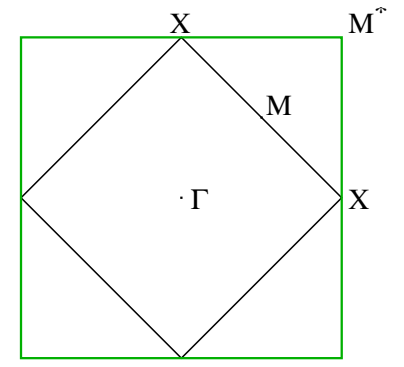
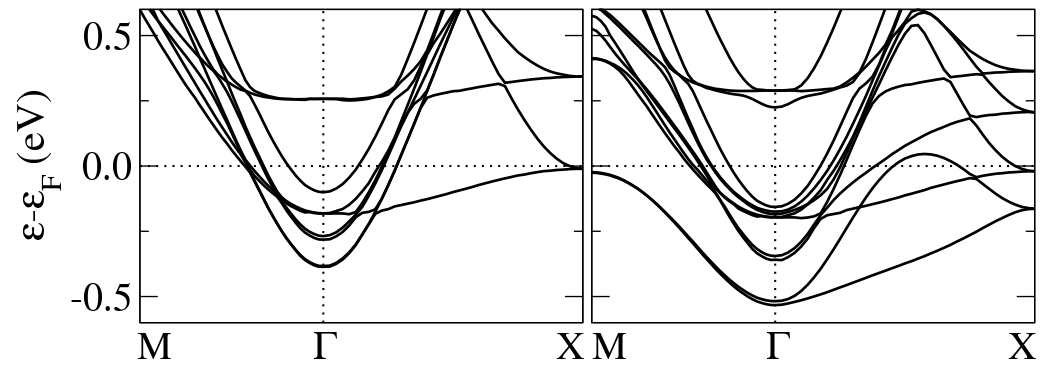
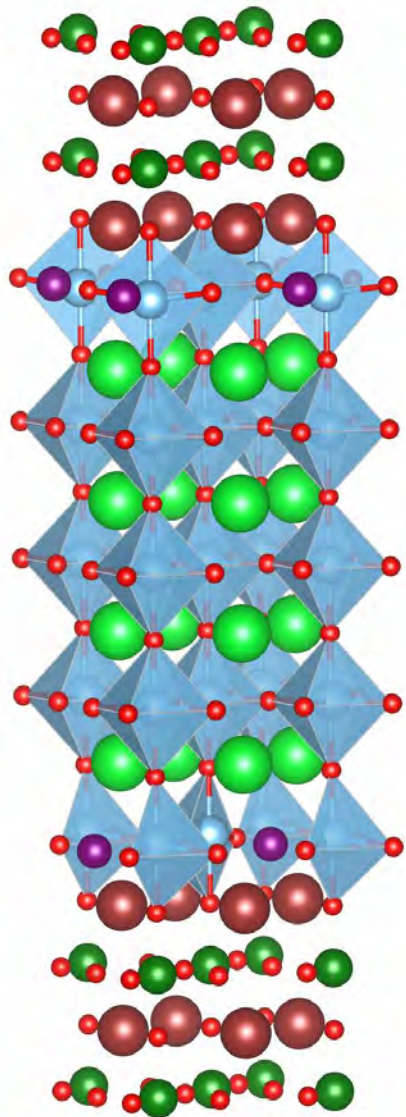
[Park et al., PRL 110, 017401 (2013)]

[Salluzzo et al., PRL 111, 087204 (2013)]



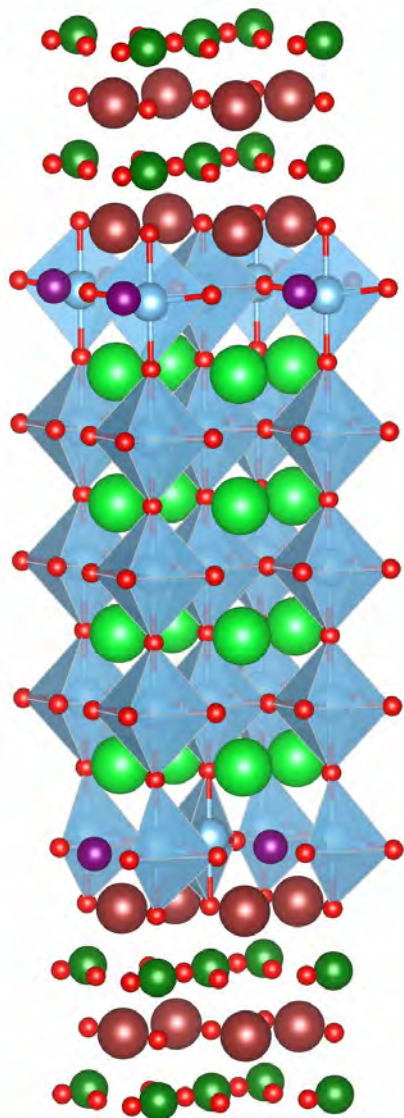
oxygen vacancies introduce electrons that localize at the nearby Ti sites ($m \sim 0.1 - 0.2\mu_B$)

LaAlO₃/SrTiO₃ interface: DFT results



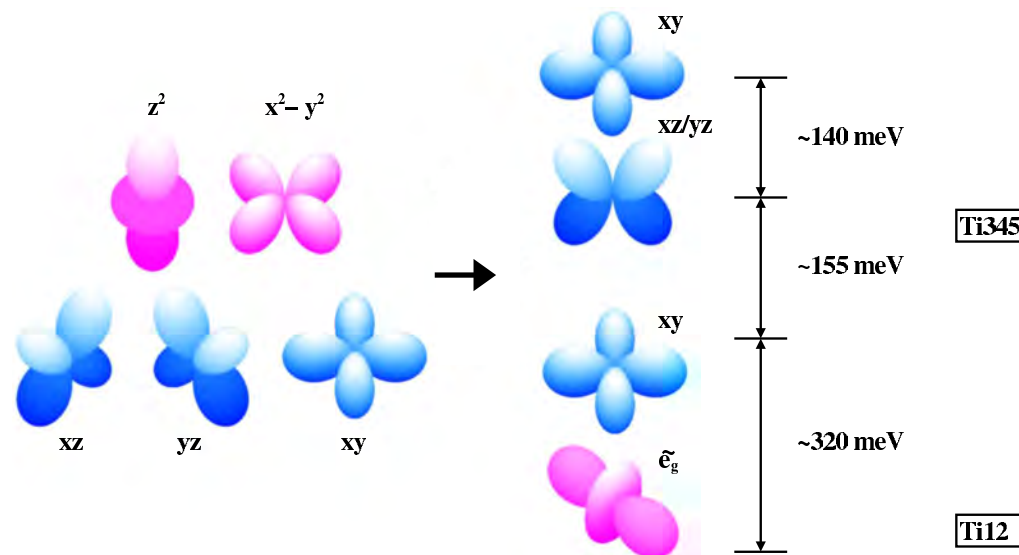
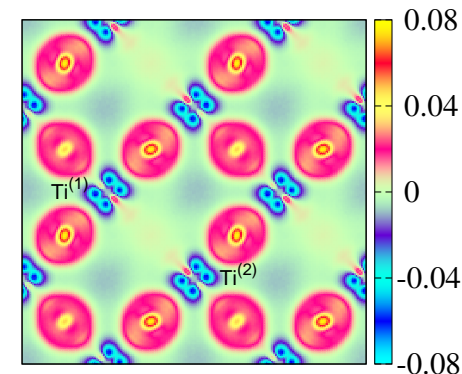
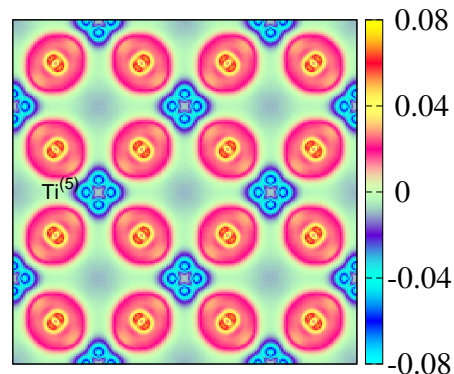
see also [Pavlenko et al., PRB 85, 020407(R) (2012); PRB 88, 201104(R) (2013)]

LAO/STO dense-defect interface : Relevant States



in-plane bond
charge density

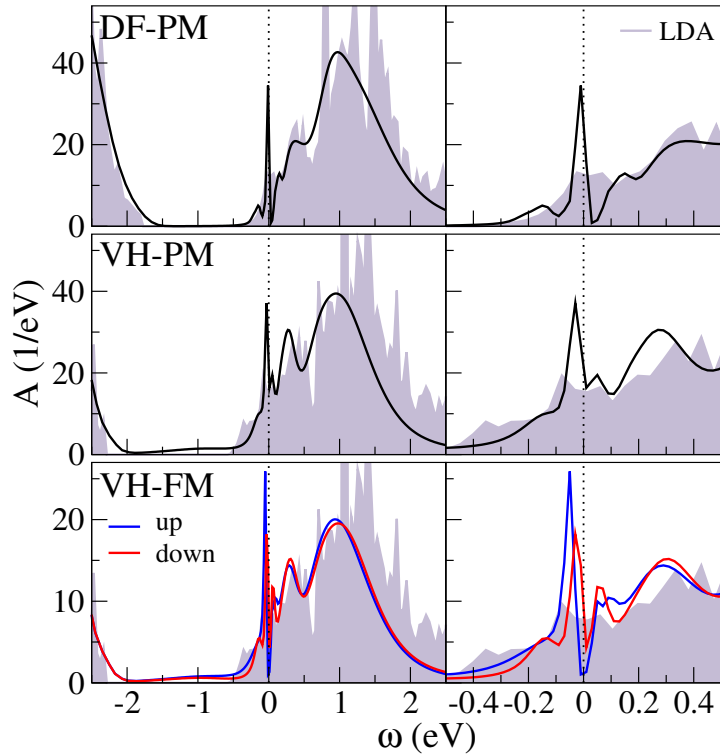
$$\rho - \rho_{at}$$



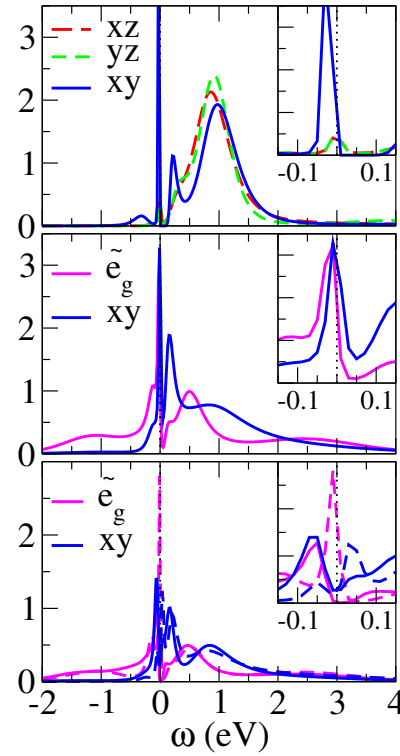
- two-orbital correlated subspace for each Ti ion from orbital projection
- interface Ti: $\tilde{e}_g - t_{2g}(xy)$ description, with $|\tilde{e}_g\rangle \sim 0.55|z^2\rangle \pm 0.84|x^2 - y^2\rangle$
- \Rightarrow **coexisting localized and itinerant electrons expected !**
- \rightarrow in line with RIXS measurements [Zhou et al., PRB 83, 201402(R) (2011)]

LAO/STO dense-defect interface : DFT+DMFT

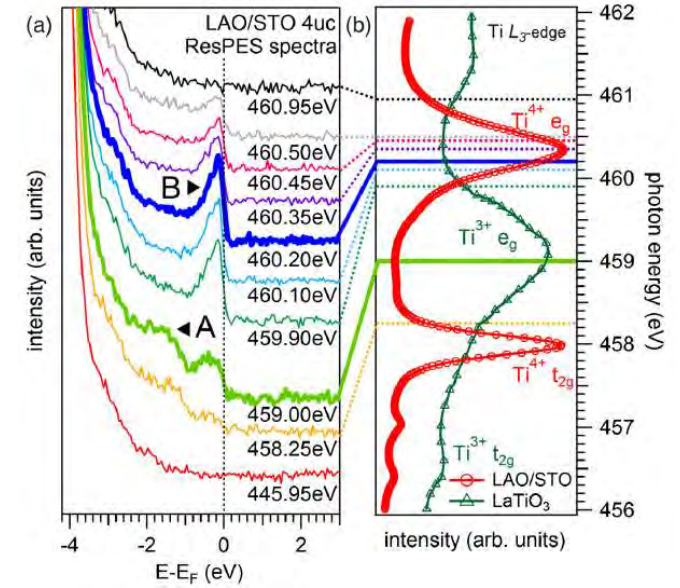
total spectrum



interface Ti spectrum



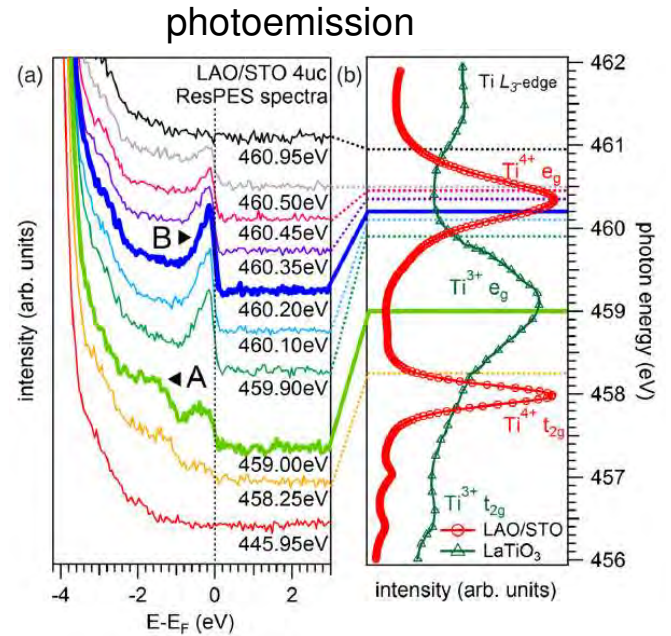
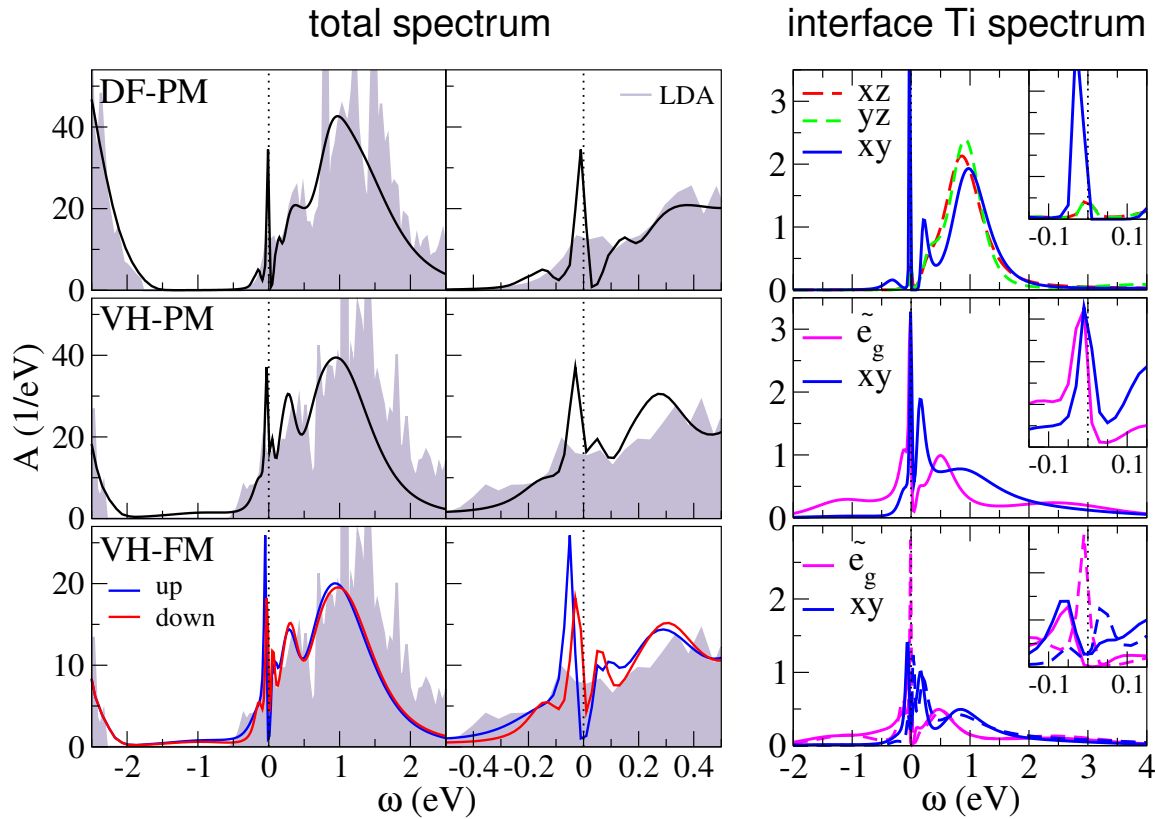
photoemission



[Cancellieri et al., PRL 110 137601 (2013)]

[Berner et al., PRL 110, 247601 (2013)]

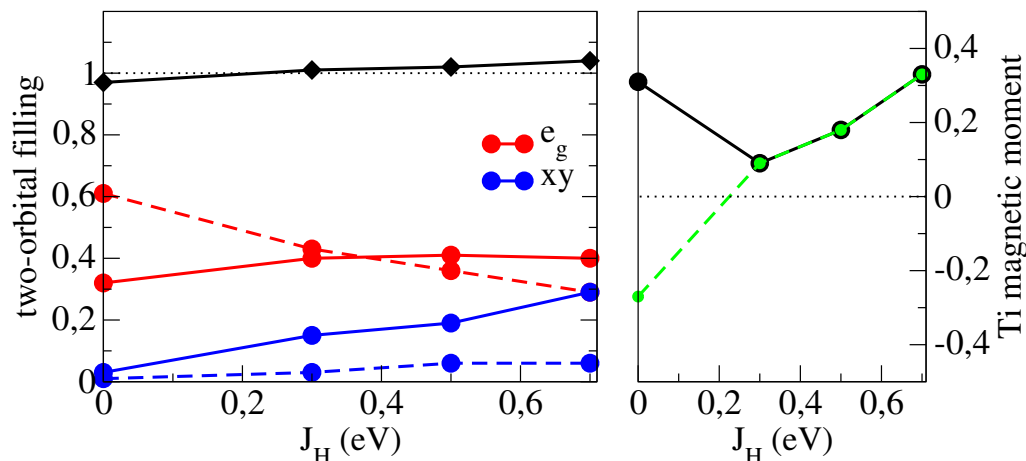
LAO/STO dense-defect interface : DFT+DMFT



[Cancellieri et al., PRL 110 137601 (2013)]

[Berner et al., PRL 110, 247601 (2013)]

local two-orbital fillings and polarizations with vacancies ($U=2.5$ eV)



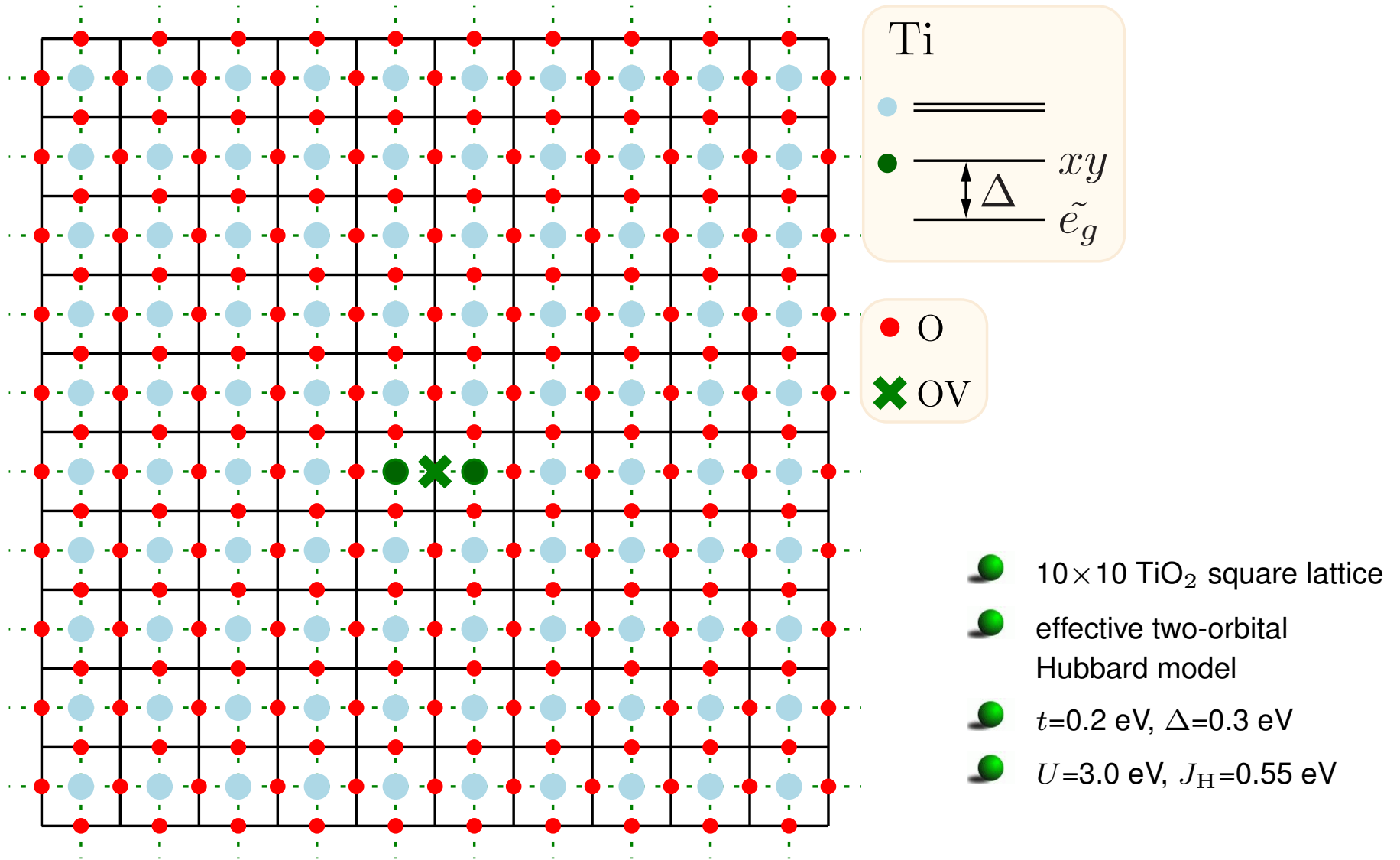
with oxygen vacancies (VH)

- one electron at interface Ti: quarter filling
- FM moment $m_{Ti12} \sim 0.2\mu_B$
- double-exchange mechanism
- **very sensitive to J_H !**

[FL, Boehnke, Grieger, Piefke, PRB 085125 (2014)]

Arbitrary OV Concentration in LAO/STO : Real-Space Modelling

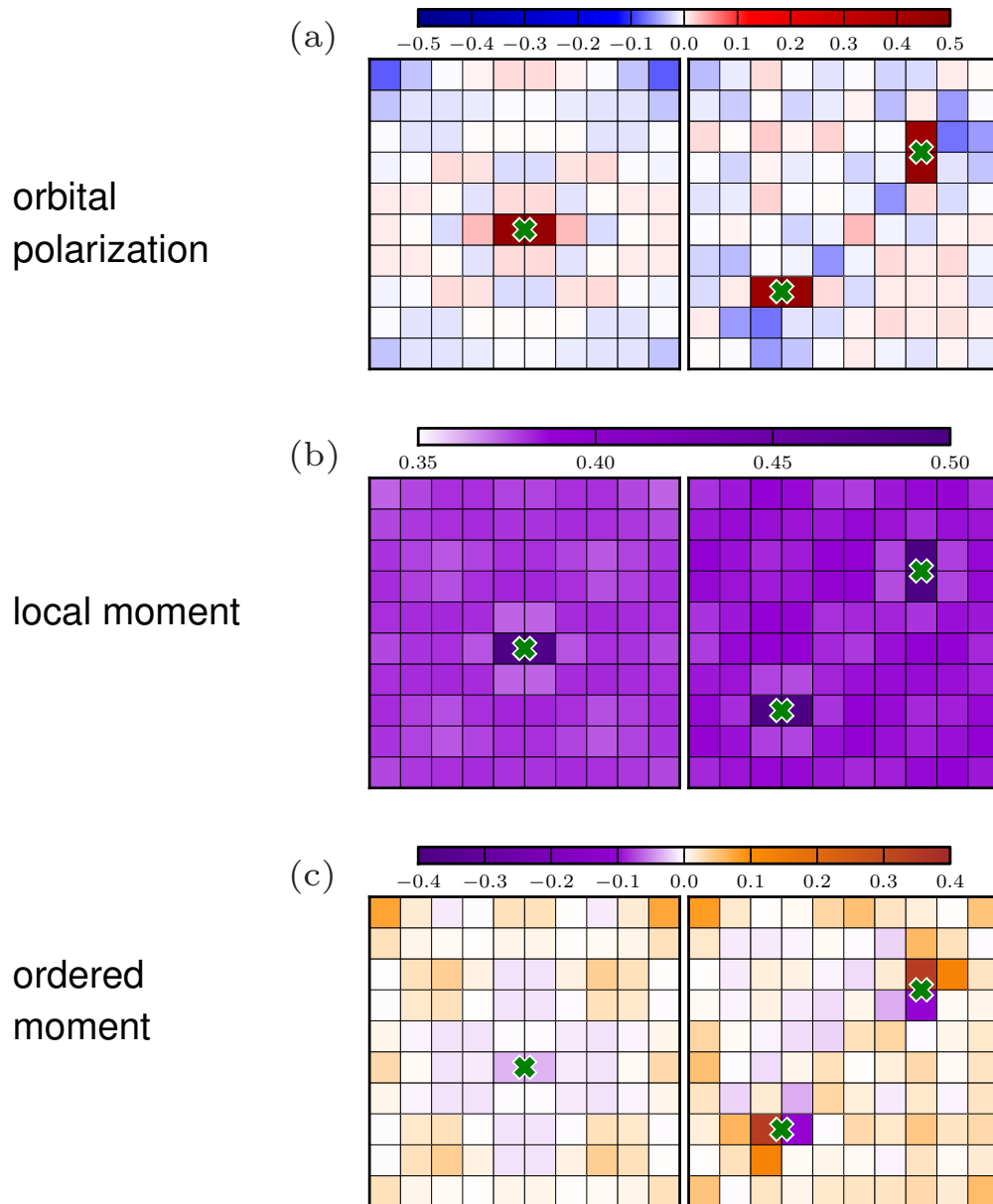
[Behrmann and FL, PRB 92, 125148 (2015)]



● Gutzwiller/Slave-Boson approach : nearest-neighbor hopping and crystal field from former DFT calculation

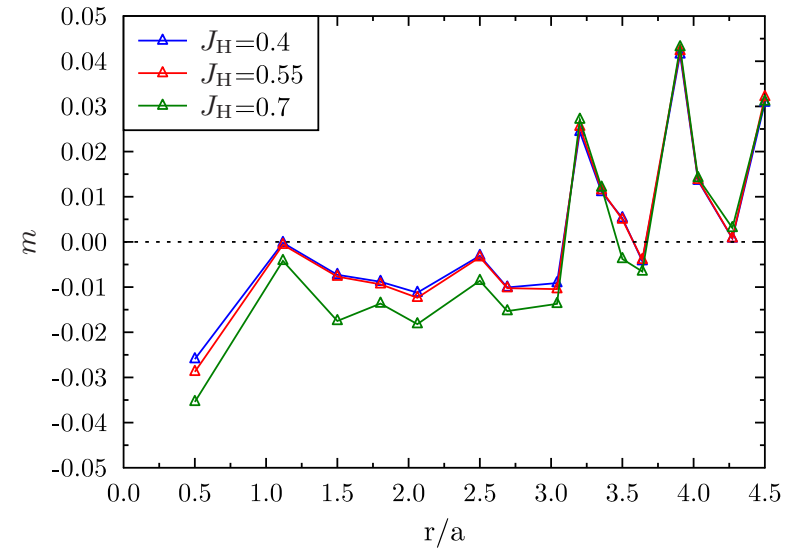
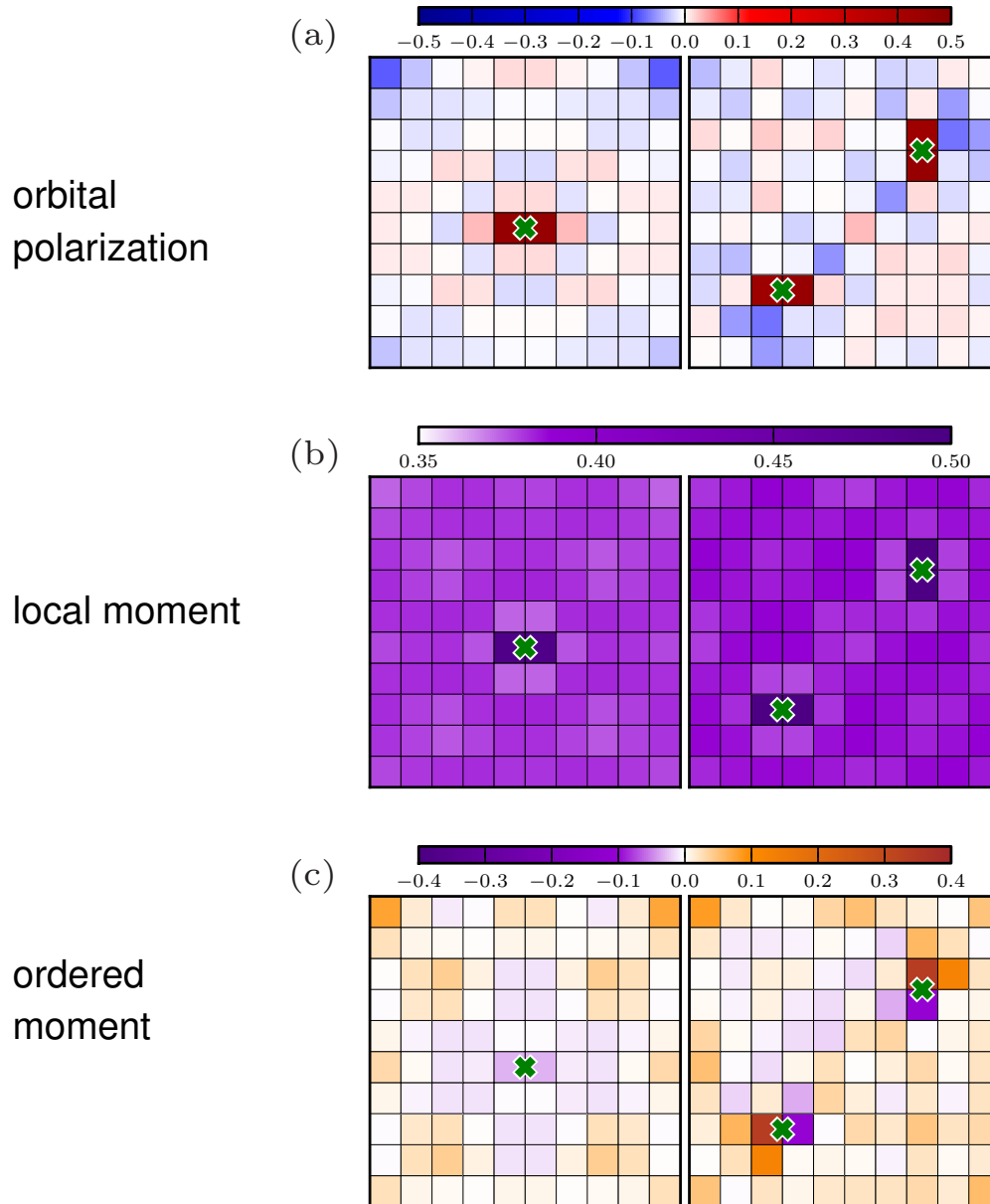
Arbitrary OV Concentration in LAO/STO : Dilute-Defect Limit

[Behrmann and FL, PRB 92, 125148 (2015)]



Arbitrary OV Concentration in LAO/STO : Dilute-Defect Limit

[Behrmann and FL, PRB 92, 125148 (2015)]



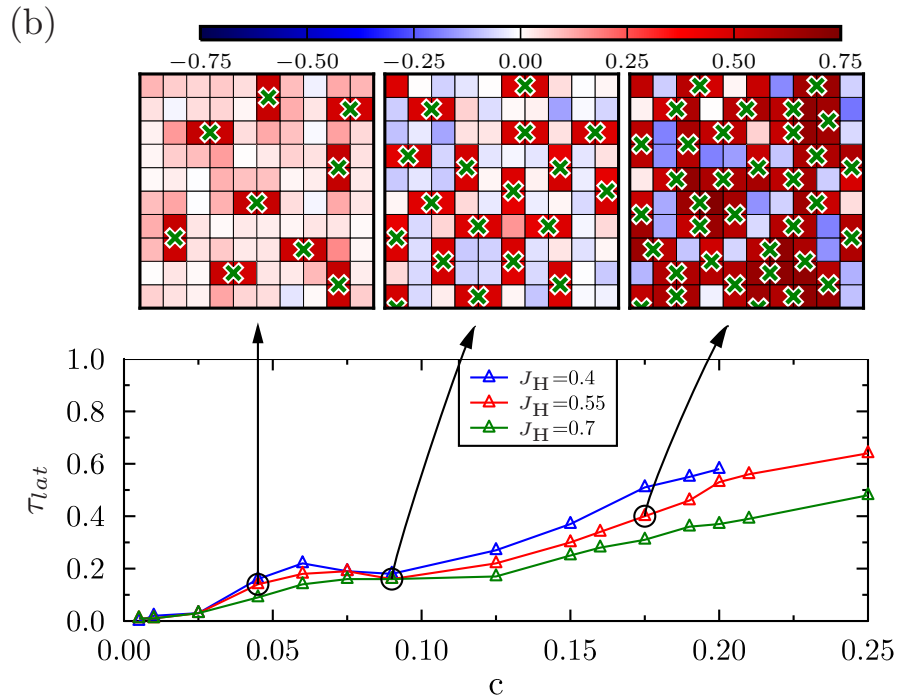
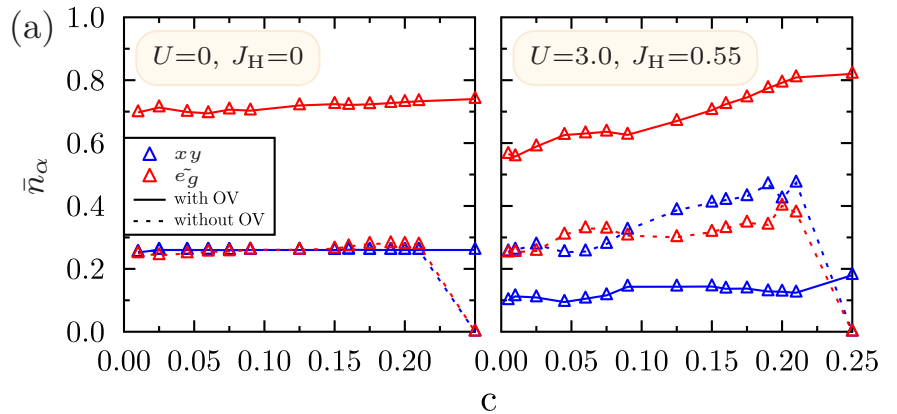
RKKY(-like) coupling

$$J(d) \cos(2k_F d) \quad , \quad d \sim 2.5 \Rightarrow k_F \sim \frac{\pi}{10}$$

→ about 1/8 filling ✓

Arbitrary OV Concentration in LAO/STO : orbital and spin polarization

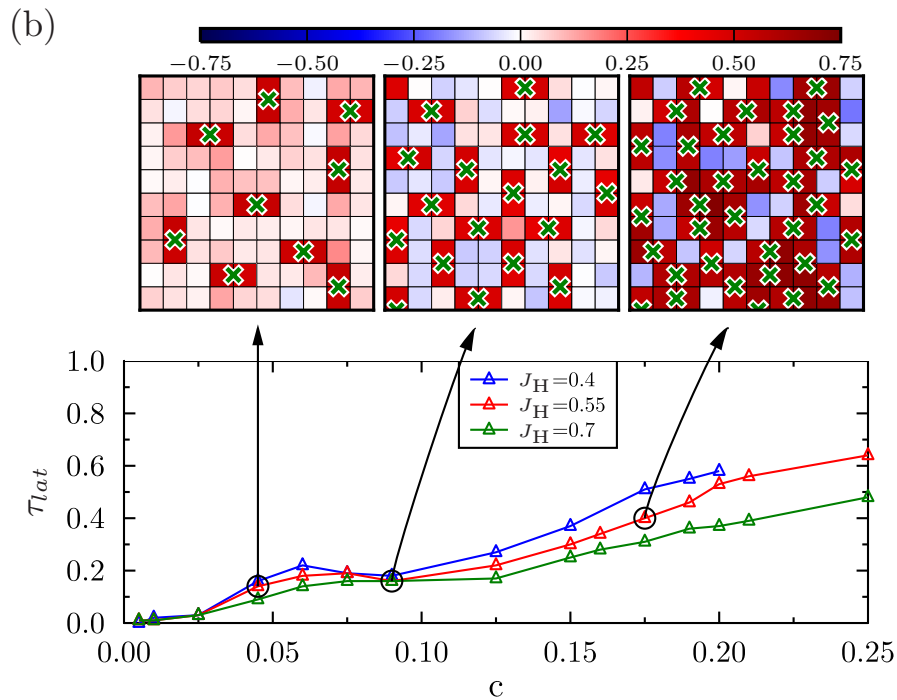
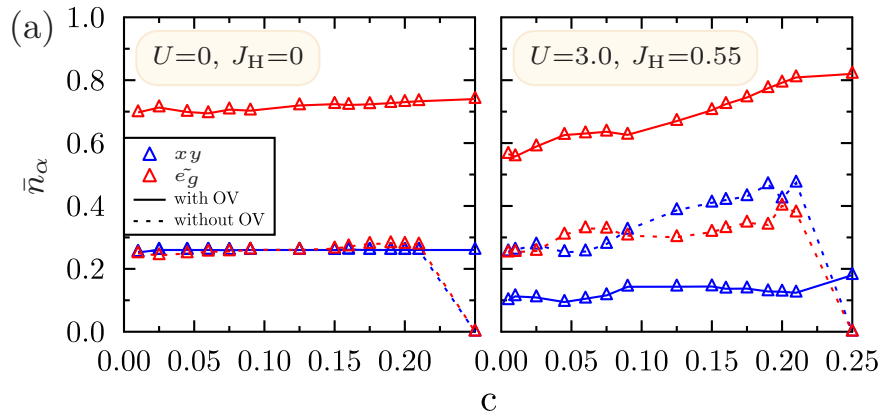
orbital occupation and orbital moment



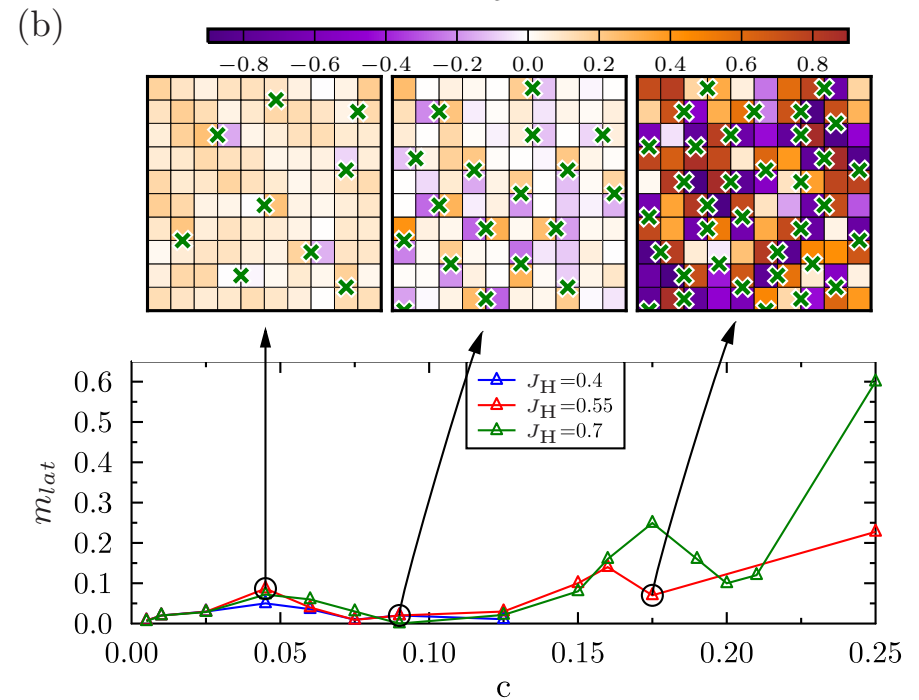
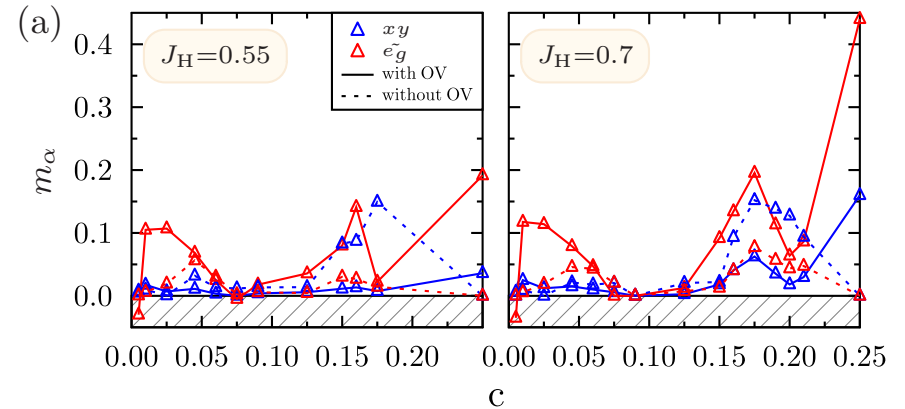
 intricate dichotomy between Ti sites with and without nearby oxygen vacancies


Arbitrary OV Concentration in LAO/STO : orbital and spin polarization

orbital occupation and orbital moment



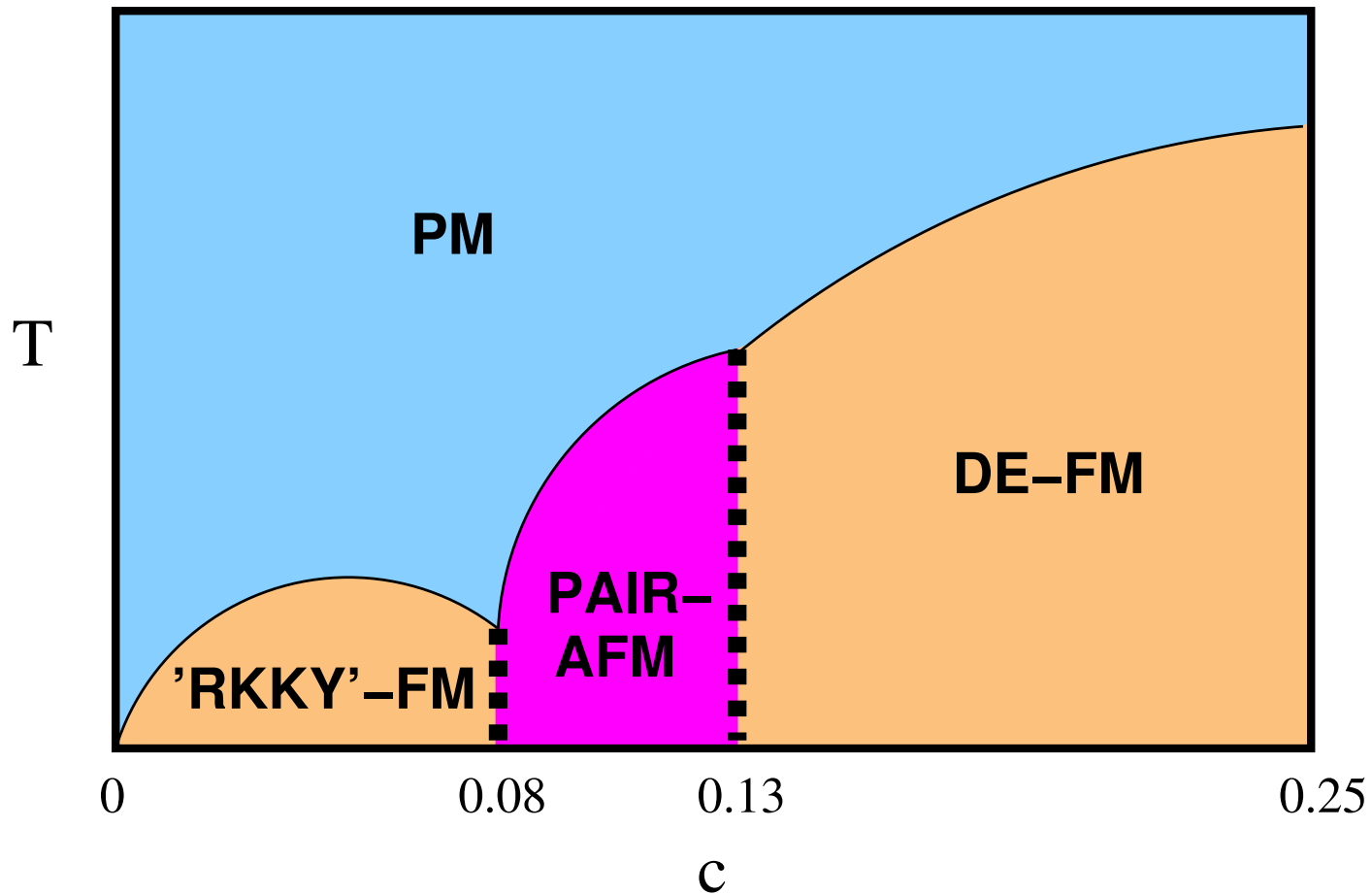
magnetic moment



 intricate dichotomy between Ti sites with and without nearby oxygen vacancies

Arbitrary OV Concentration in LAO/STO : Magnetic Phase Diagram

[Behrmann and FL, PRB 92, 125148 (2015)]

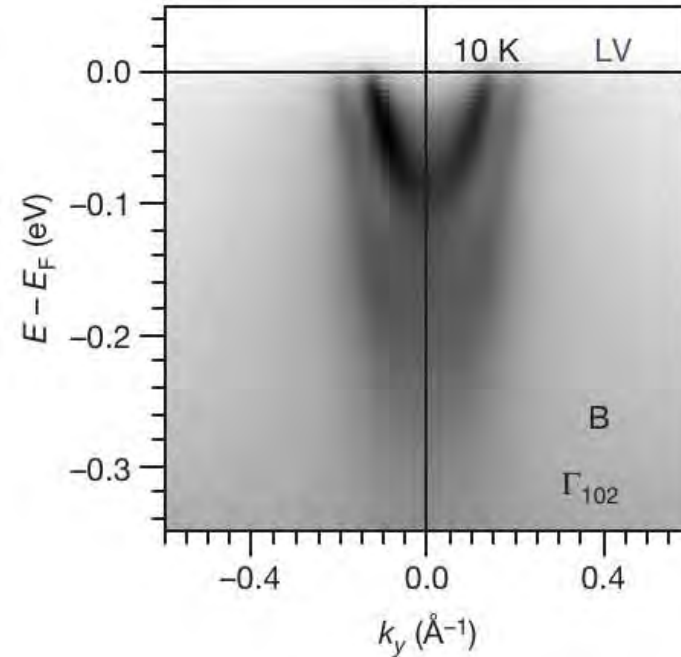
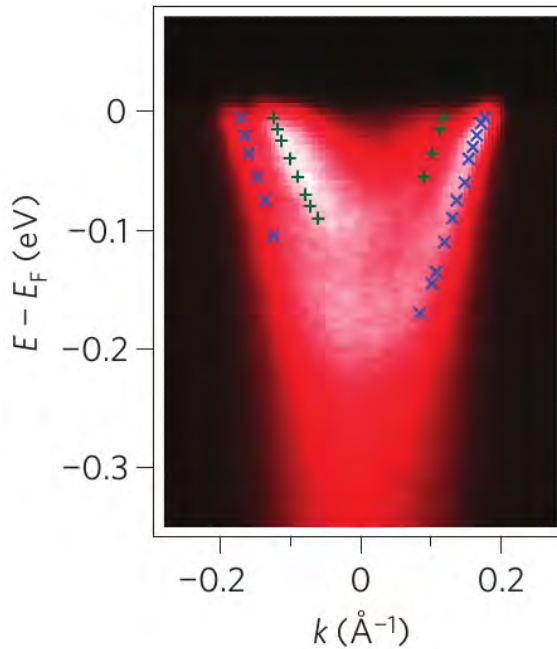
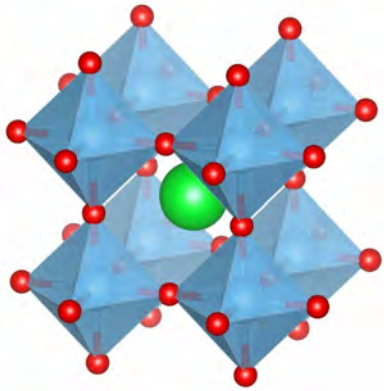


- nontrivial ferrromagnetism triggered by even few oxygen vacancies
- double-exchange-like ferromagnetism for higher oxygen vacancy concentration

And now
for something
completely different...



Oxygen-Deficient SrTiO₃ surface : Experiment



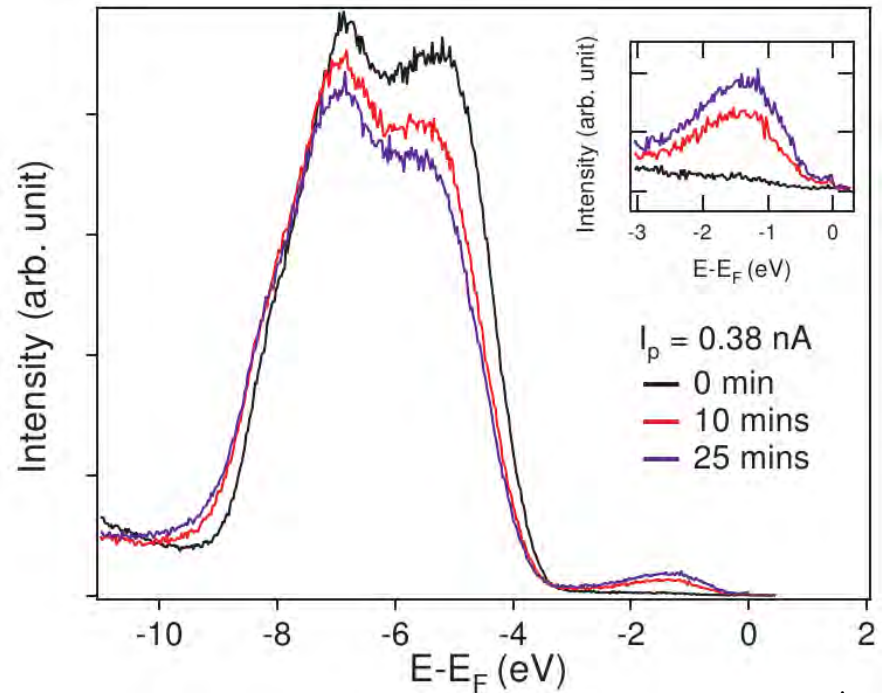
two-dimensional electron system on surface via

 exposure to UV light

[Meevasana et al., Nature Mat. 10, 114 (2011)]

 in-situ crystal fracture

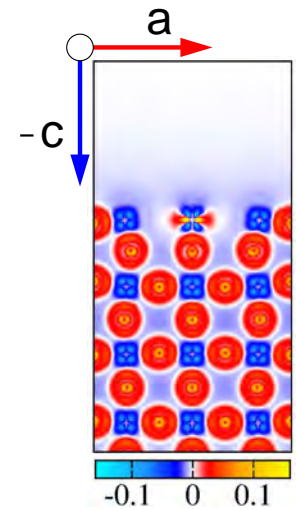
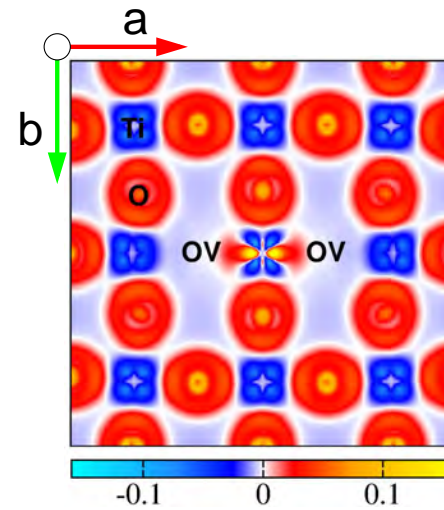
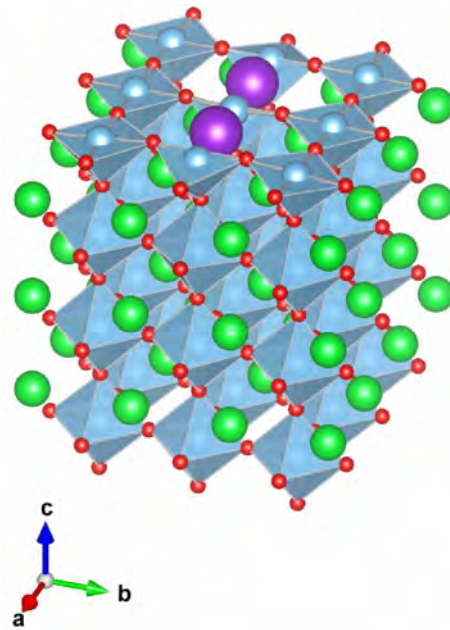
[Santander-Syro et al., Nature 469, 189 (2011)]



Oxygen-Deficient SrTiO₃ surface : DFT+DMFT

[FL, Jeschke, Kim, Backes and Valentí, PRB 93, 121103(R) (2016)]

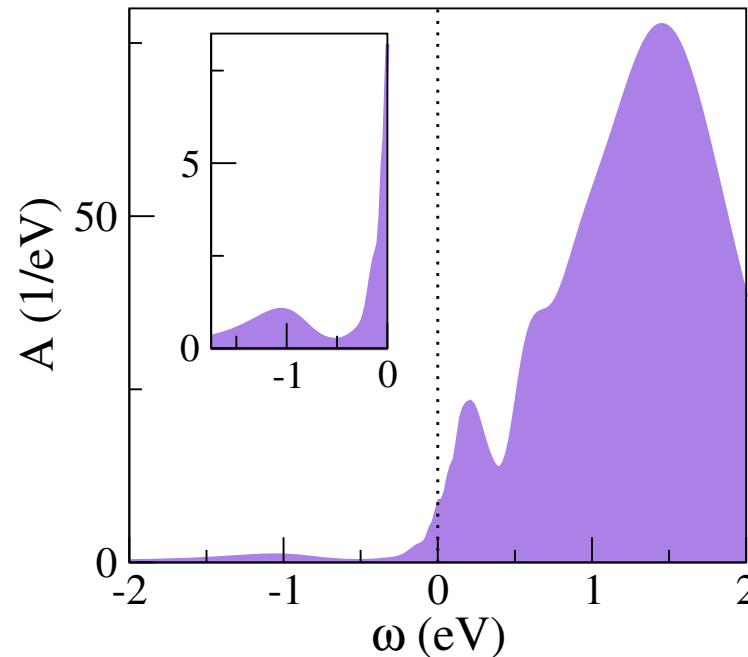
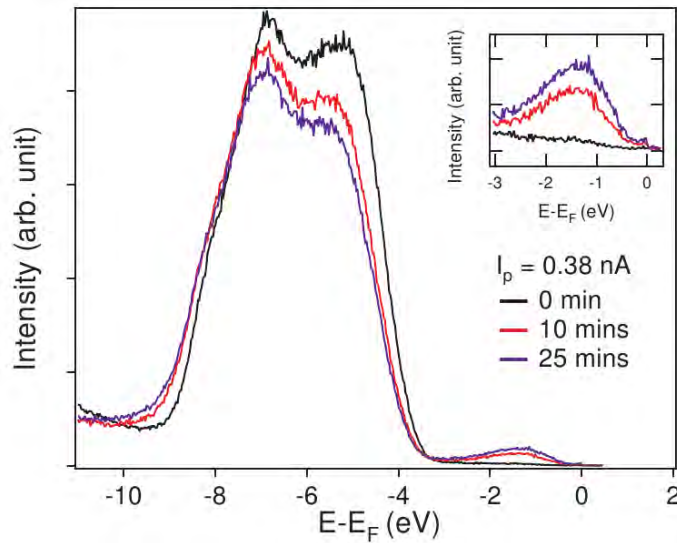
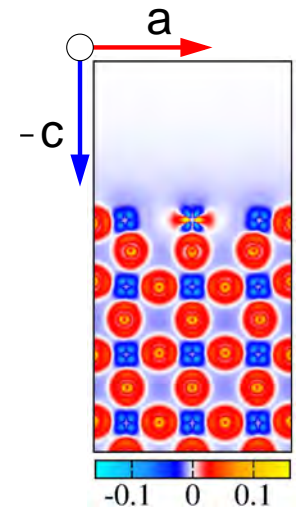
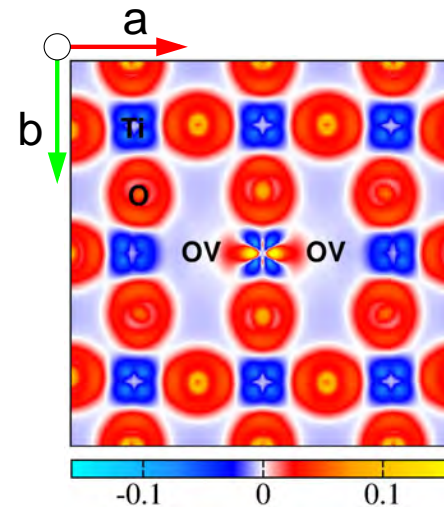
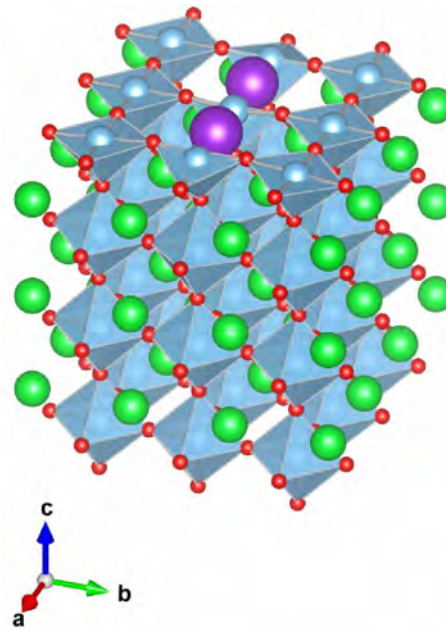
- ★ 180-atom slab
- ★ two oxygen vacancies in TiO₂ surface layer
- ★ projections allow for e_g and t_{2g} states
- ★ $U=3.5\text{eV}$, $J_H=0.5\text{eV}$



Oxygen-Deficient SrTiO₃ surface : DFT+DMFT

[FL, Jeschke, Kim, Backes and Valentí, PRB 93, 121103(R) (2016)]

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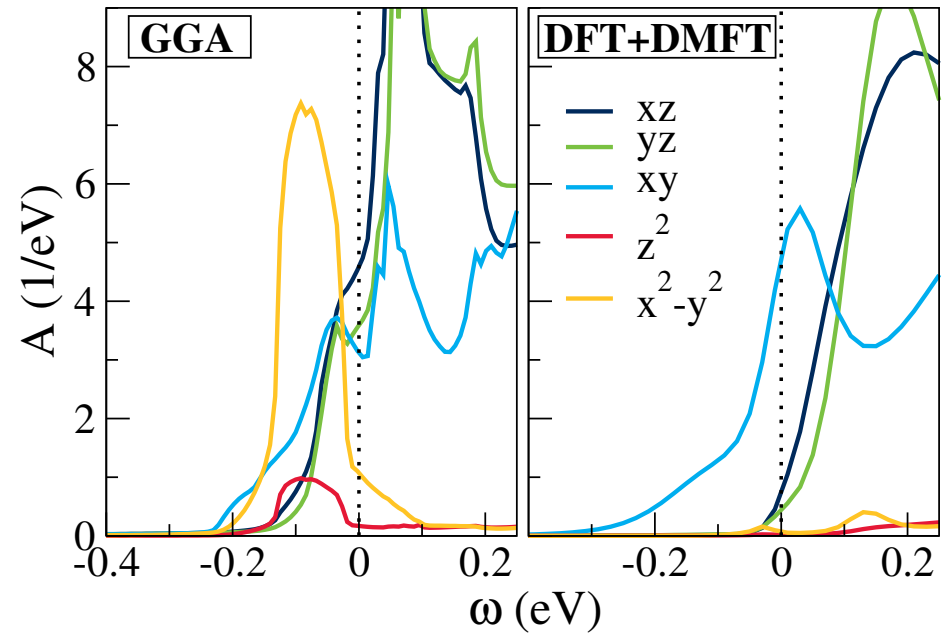
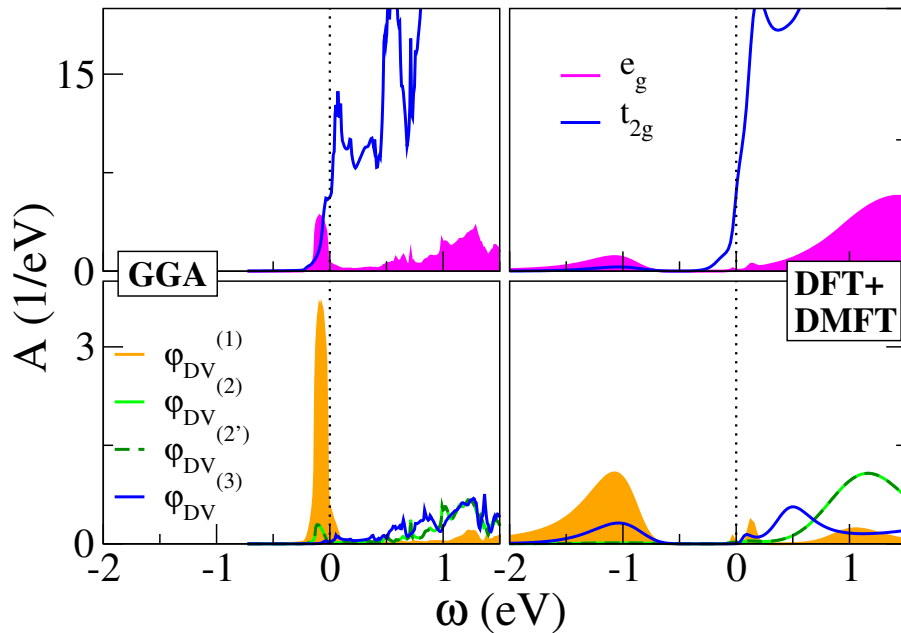
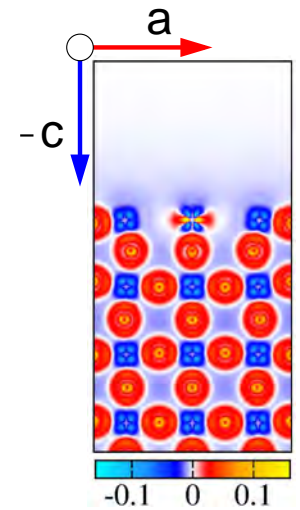
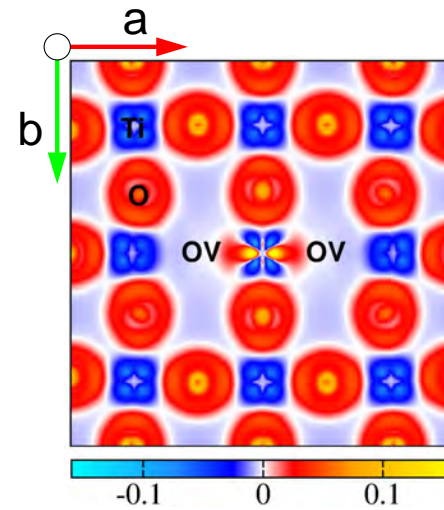
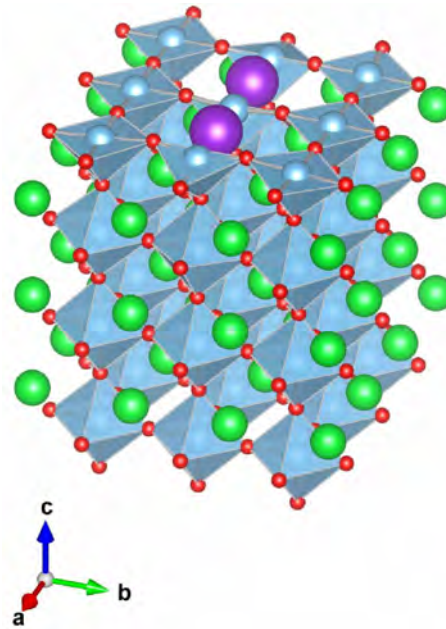


[Meevasana et al., Nature Mat. 10, 114 (2011)]

Oxygen-Deficient SrTiO₃ surface : DFT+DMFT

[FL, Jeschke, Kim, Backes and Valentí, PRB 93, 121103(R) (2016)]

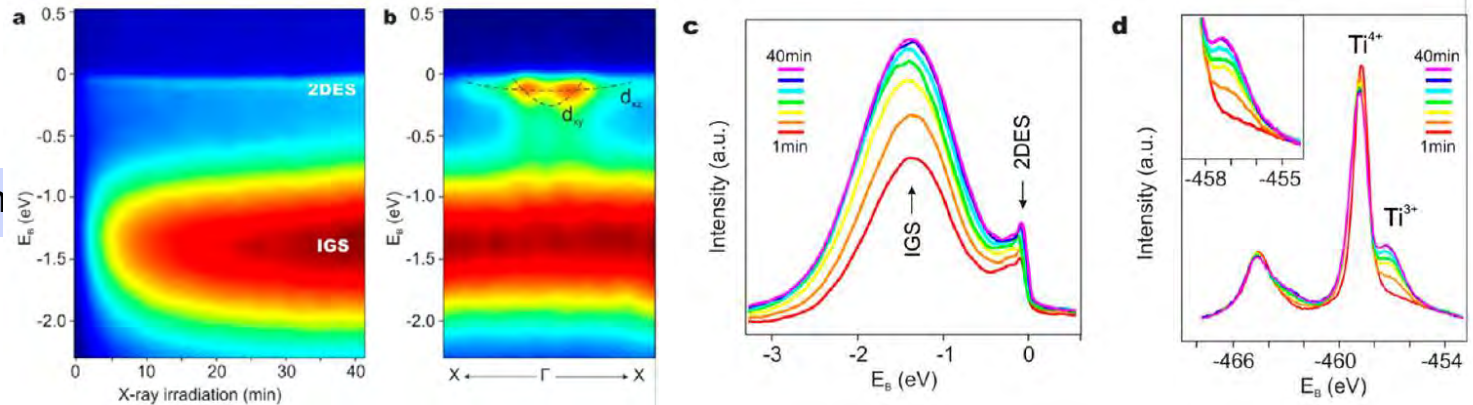
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- ★ $U=3.5\text{eV}$, $J_H=0.5\text{eV}$





LaAlO₃/SrTiO₃ revisited : experiment and DFT+DMFT

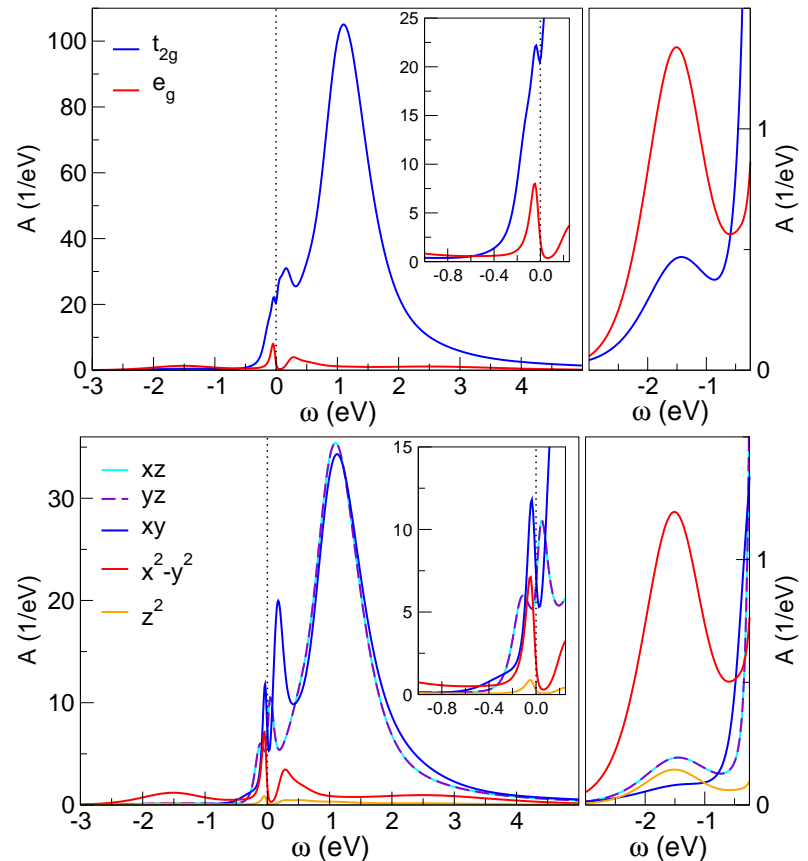
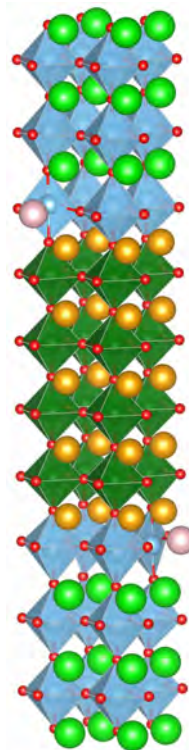
[FL, Boehnke, Grieger and Piefke, PRB 90, 085125 (2014)] [Chikina, FL, Husanu et al., ACS Nano 12, 7927 (2018)]

experimental spectrum



theoretical description

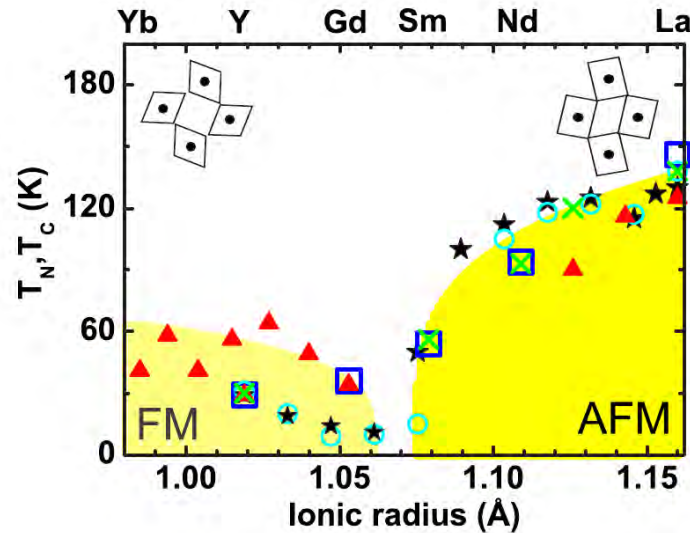
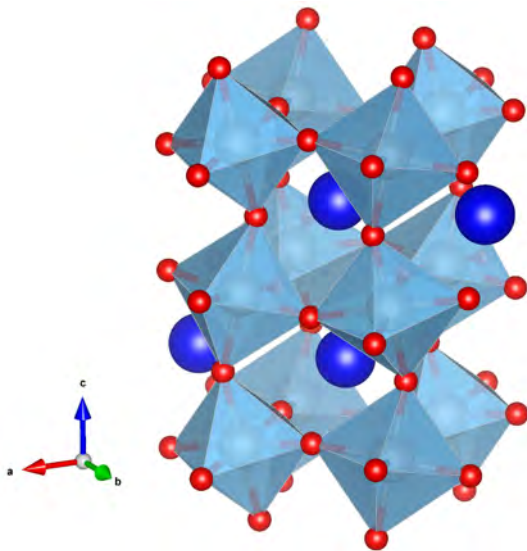
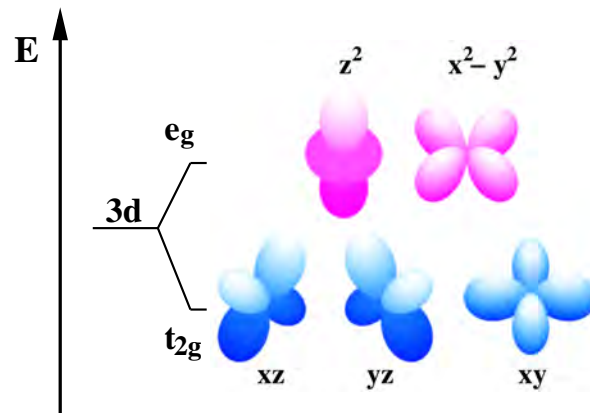
-  200-atom supercell
-  Hubbard $U = 3$ eV,
Hund's exchange $J_H = 0.5$ eV



Doping the Mott-Insulating State : Phenomenology of Bulk Titanates

Mott-insulating distorted perovskites $RTiO_3$

→ Ti^{3+} with $3d^1(t_{2g})$



[Komarek et al., PRB 75, 224402 (2007)]

$LaTiO_3$

AFM $T_N=146K$, apical tilt $\Theta=13^\circ$, charge gap $\Delta=0.20$ eV

$GdTiO_3$

FM $T_C=36K$, apical tilt $\Theta=18^\circ$, charge gap $\Delta=0.75$ eV

$SmTiO_3$

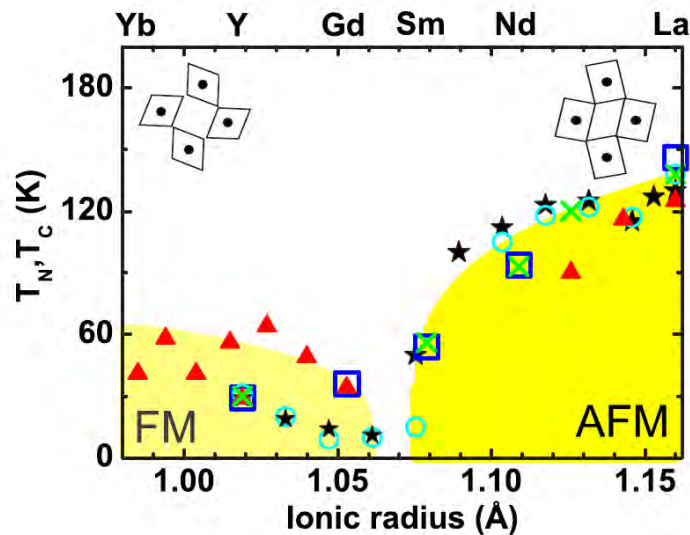
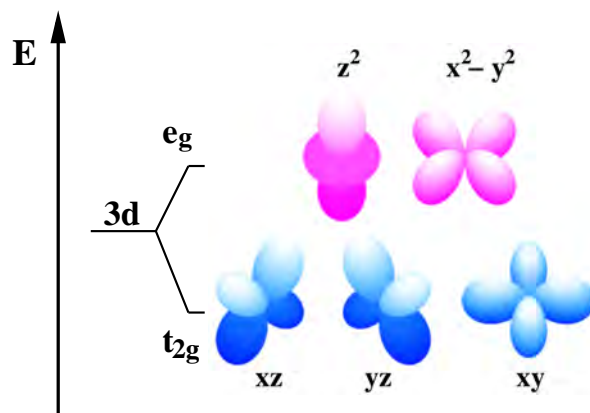
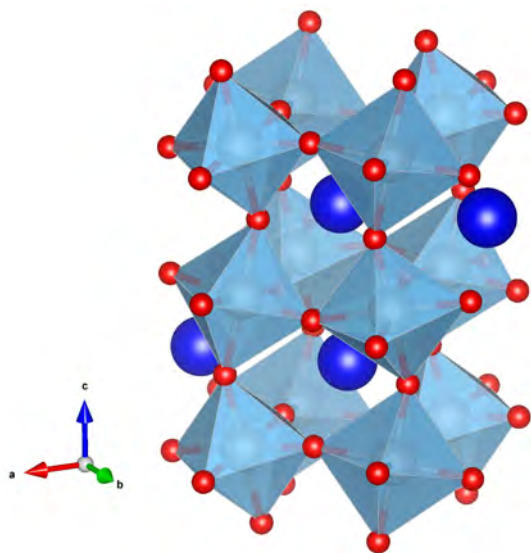
AFM $T_N=48K$, apical tilt $\Theta=17^\circ$, charge gap $\Delta=0.50$ eV

[Crandles et al., Physica C, 407 (1992)] [Okimoto et al., PRB 95, 9581 (1995)]

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LaTiO₃

AFM $T_N=146K$, apical tilt $\Theta=13^\circ$, charge gap $\Delta=0.20$ eV

GdTiO₃

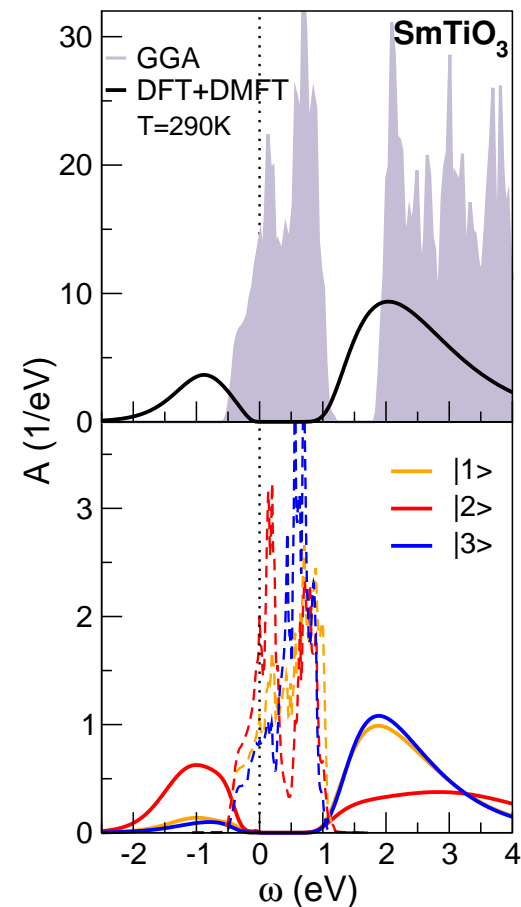
FM $T_C=36K$, apical tilt $\Theta=18^\circ$, charge gap $\Delta=0.75$ eV

SmTiO₃

AFM $T_N=48K$, apical tilt $\Theta=17^\circ$, charge gap $\Delta=0.50$ eV

[Crandles et al., Physica C, 407 (1992)] [Okimoto et al., PRB 95, 9581 (1995)]

$U=5$ eV $J_H=0.64$ eV

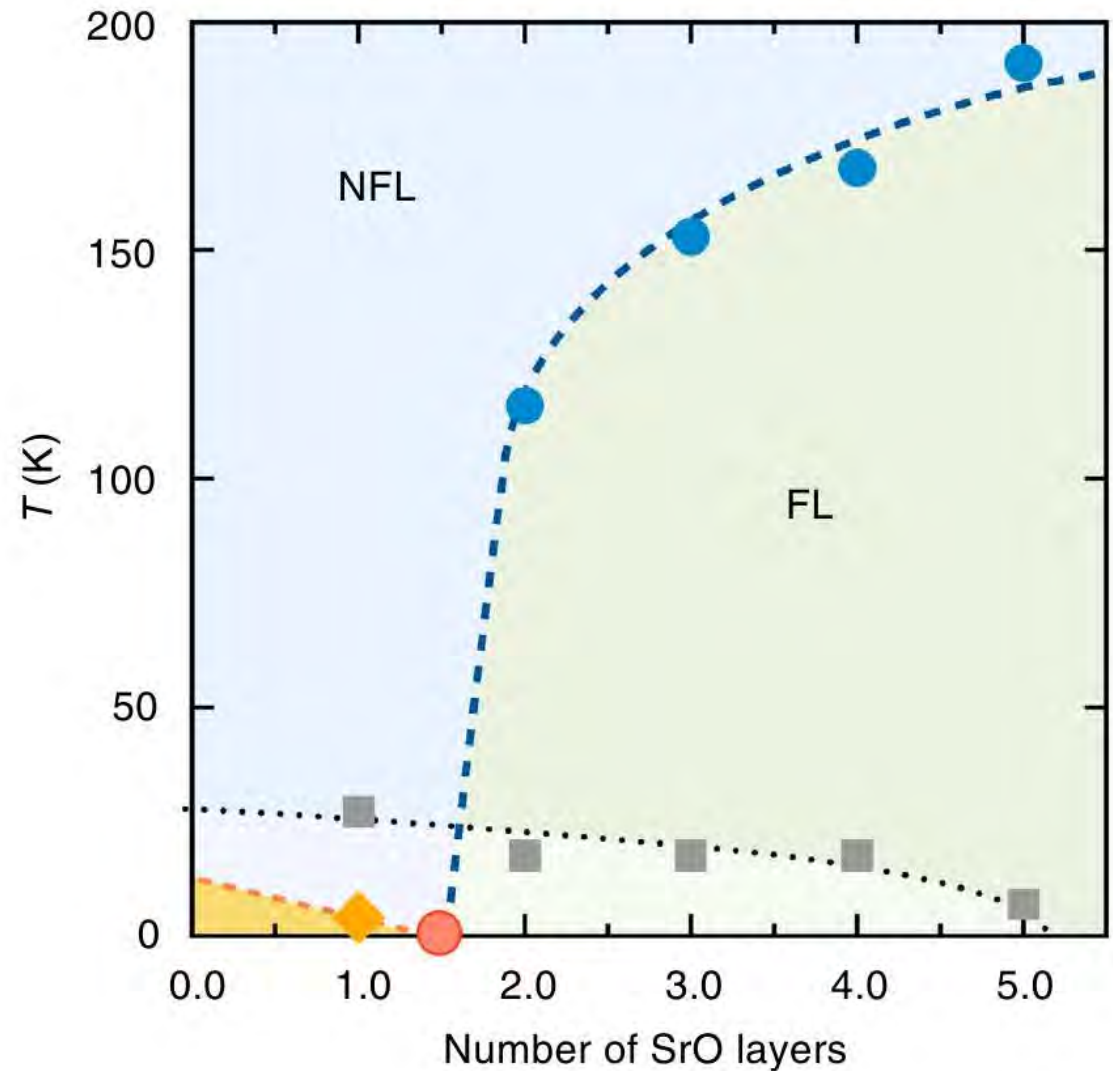
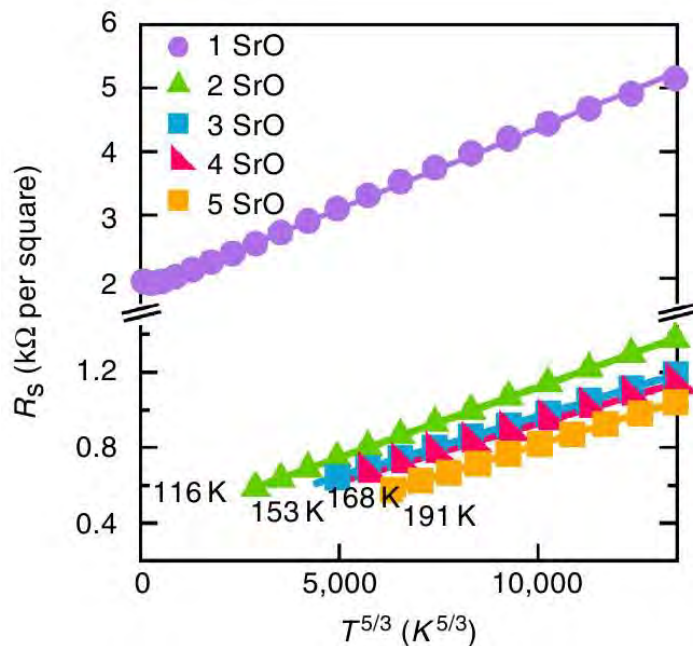
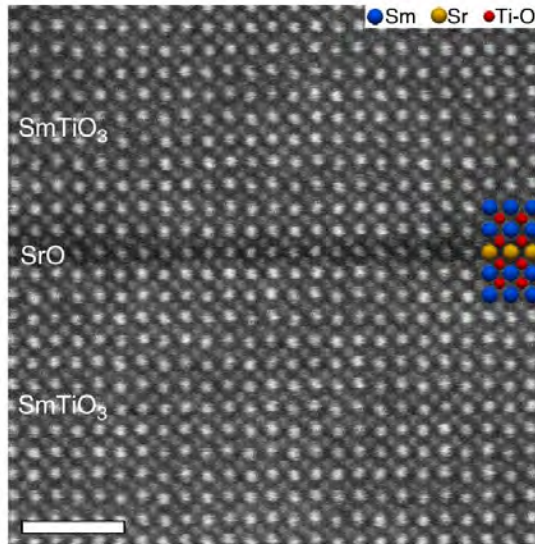


SmTiO₃ is close to an AFM-to-FM quantum phase transition

δ -doped SmTiO_3 : Experimental Non-Fermi-Liquid Behavior

[Jackson et al., Nat. Commun. 5, 4258 (2014)]

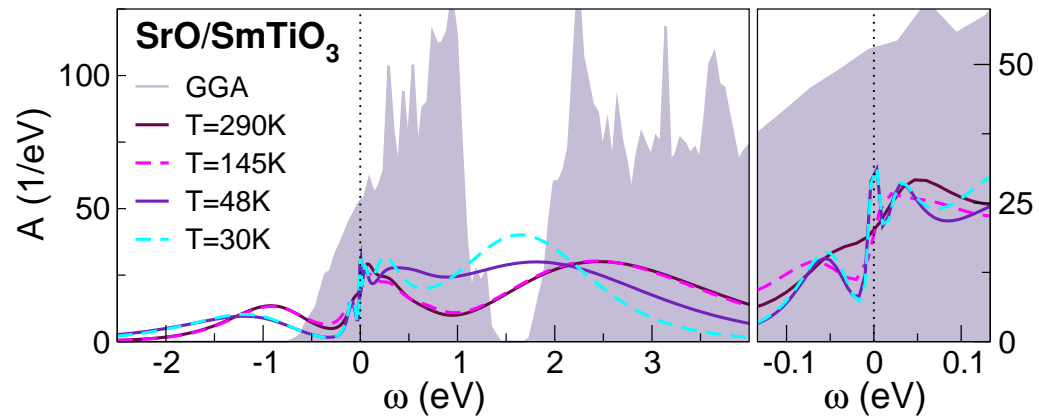
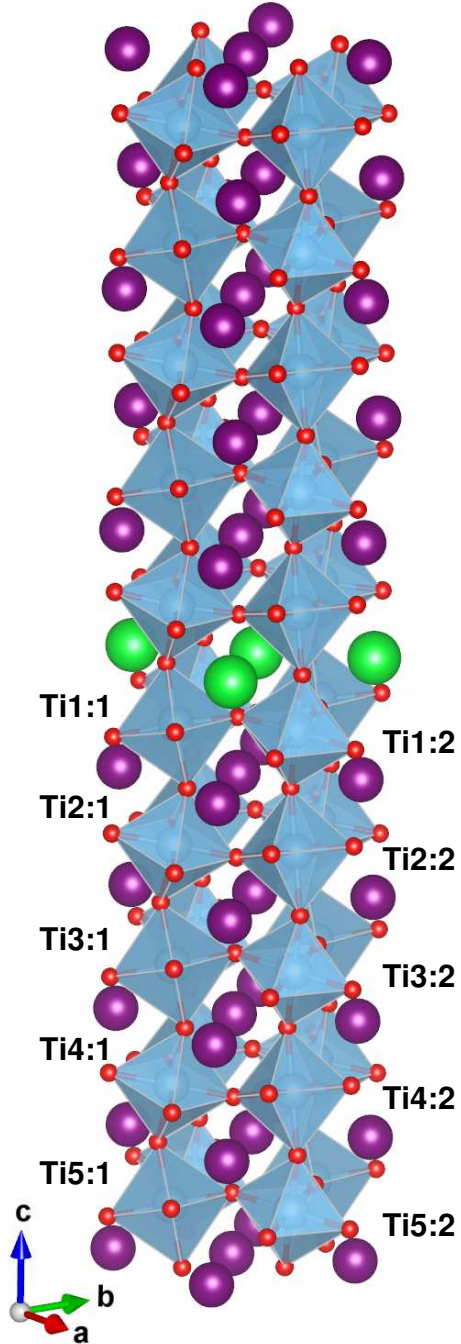
well-defined doping with layers of SrO through molecular-beam epitaxy



$T^{5/3}$ -law for the resistivity

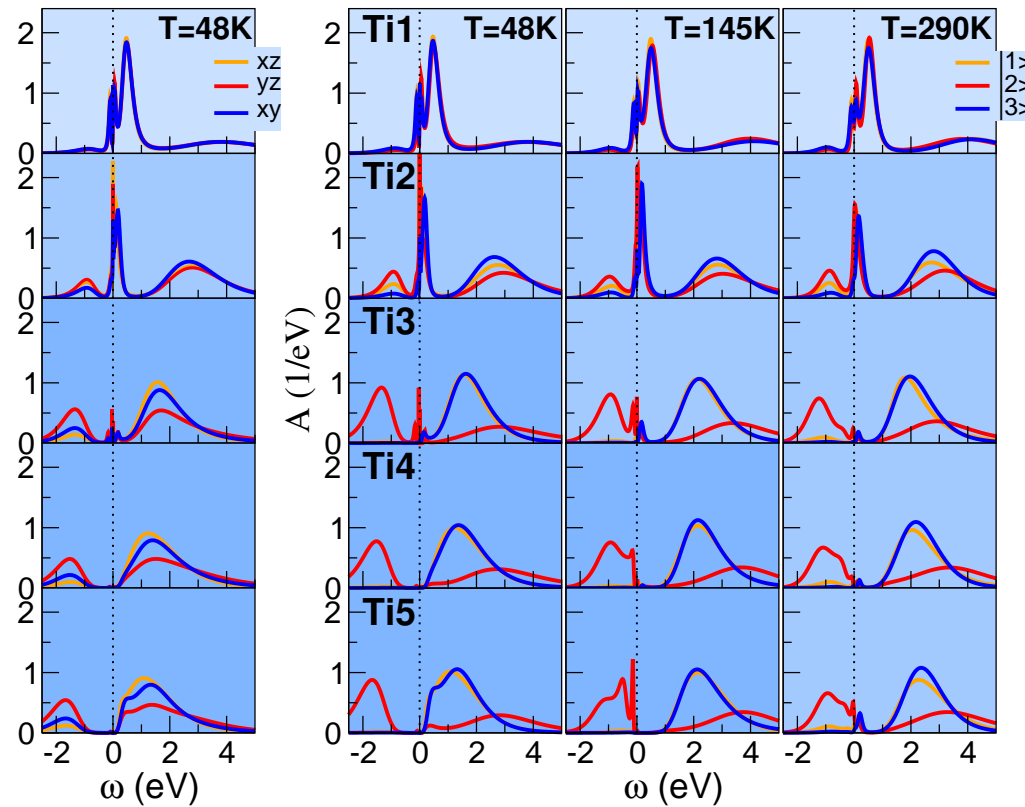
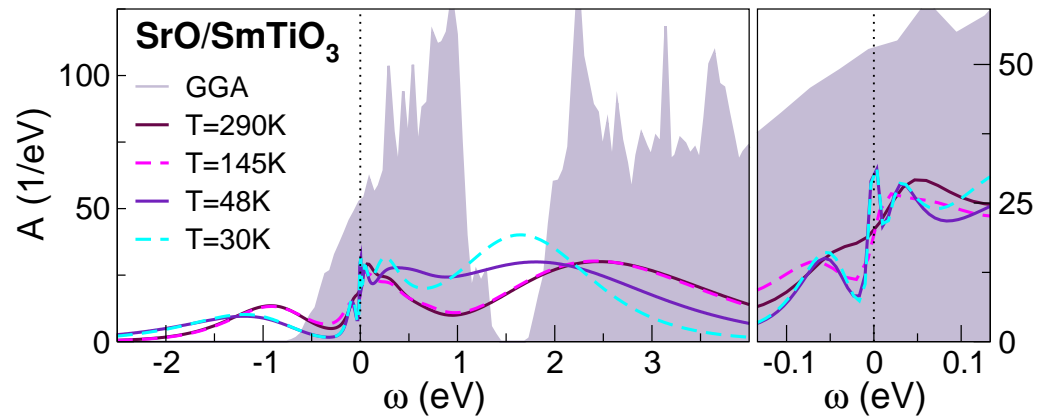
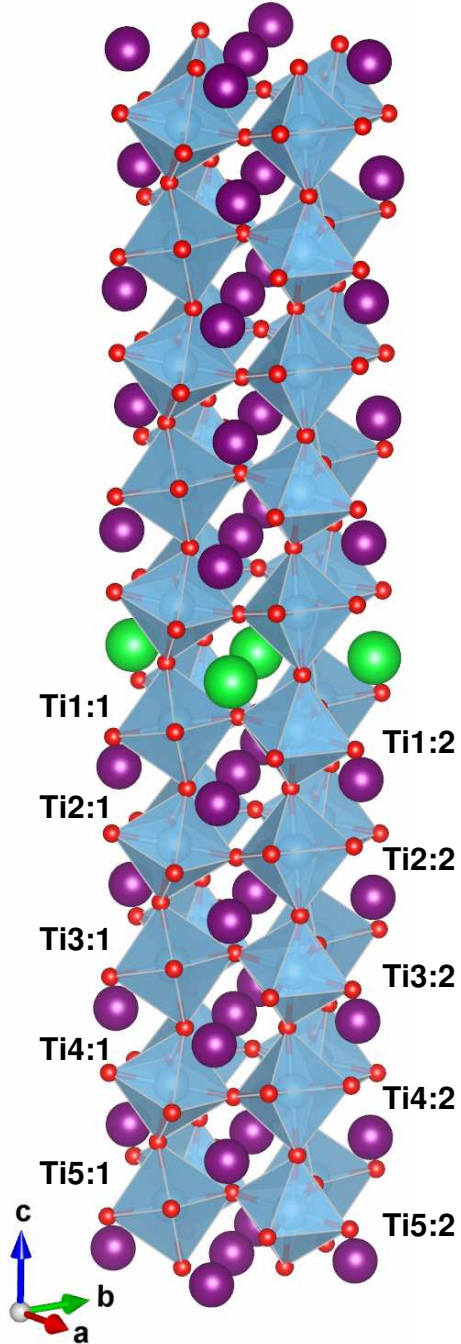
δ -doped SmTiO_3 : Spectral Function from DFT+DMFT

[FL, Sci. Rep. 7, 1565 (2017)]



δ -doped SmTiO_3 : Spectral Function from DFT+DMFT

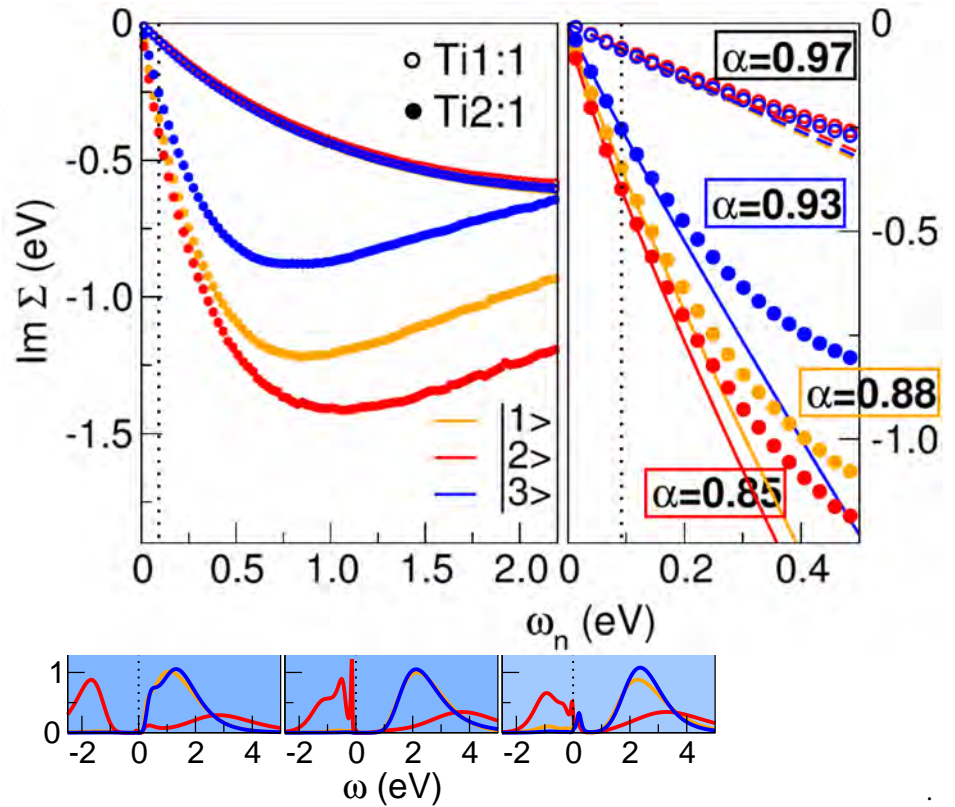
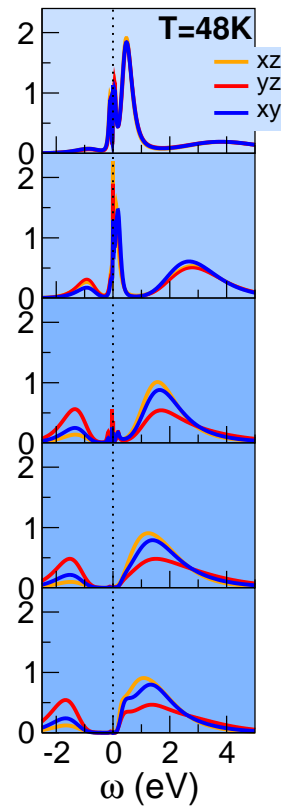
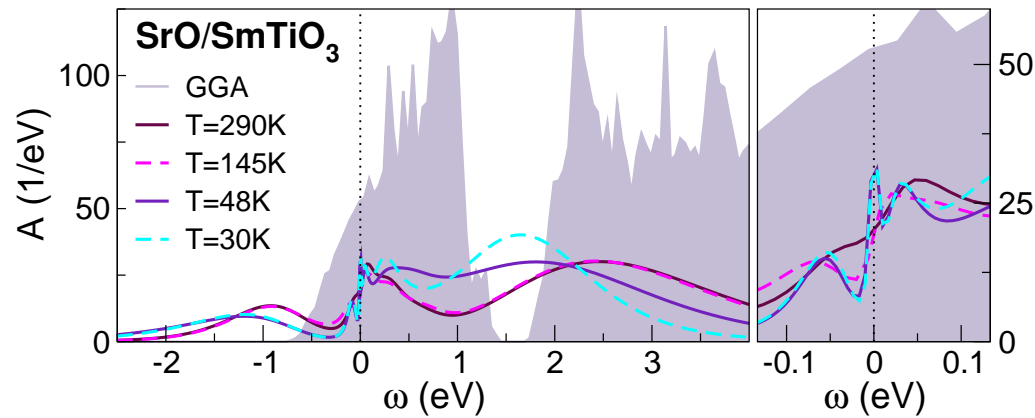
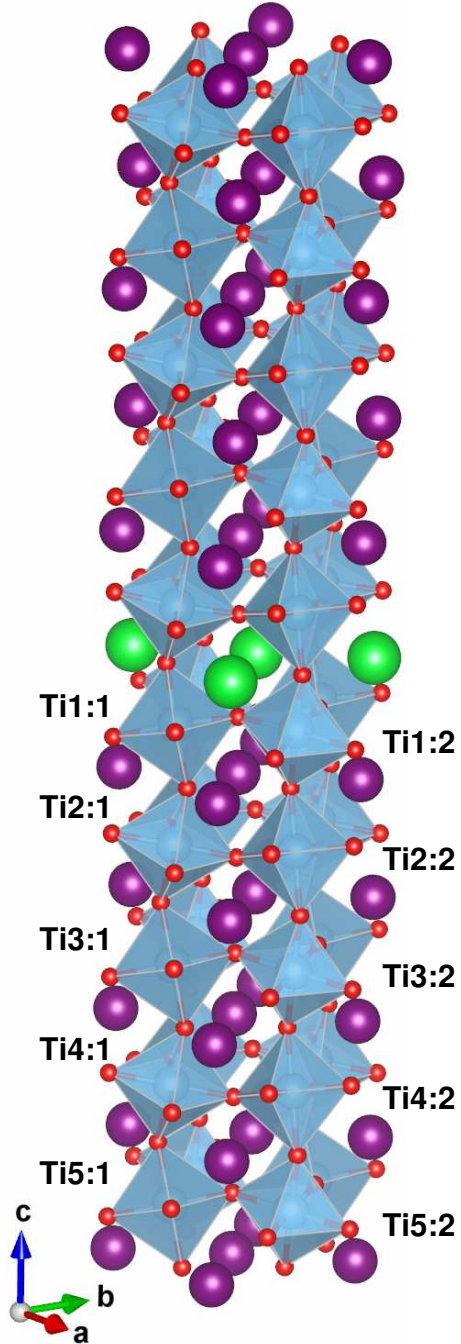
[FL, Sci. Rep. 7, 1565 (2017)]



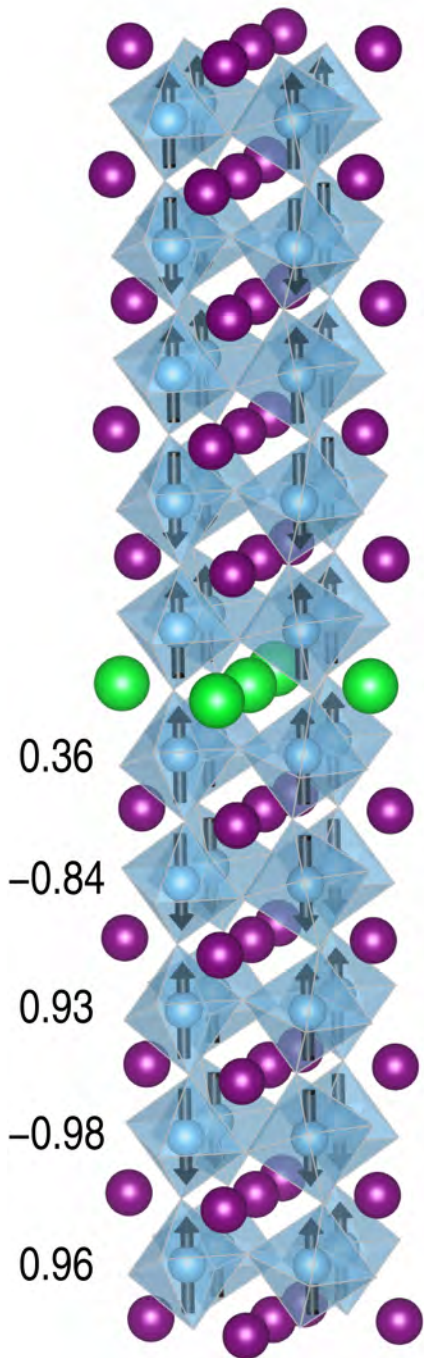
↓
layer-dependent
multi-orbital
Mott transition

δ -doped SmTiO_3 : Spectral Function from DFT+DMFT

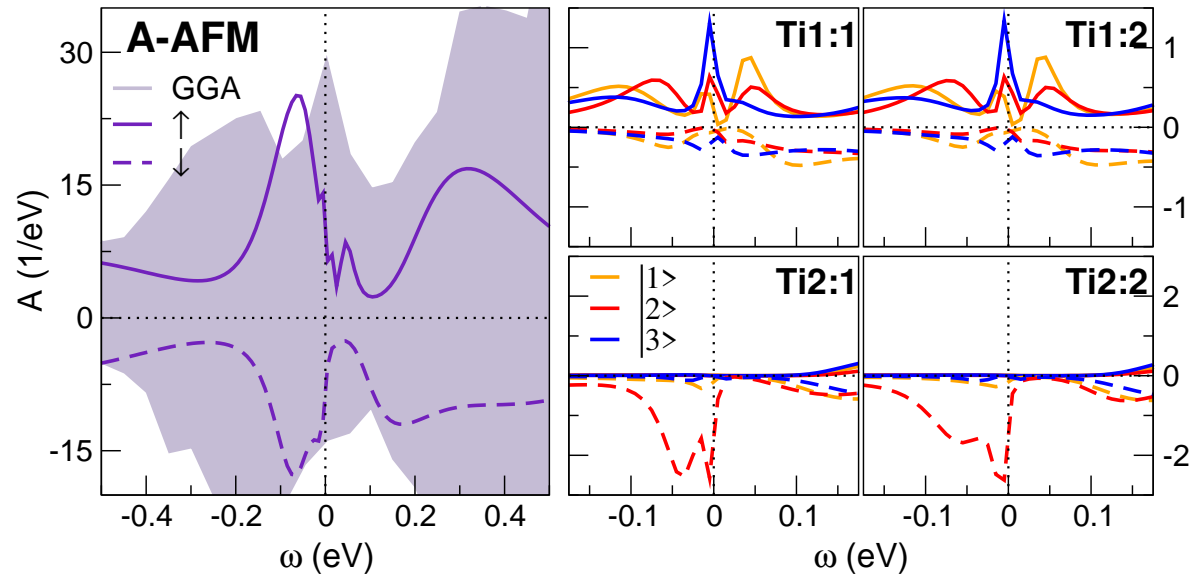
[FL, Sci. Rep. 7, 1565 (2017)]



δ -doped SmTiO_3 : Magnetic Order and Pseudogap Fingerprint

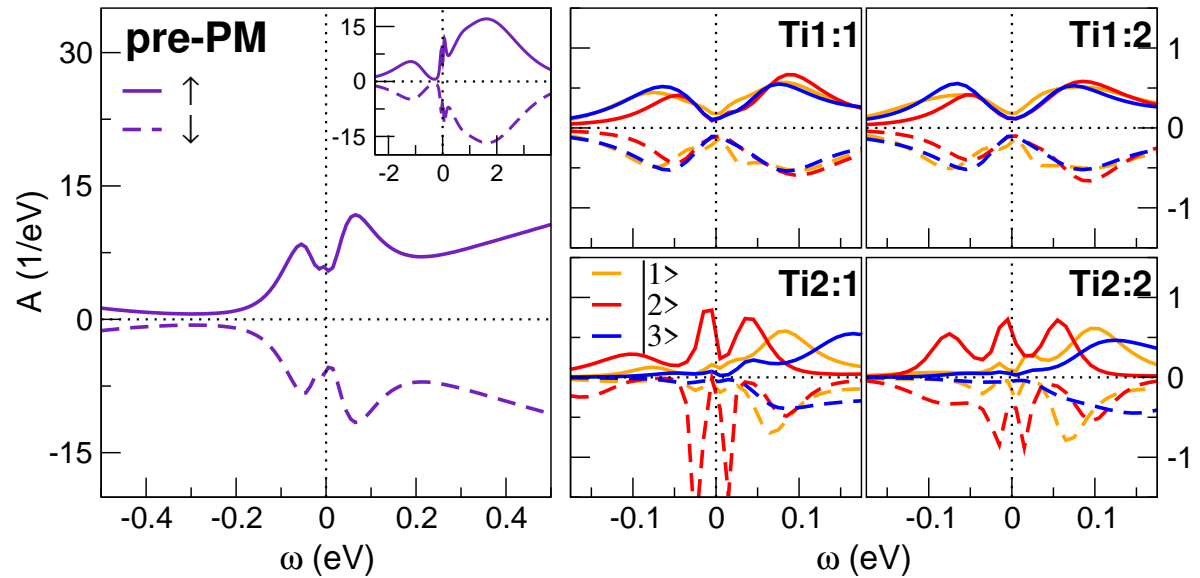


A-type AFM phase



pre-converged PM phase

with pseudogap signature

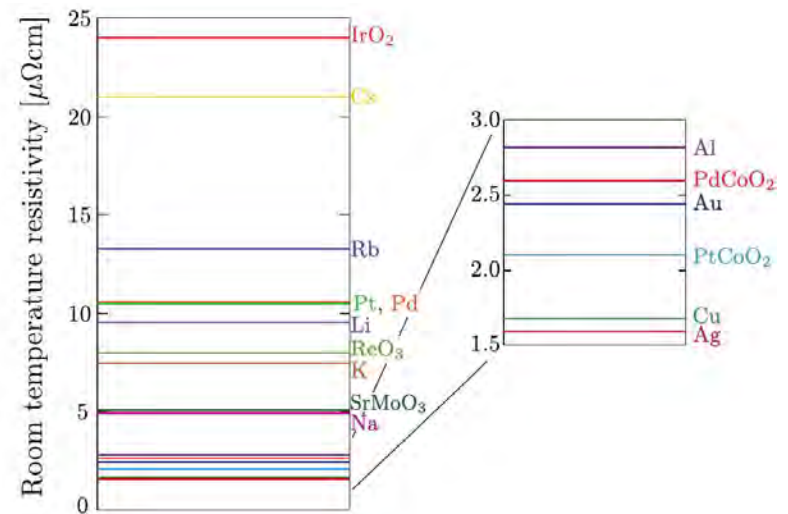
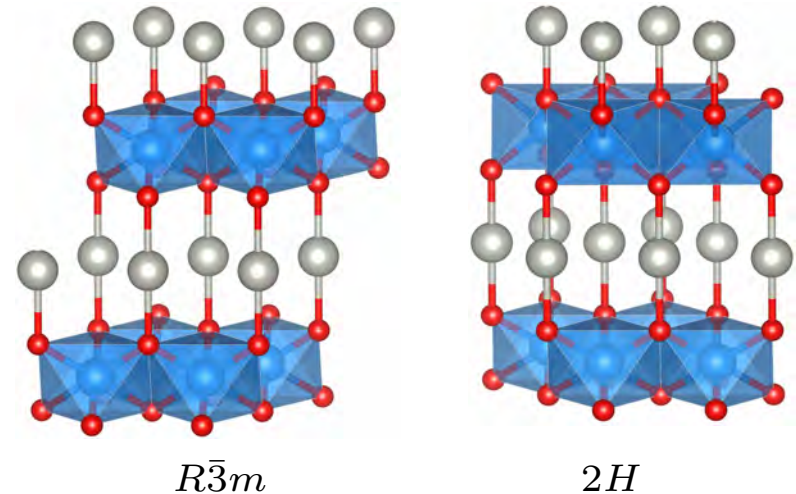


Delafossites: Introduction

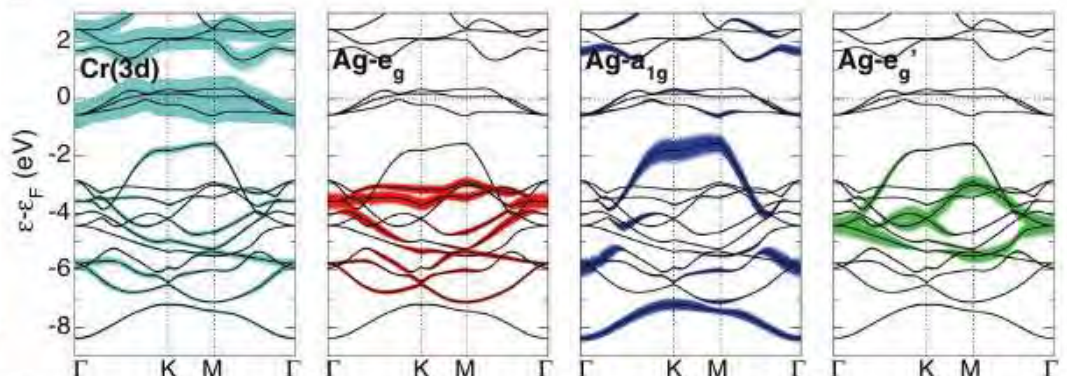
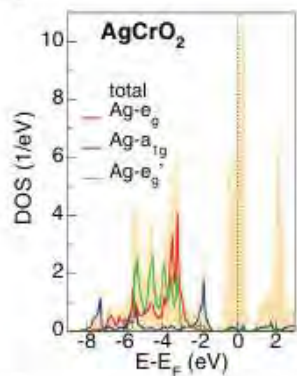
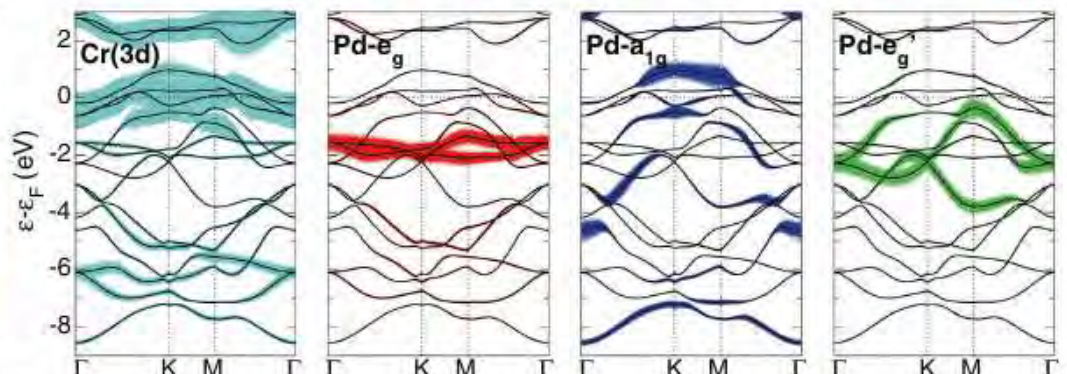
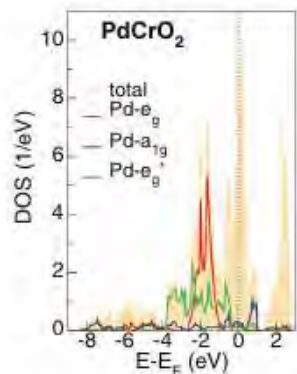
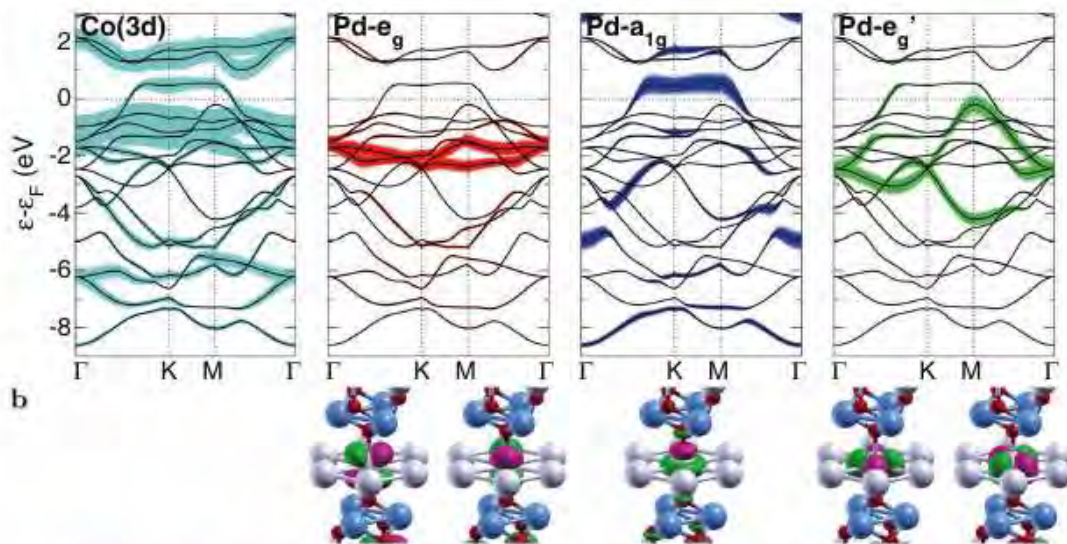
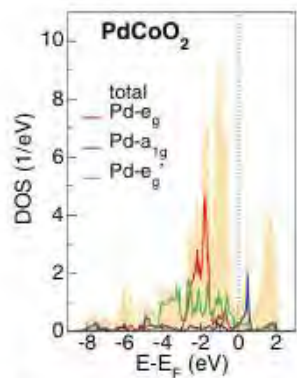
[Mackenzie, Rep. Prog. Phys. 80, 032501 (2017)]

[FL, npj Comput Mater 7, 120 (2021)]

- natural oxide heterostructure ABO_2 , with metallic A^+ ions and metallic B^{3+} ions
 - alternating A layers and BO_2 layers
 - A layer forms triangular lattice
 - unique dumbbell O-A-O structure of connecting layers
 - two stackings
 - $R\bar{3}m$: e.g. $PdCoO_2$, $CuCrO_2$, etc. (most)
 - $2H$: e.g. $AgNiO_2$
 - wide variety of compounds, insulators and metals, with metals among highest conducting materials!
- playground to design correlation physics

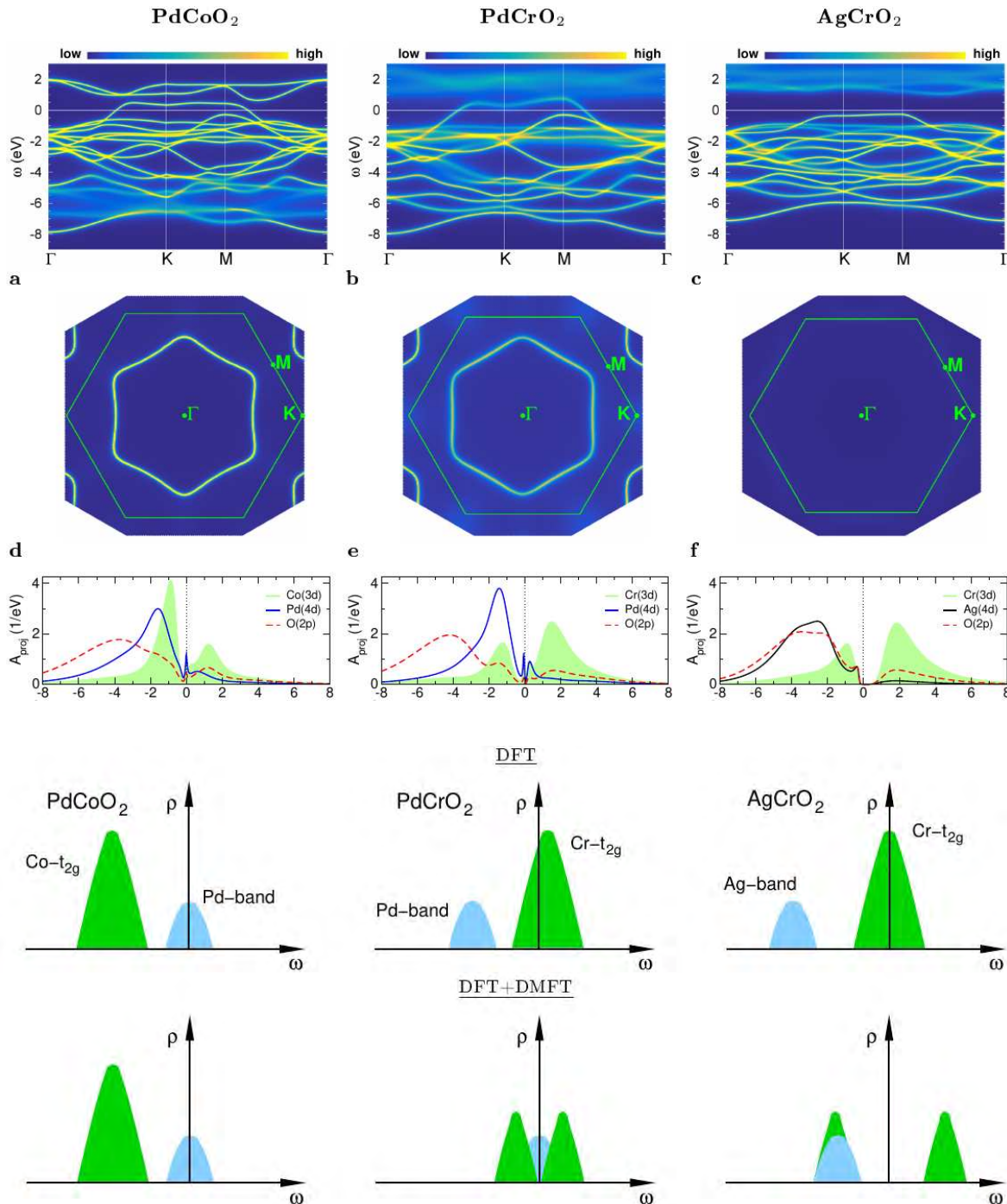


Delafossites: DFT picture

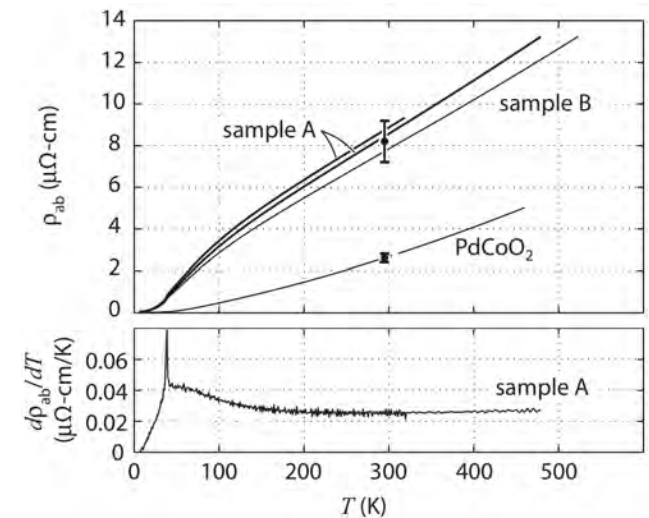


Delafossites: Correlation Scenarios

[FL, npj Comput Mater 7, 120 (2021)]



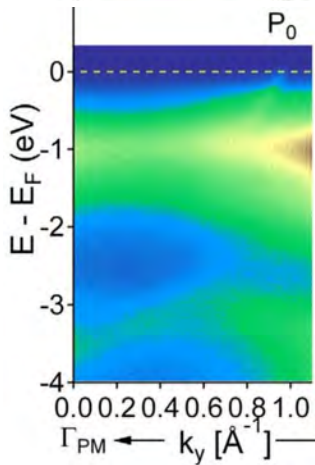
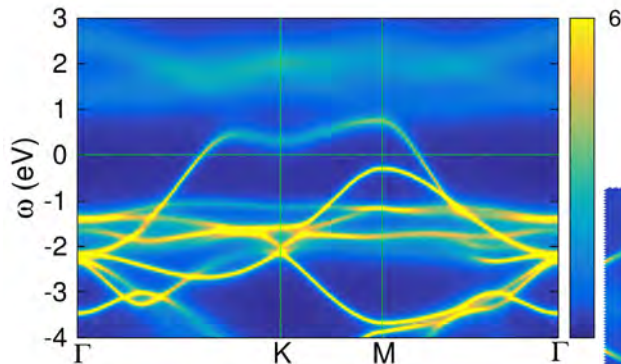
$U(\text{Co,Cr}) = 3 - 4 \text{ eV},$
 $J_H = 0.7 \text{ eV}$



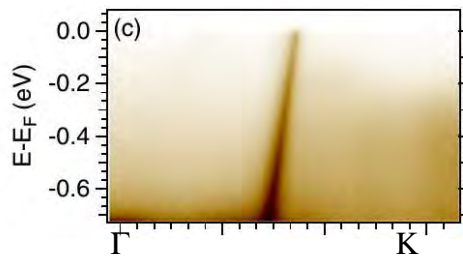
[Hicks et al, PRB 92, 014425 (2015)]

Delafossites: PdCrO₂ Correlated Electronic Structure

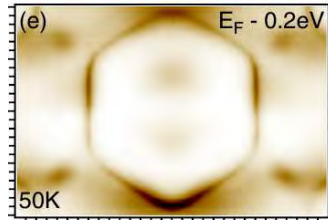
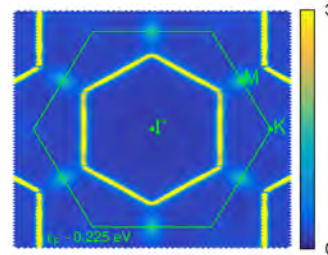
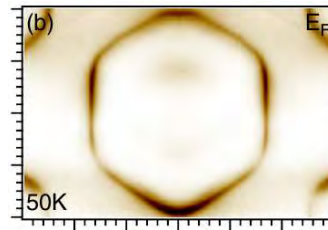
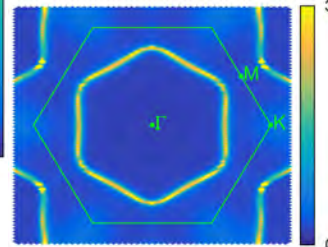
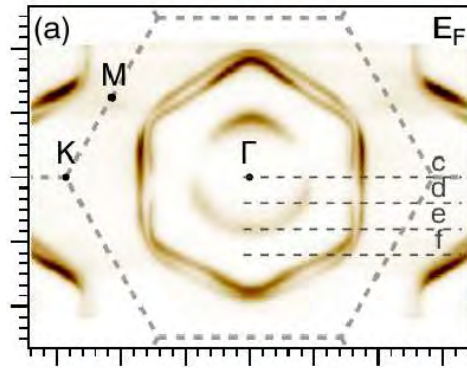
ARPES comparison in PM phase



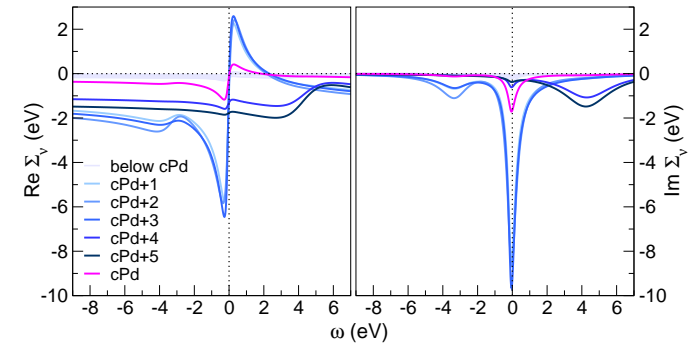
100 K [Noh et al., Sci. Rep. 4, 3680 (2014)]



10 K [Sato et al. PRB 88, 125109 (2013)]



Coupling between itinerant and localized electrons



from upfolded Cr self energy

$$\Sigma_{\nu\nu'}(\mathbf{k}, \omega) = \sum_{\mathbf{R}, m m'} \bar{P}_{\nu m}^{\mathbf{R}*}(\mathbf{k}) \tilde{\Sigma}_{m m'}^{\mathbf{R}}(\omega) \bar{P}_{m' \nu'}^{\mathbf{R}}(\mathbf{k})$$

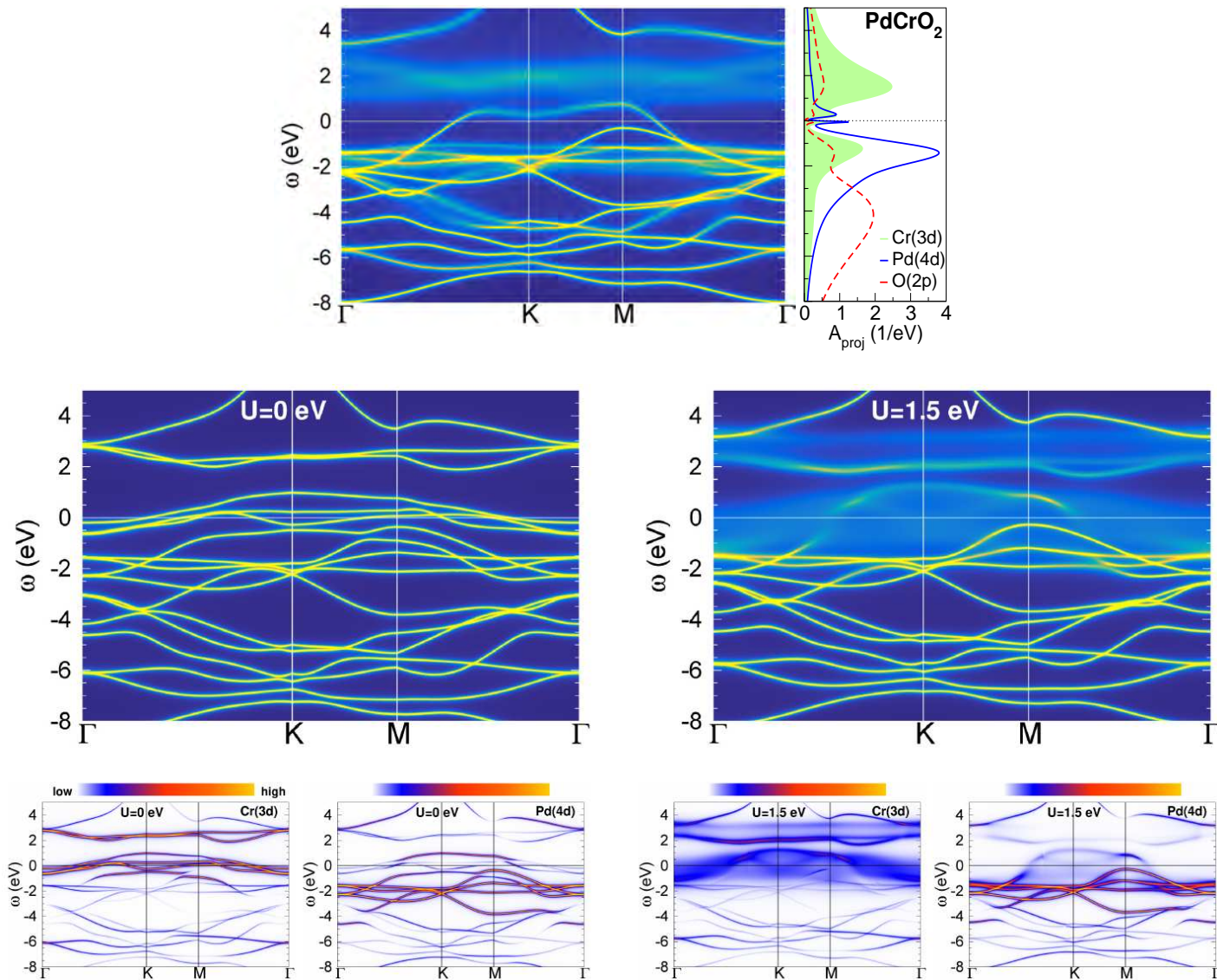
significant self-energy contribution to conducting Pd band (no Hubbard interactions on Pd)

relevant ee-scattering between Pd and CrO₂ layers

Periodic Anderson model? Kondo model?

Delafossites: PdCrO₂ from weak to strong coupling

[FL, npj Comput Mater 7, 120 (2021)]

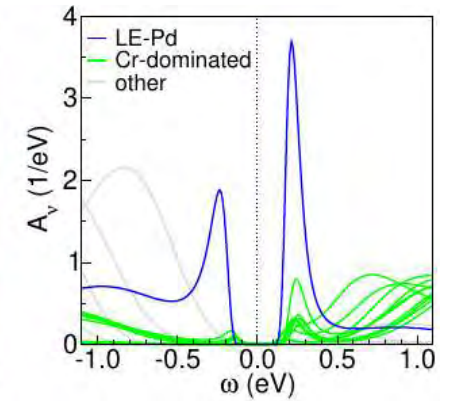
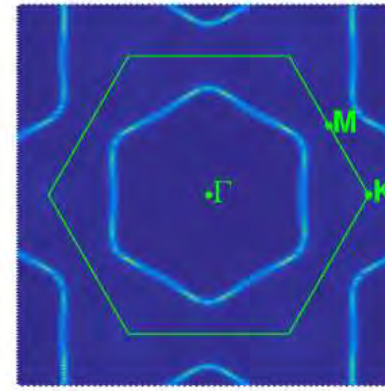
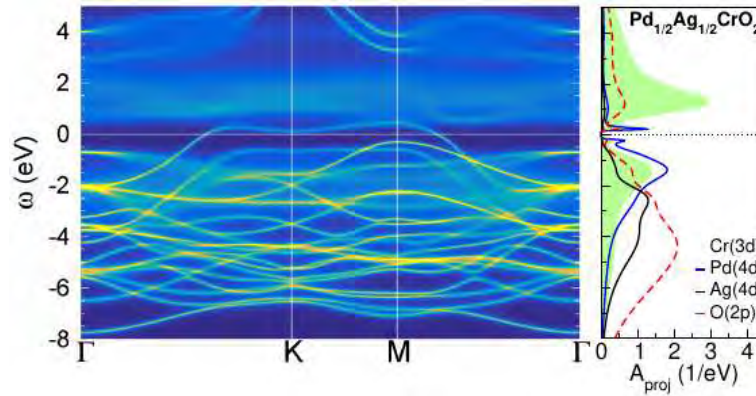
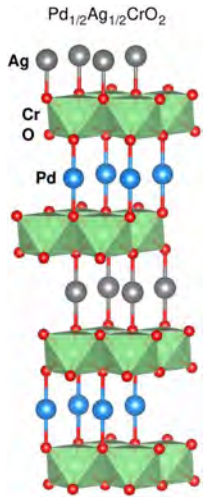


● metal-to-metal transition with growing interaction strength

● doping with isovalent Mo(4d) for Cr(3d)?

Delafossites: PdCrO₂ Mott Design I

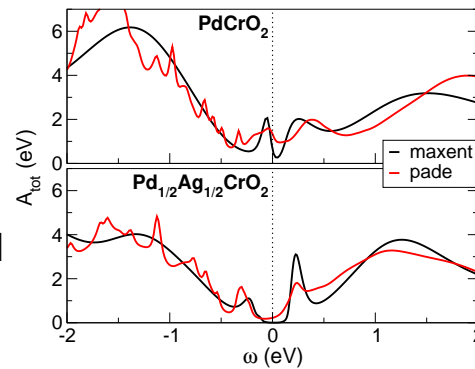
[FL, npj Comput Mater 7, 120 (2021)]



prediction

very weak dispersion crossing Fermi level

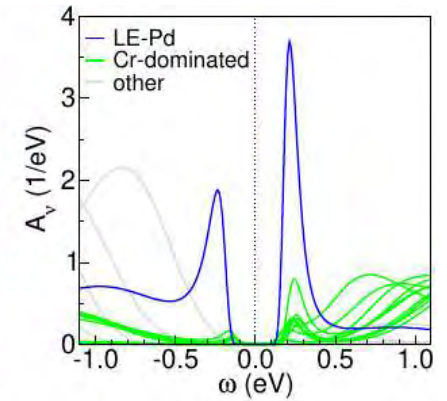
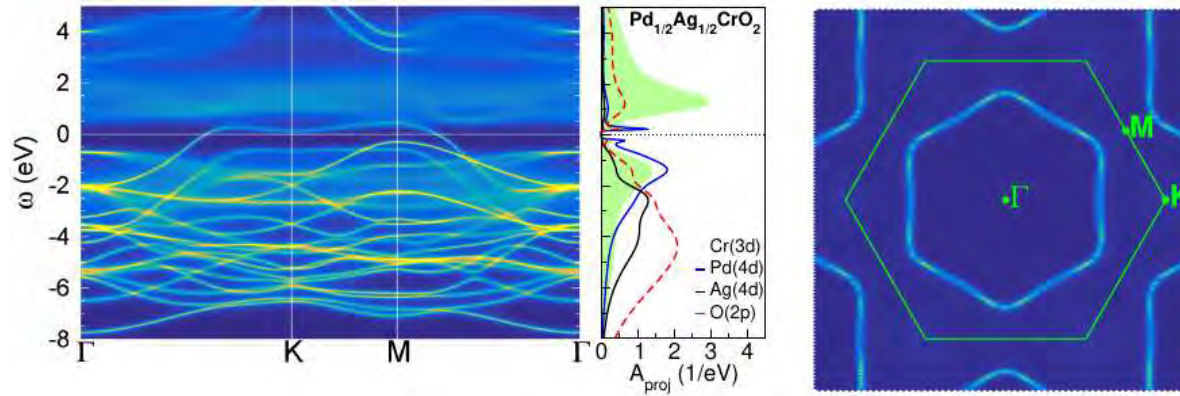
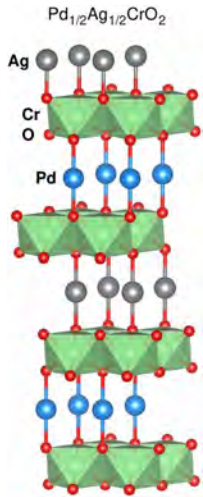
→ correlation-induced semimetal (CIS)



(no maxent artifact!)

Delafossites: PdCrO₂ Mott Design I

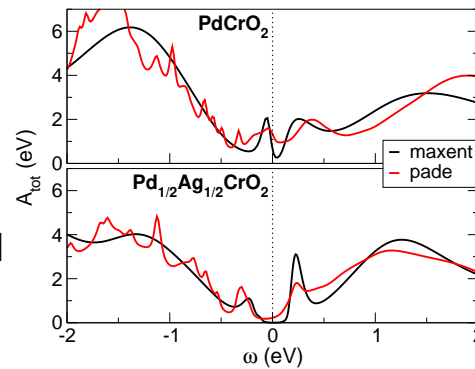
[FL, npj Comput Mater 7, 120 (2021)]



prediction

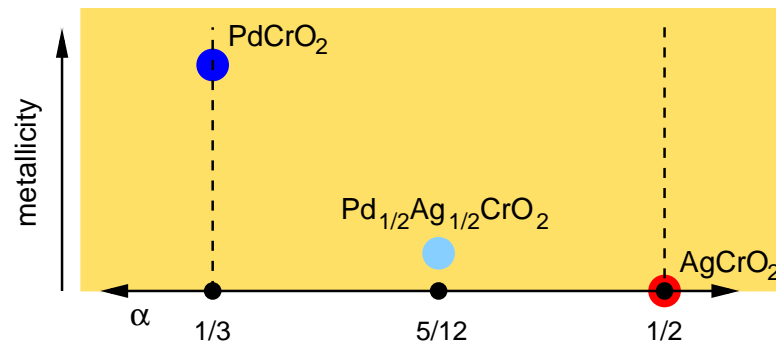
very weak dispersion crossing Fermi level

→ correlation-induced semimetal (CIS)



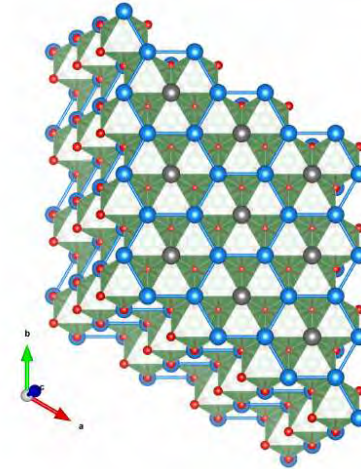
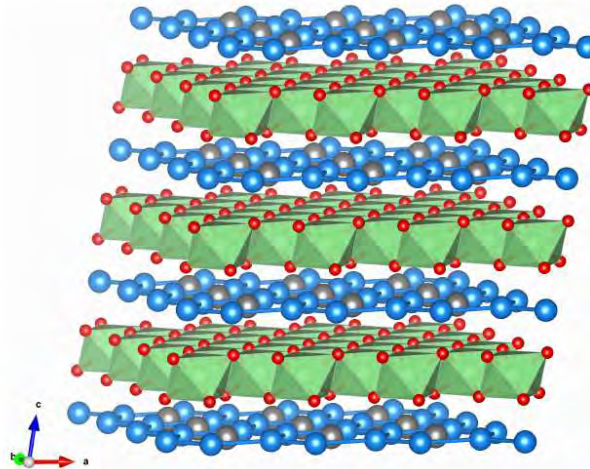
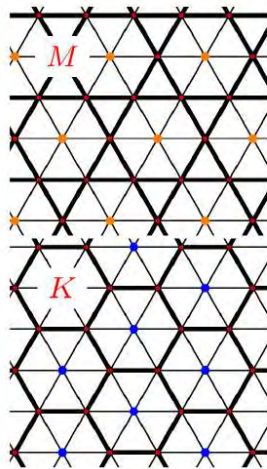
(no maxent artifact!)

strange metal inbetween hidden-Mott and Mott insulator



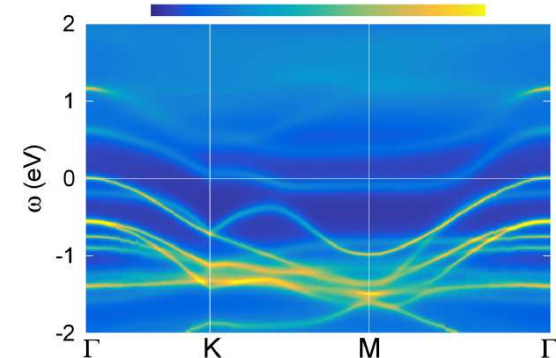
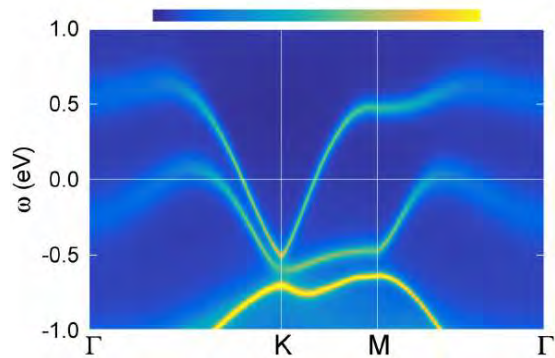
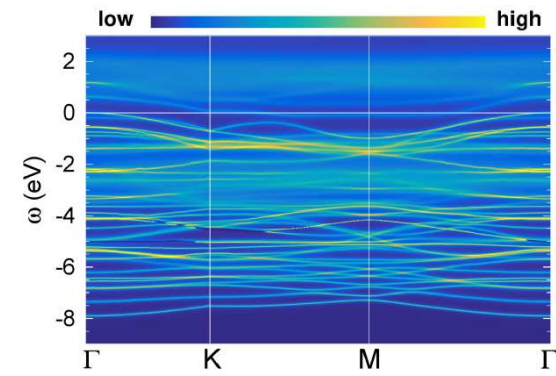
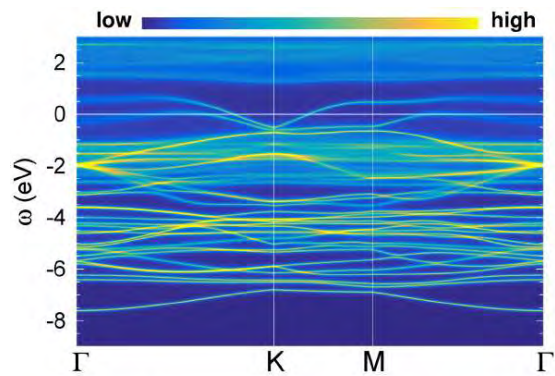
Delafossites: PdCrO₂ Mott Design II

[FL, npj Comput Mater 7, 120 (2021)]



honeycomb-type (K)

kagomé-type (M)



Summary

- DFT+Many-Body methodology is capable of addressing oxide-heterostructure challenges on a realistic level
- LaAlO₃/SrTiO₃ interface
 - **electronic correlations and defects are relevant to allow for intricate physics**
 - **2DES vs. in-gap weight dichotomy**
- δ -doping of titanate Mott insulators
 - **layer-selective Mott transitions**
 - **triggering the AFM-to-FM criticality**
- delafossites
 - **natural heterostructures (“model system”)**
 - **various design options, to be further explored**

Density Functional Theory (DFT) + Dynamical Mean-Field Theory (DMFT)

